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Improved measurements and analysis of the current profile in Tokamak fusion plasmas

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Outline

- Introduction
 - Flux surfaces and current profiles
 - Magnetic equilibrium
 - Bayesian analysis

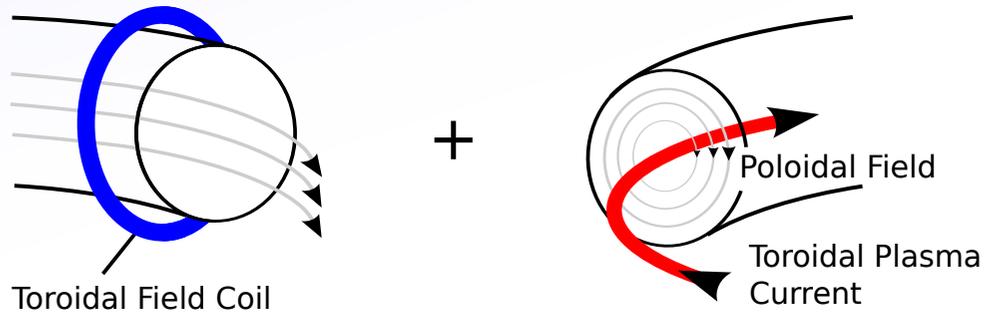
- Bayesian equilibrium
 - Current-tomography
 - Current tomography + Grad-Shafranov
 - L-Mode reconstructions
 - H-Mode results

- Internal measurements
 - Motional Stark effect.
 - Coherence Imaging
 - Imaging MSE
 - Direct j_φ imaging.

- Integrated Data Analysis
 - Current diffusion
 - Imaging MSE comparison
 - Sawtooth models

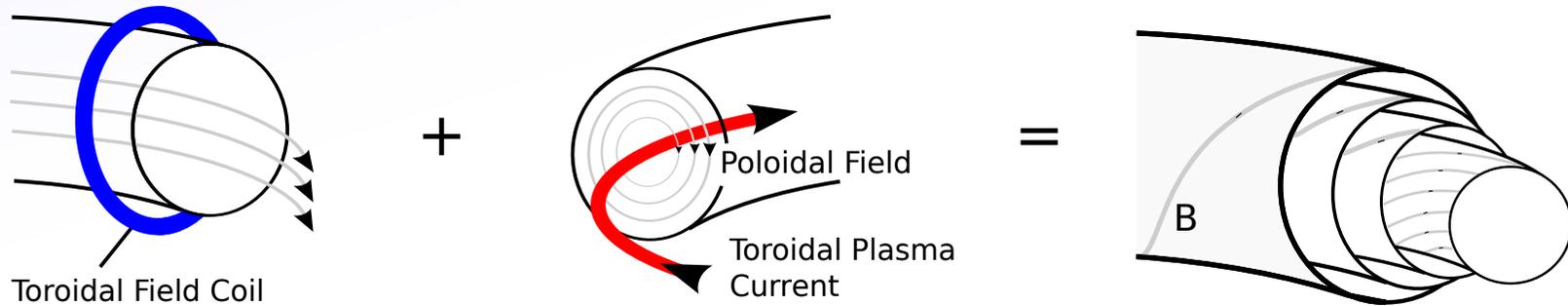
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- The Tokamak: External toroidal field coils and a large current in the plasma result in a helical magnetic field.



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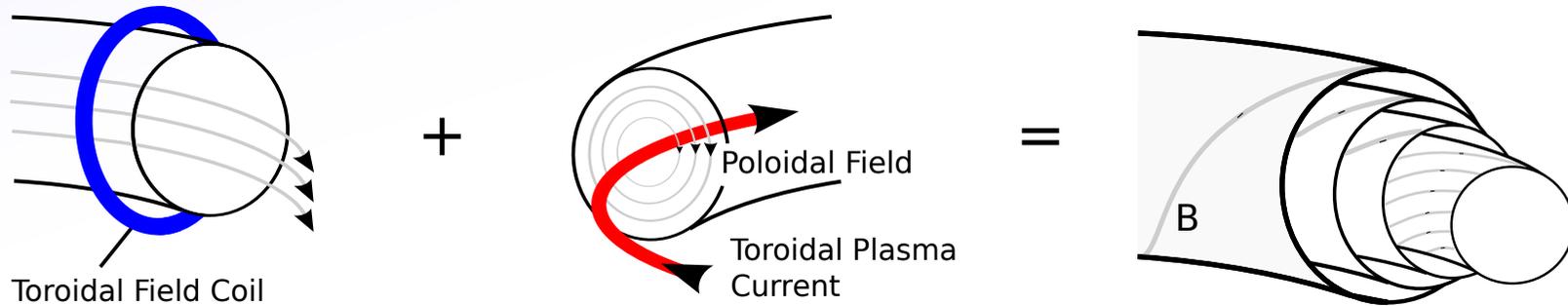
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- Flux surfaces are the basis of our knowledge and used for:
 - Comparing/combining measurements ('mapping')
 - Basis of 1D transport calculations

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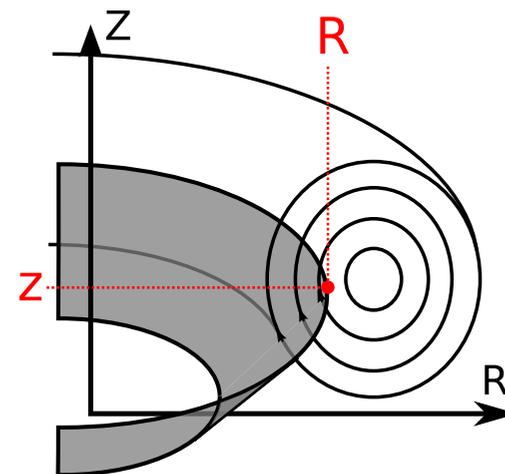


- Field lines form surfaces of constant magnetic flux

Poloidal magnetic flux:

$$\psi(R, Z) = \int_{0,0}^{R,Z} B_{\theta}(\mathbf{r}) dr$$

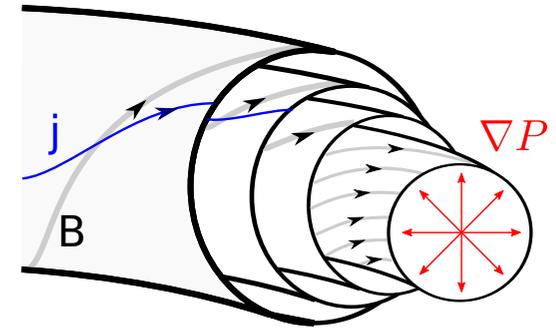
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Current and stability

- The field and current balance the kinetic pressure:

$$j \times B = \nabla P$$

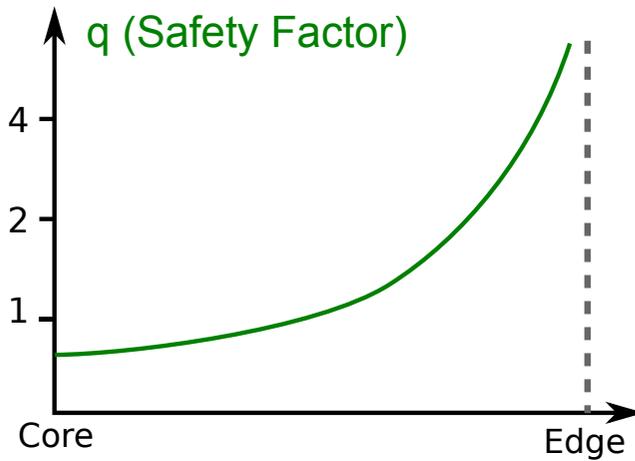
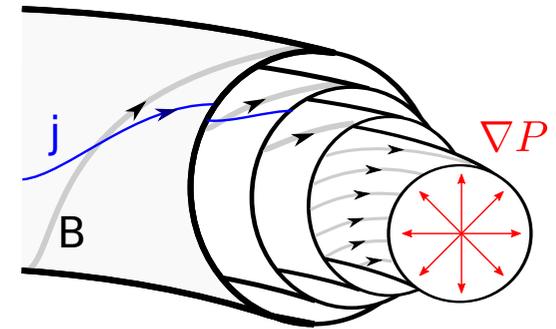


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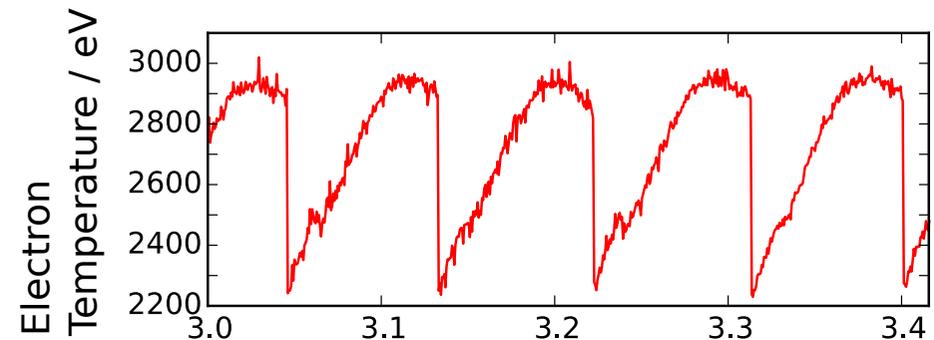
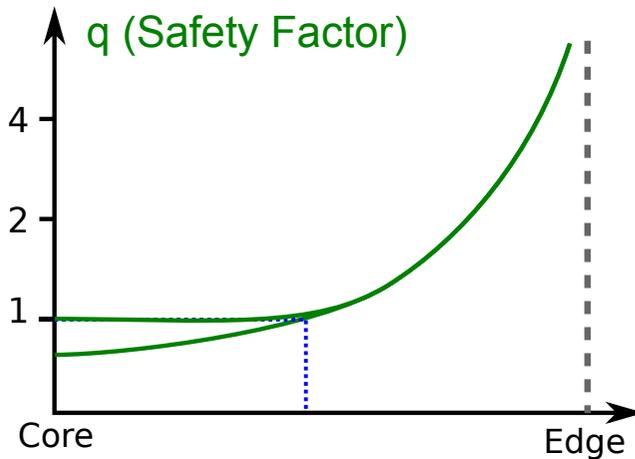
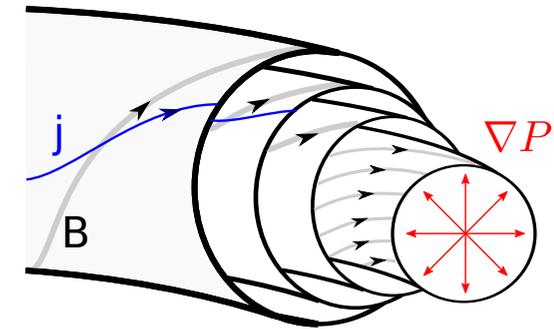


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e.g.: When the central q value falls below 1.0, the plasma core periodically suddenly expels particles and energy - known as a 'sawtooth' crash. The crash is a magnetic reconnection event, which occurs far more rapidly than explained by simple theoretical models.

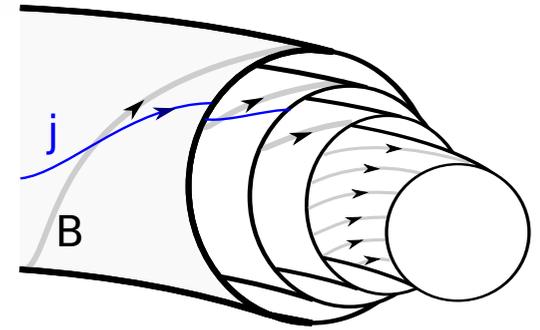
...(we'll return to this later)

Magnetic Equilibrium

- How do we know \mathbf{j} and \mathbf{B} ?

Assume: Axisymmetry + Isotropic pressure + No flow
Define the 'poloidal current flux' f .

$$f(R, Z) = \int^{R, Z} j_{\theta}(\mathbf{r}) d\mathbf{r} \qquad f = RB_{\phi}$$



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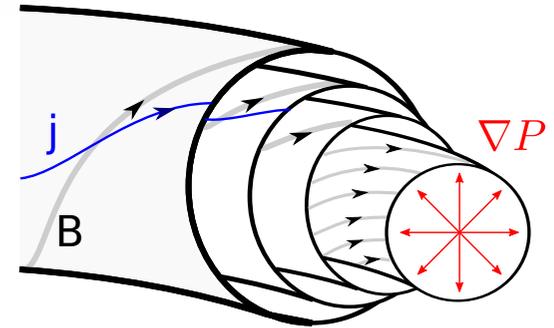
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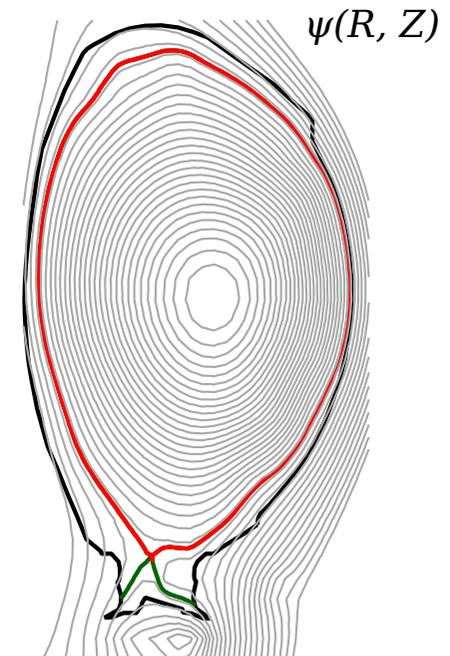
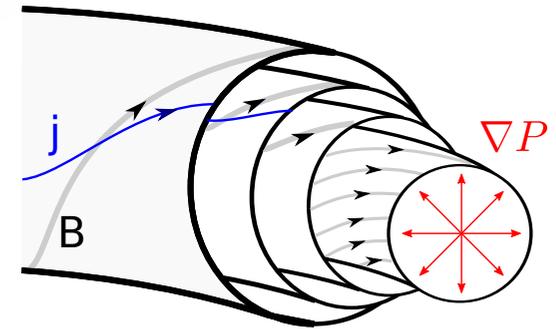
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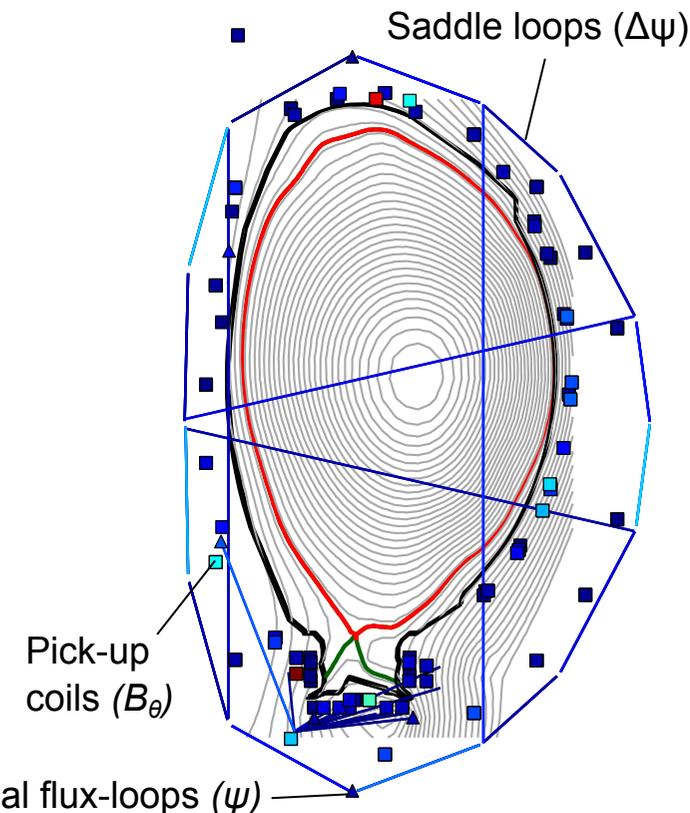
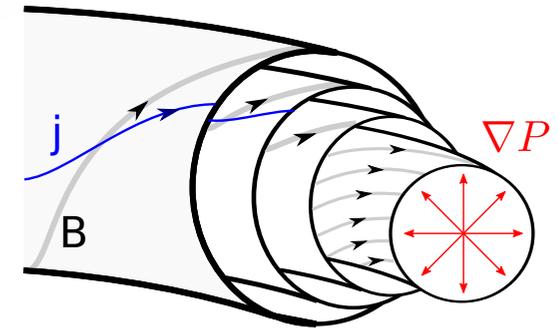
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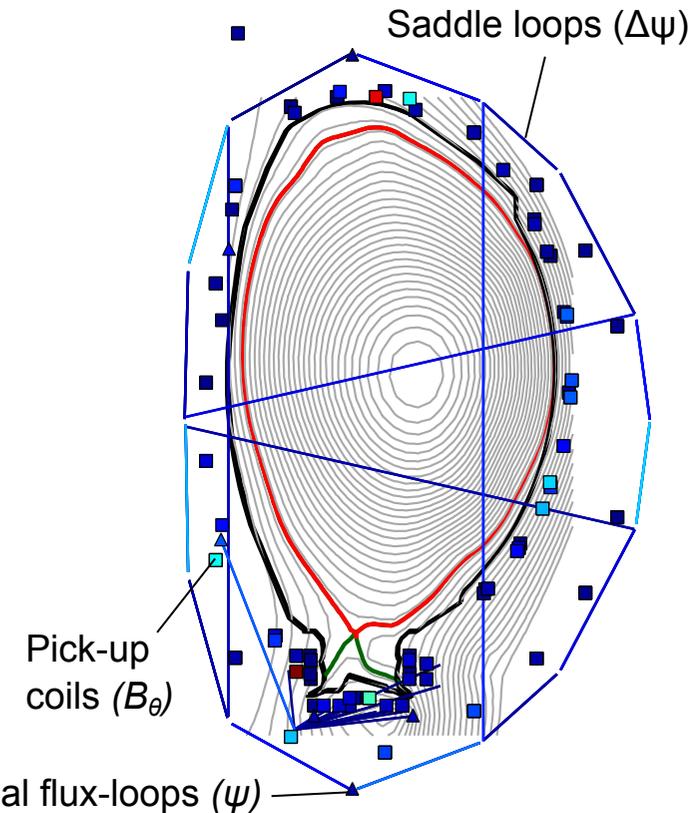
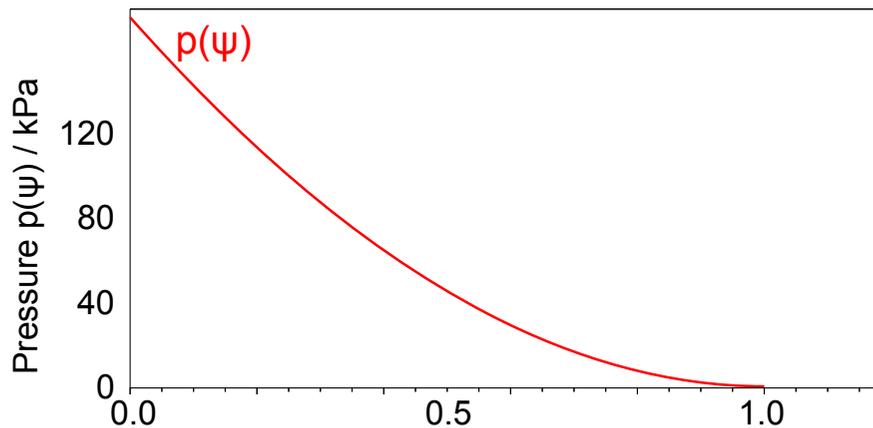
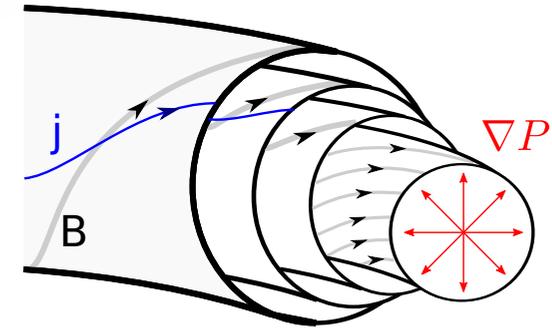
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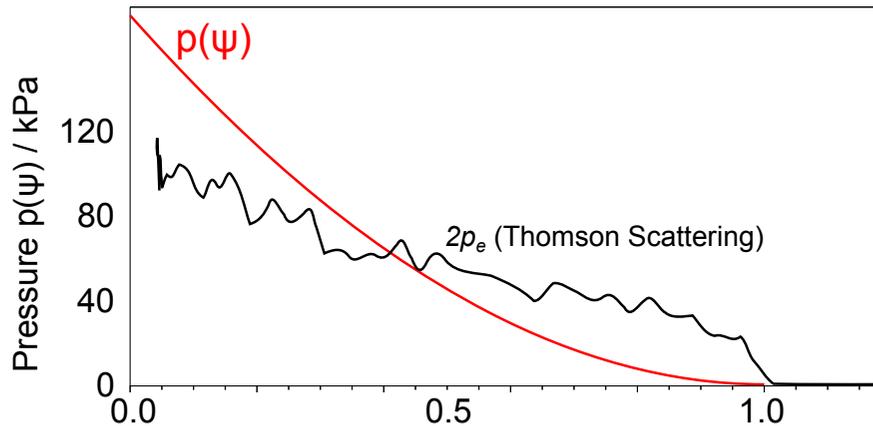
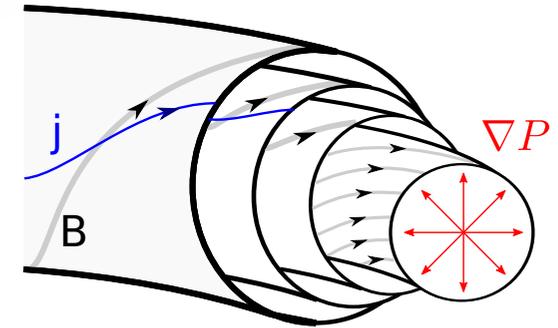
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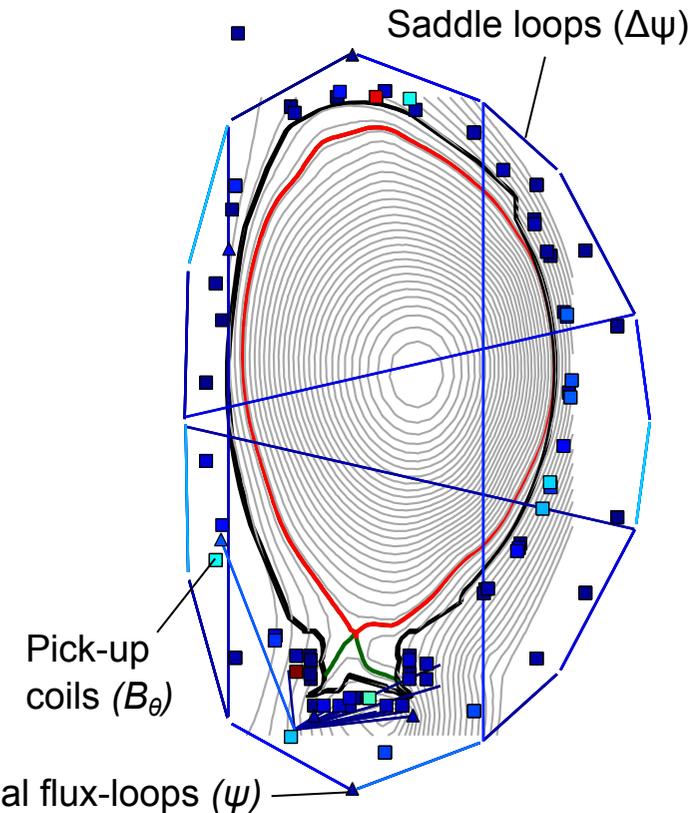
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but...

- Is the converged solution the only solution?
- Are the simplified p, f profiles over-constrained / under-constrained?



Are the data consistent with the assumptions?



Bayesian Inference

- Rigorous framework for dealing with the question:

What can we know about the plasma, given the data we measured?



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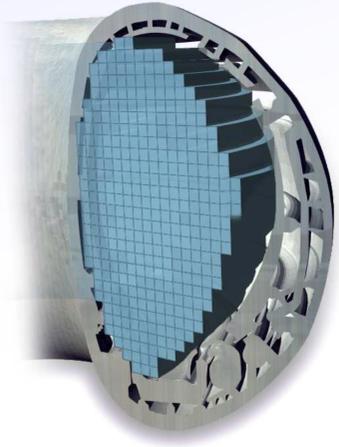
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Any combination of diagnostics: $P(D | \mu) = \prod_i P(D_i | \mu)$

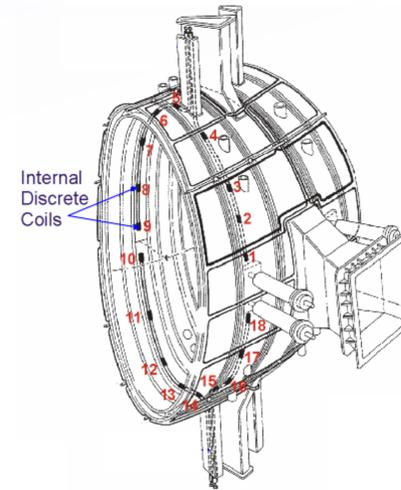
(More explanation and examples available for Q&A)

Current Tomography

- How do we apply this to current distribution?



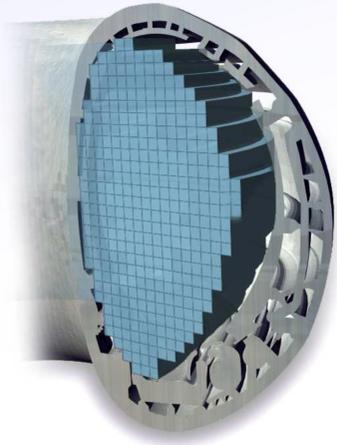
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grid of axisymmetric
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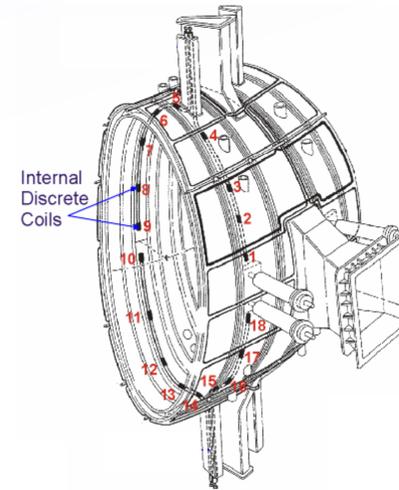
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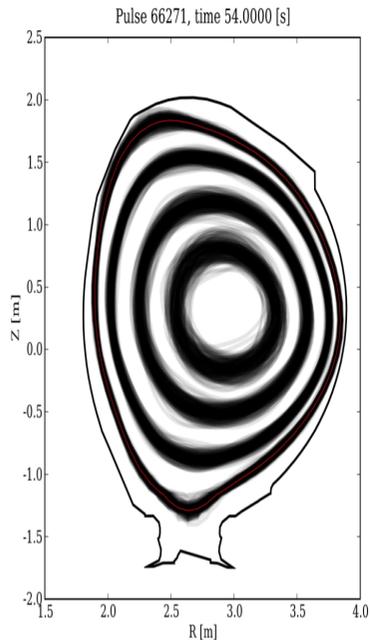


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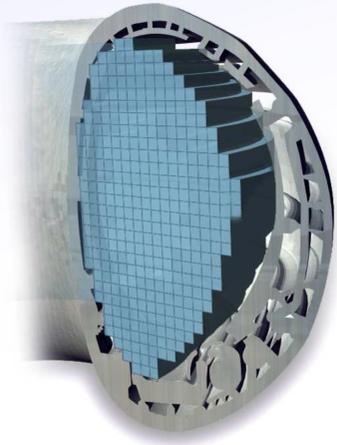
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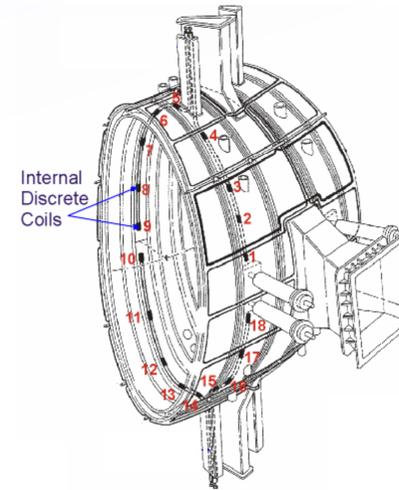
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<--- Samples of flux surfaces
Shows uncertainty and degeneracy
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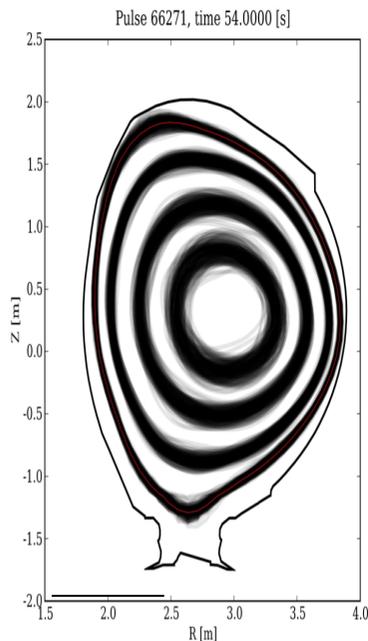


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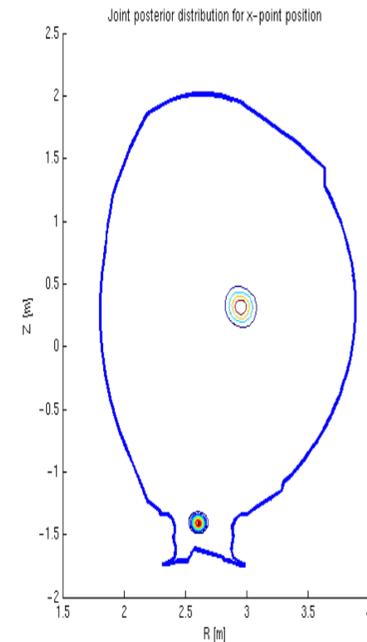
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However, some quantities are well known --->
- Axis position
- X-point position





Current Tomography + Equilibrium

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Sum over current beams

How good should our equilibrium be?

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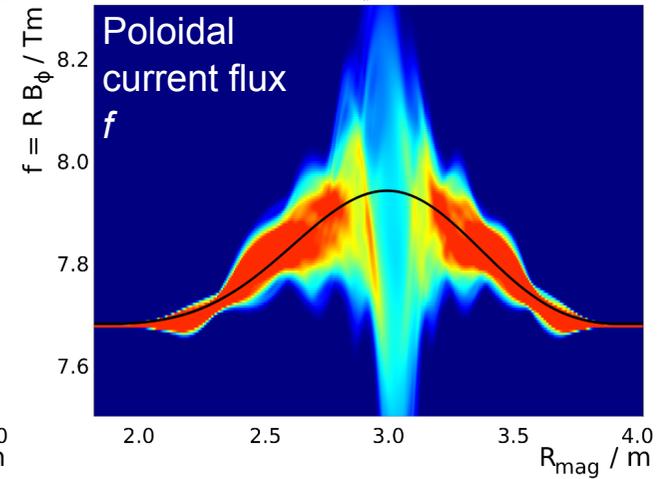
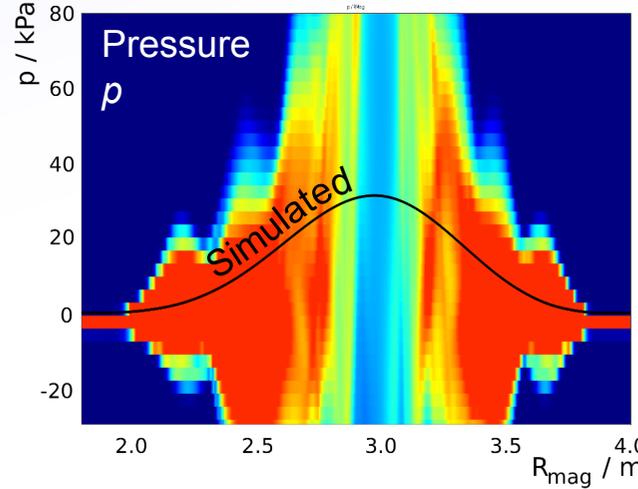
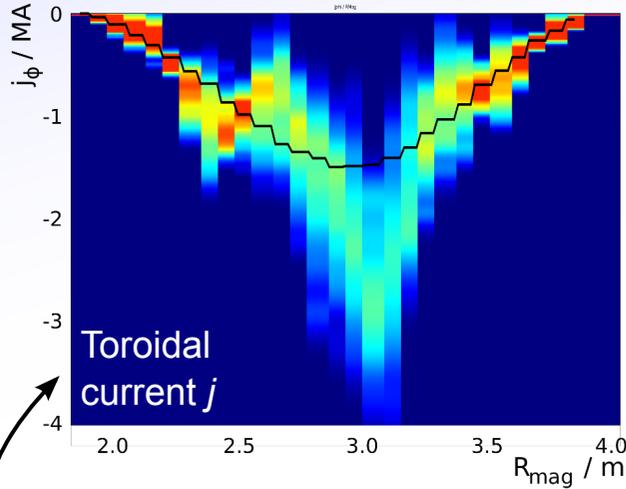
- Now we can ask the question:

$$P(j_\phi, f, p \mid D_m, \text{stable})$$

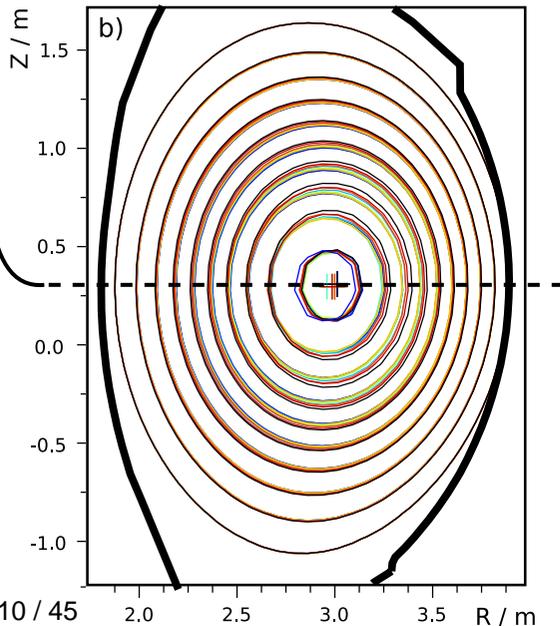
What space of plasma currents and pressures are consistent with the measurements and are close to equilibrium?

Bayesian Equilibrium

L-Mode plasmas: Low resolution current beam grid, fully explored posterior distribution:

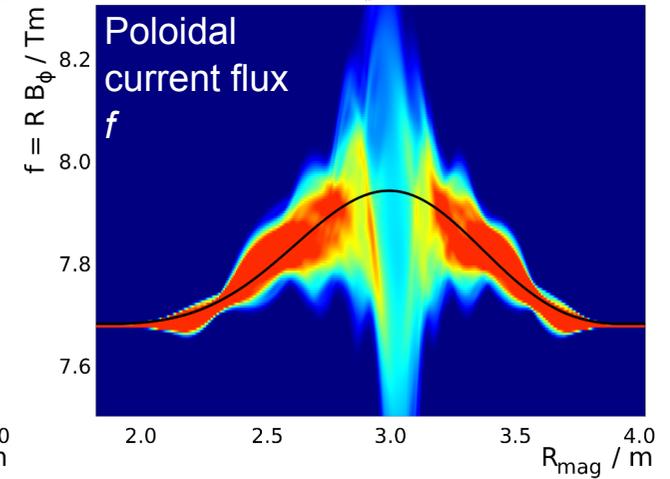
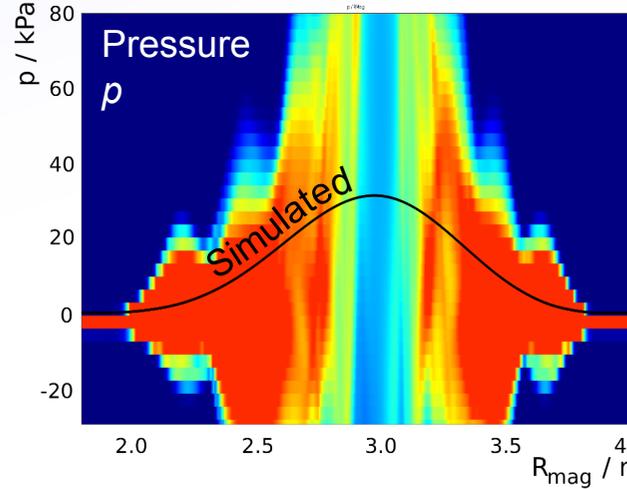
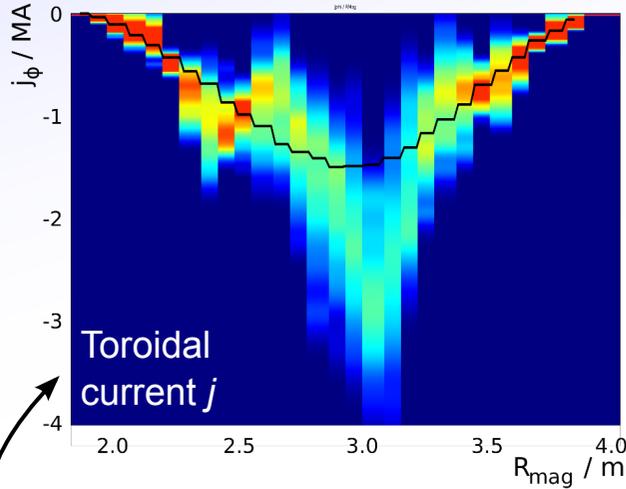


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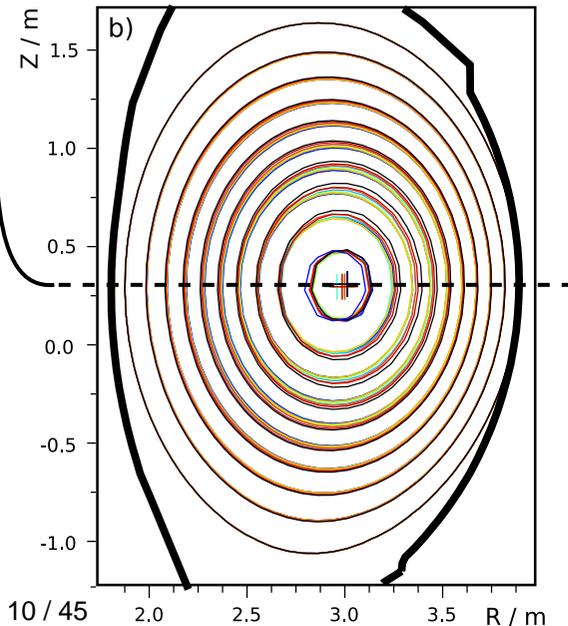


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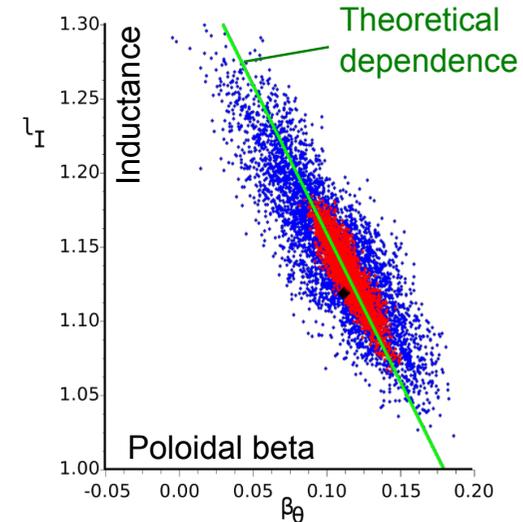
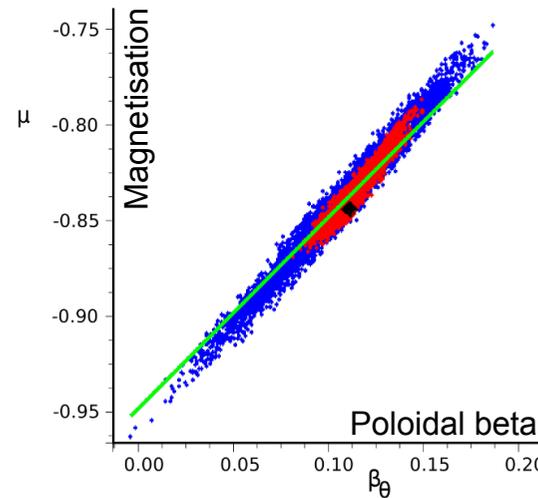
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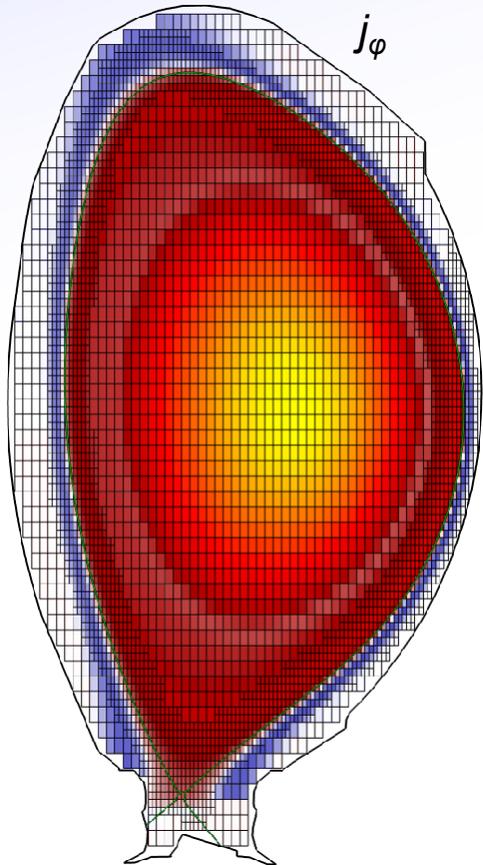


Samples of integral quantities reveal relations between 'Shafranov Integrals' that are well determined in simple analytical equilibrium solutions:



Bayesian Equilibrium

H-Mode plasmas: Very sharp changes in j and p require high-resolution current beam grid:



Too many parameters to explore the posterior (Monte-Carlo algorithm). Needs:

- More computation power
- Better algorithms

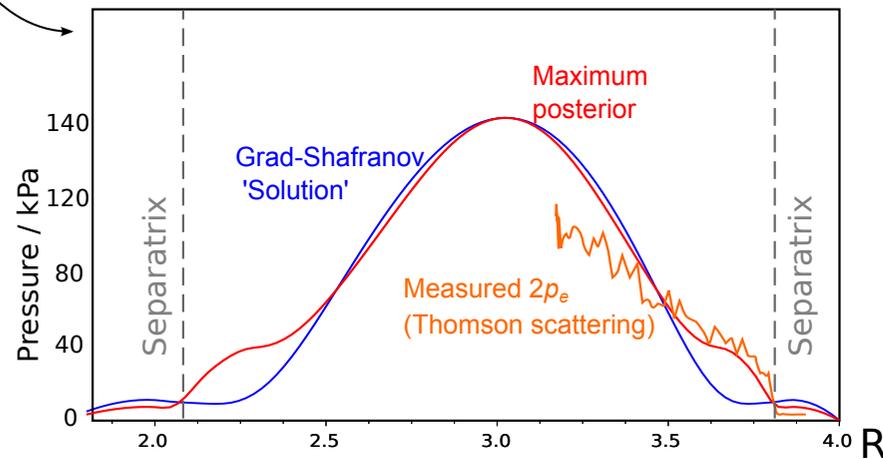
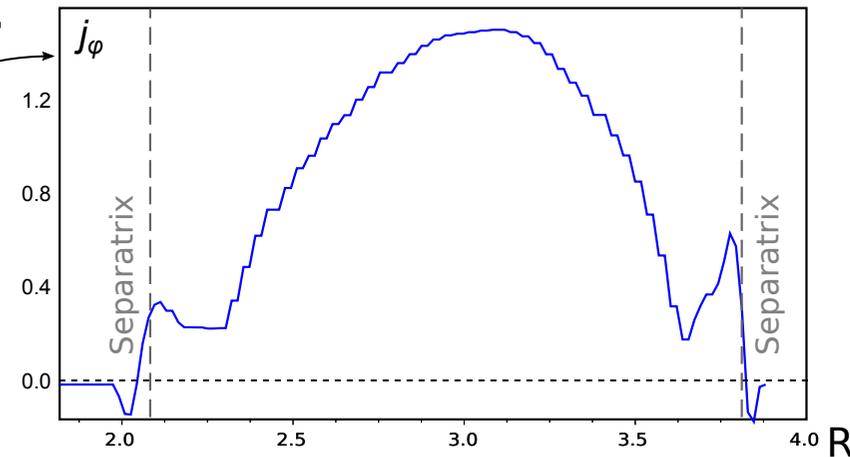
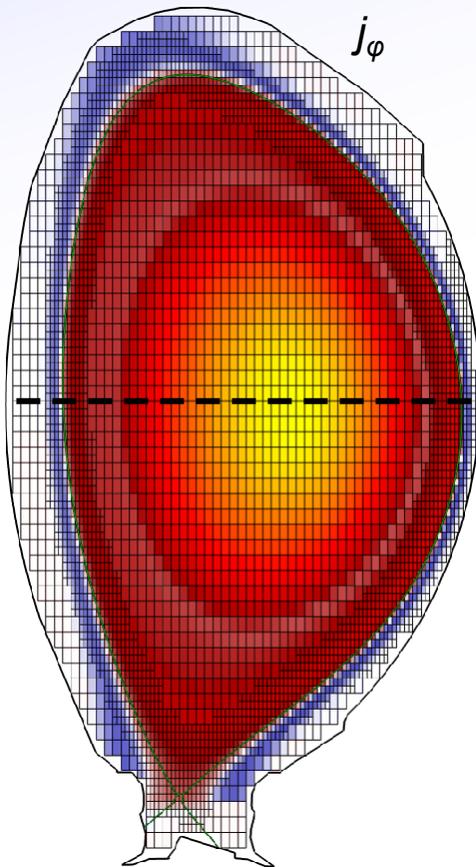
Bayesian Equilibrium

H-Mode plasmas: Very sharp changes in j and p require high-resolution current beam grid:

Too many parameters to explore the posterior (Monte-Carlo algorithm). Needs:

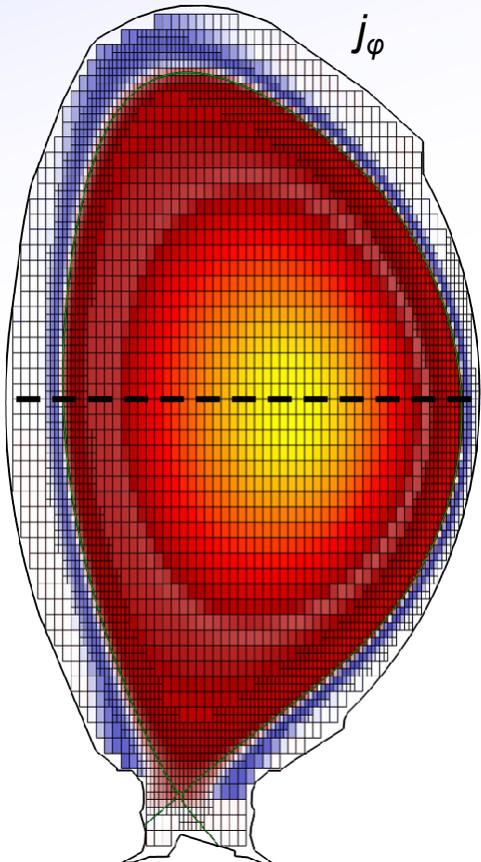
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but, maximum posterior can be calculated:



Bayesian Equilibrium

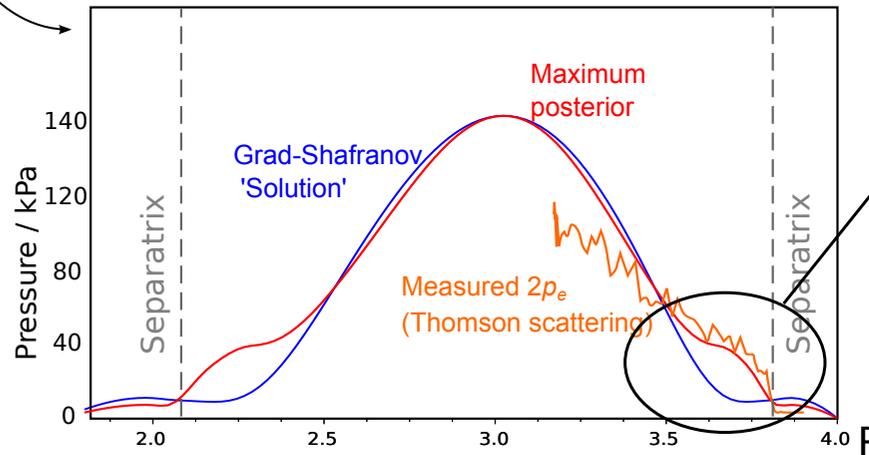
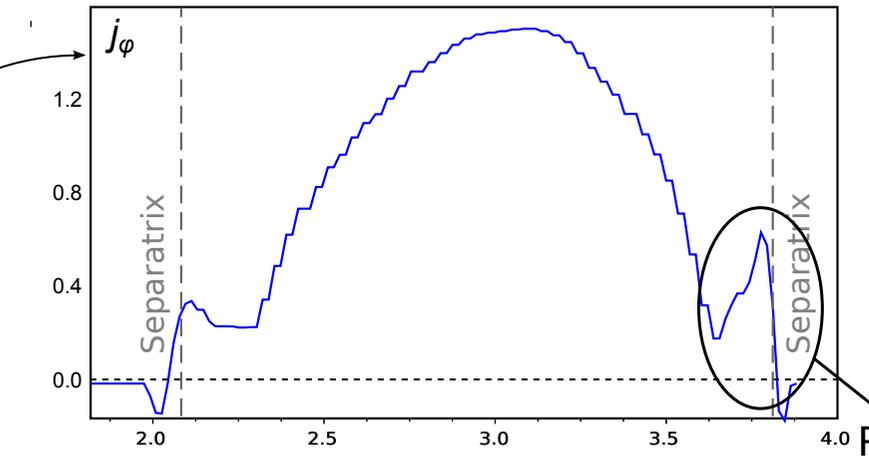
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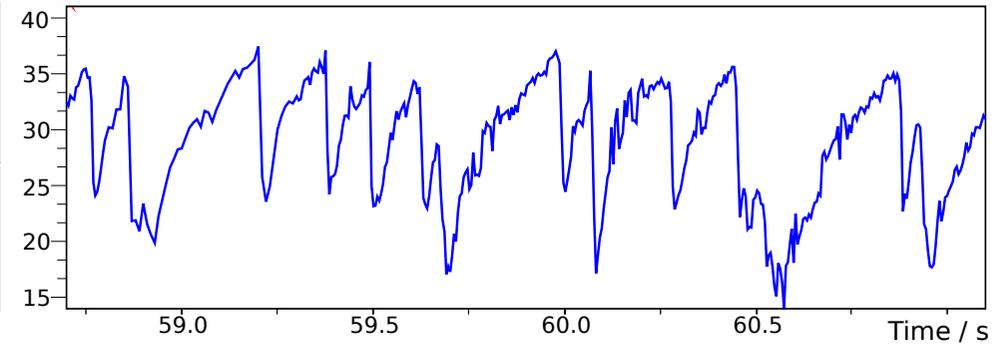
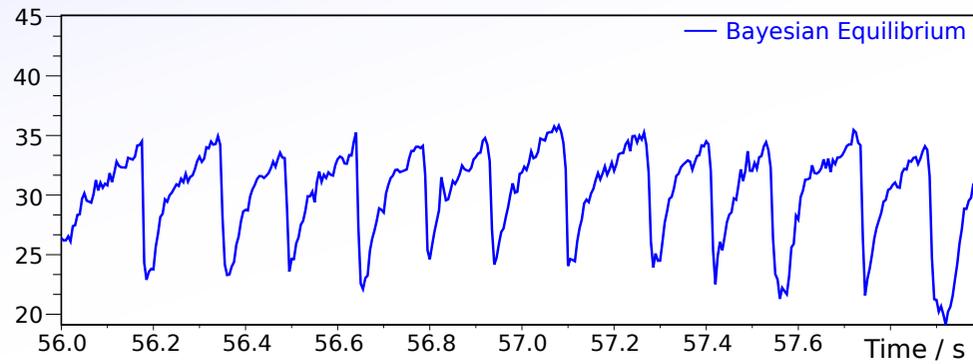
but, maximum posterior can be calculated:



Pedestal current and pressure resolved from external magnetic measurements alone.

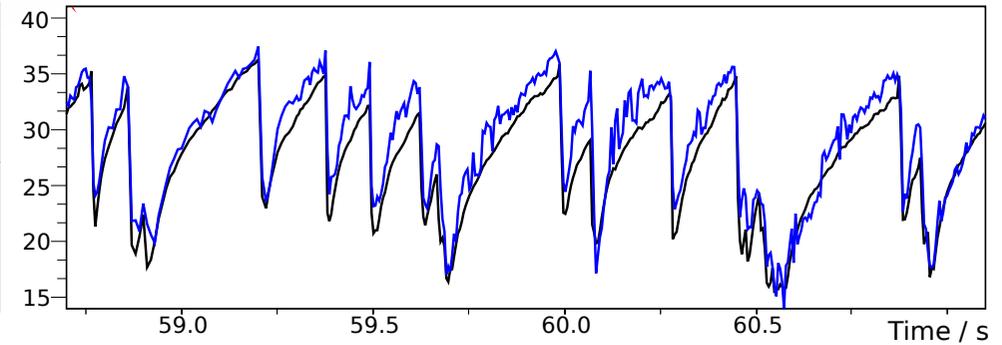
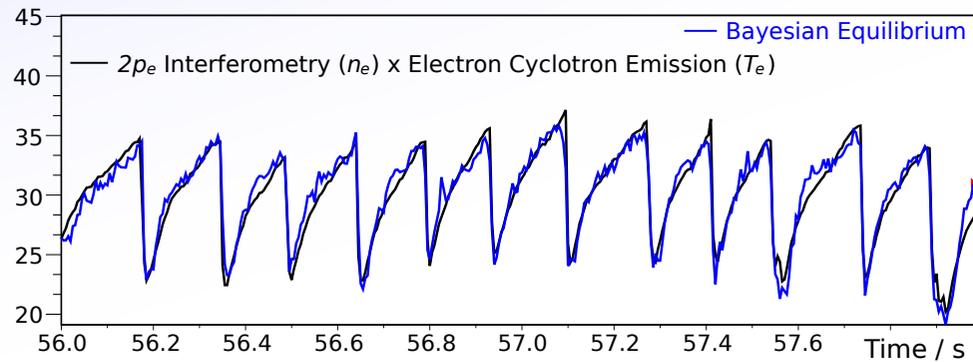
Pedestal Pressure

Flexible p , f profiles show that pedestal pressure can be very accurately measured with magnetic coils.



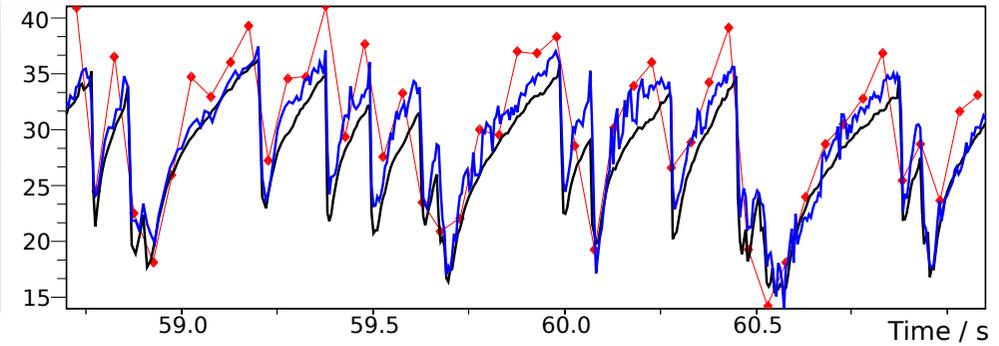
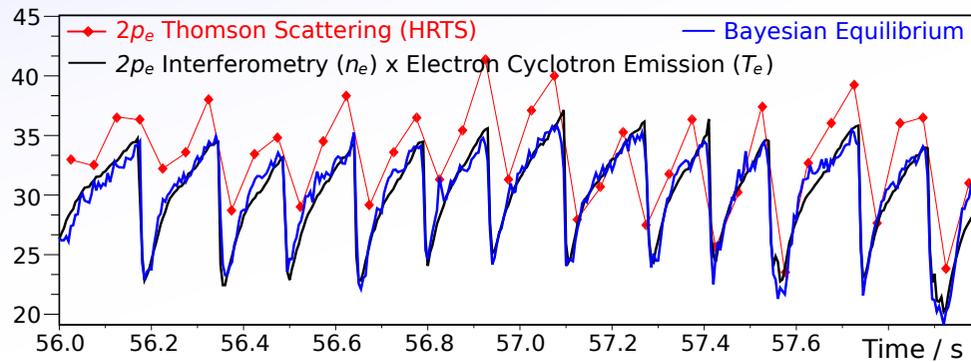
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Pedestal Pressure

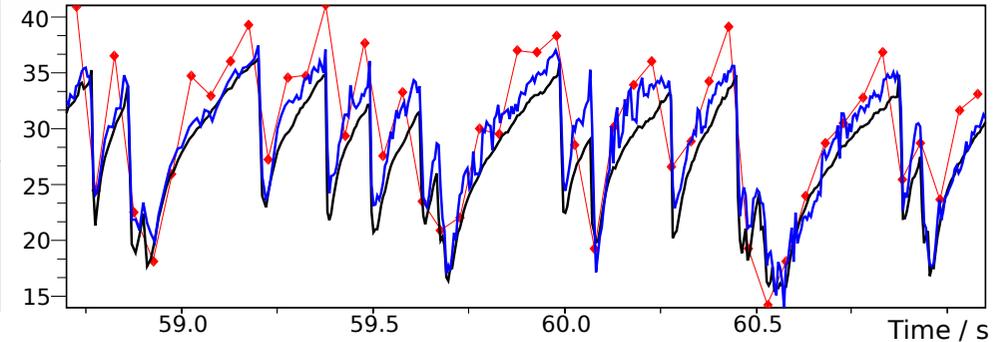
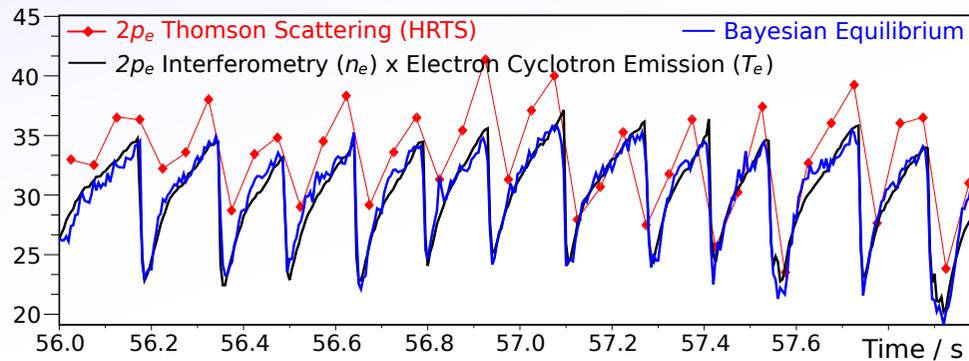
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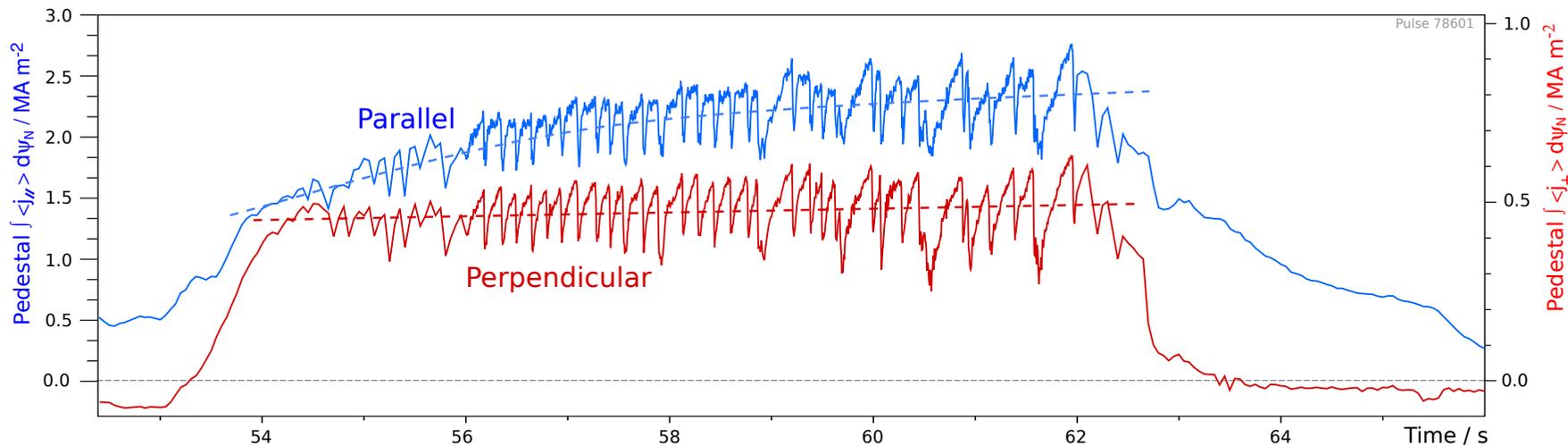
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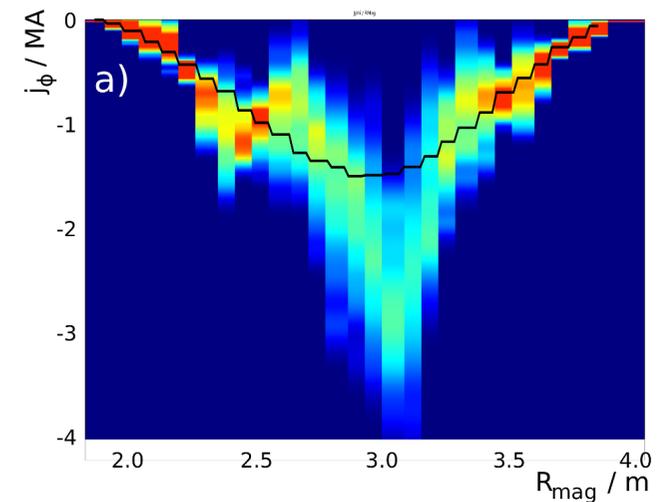
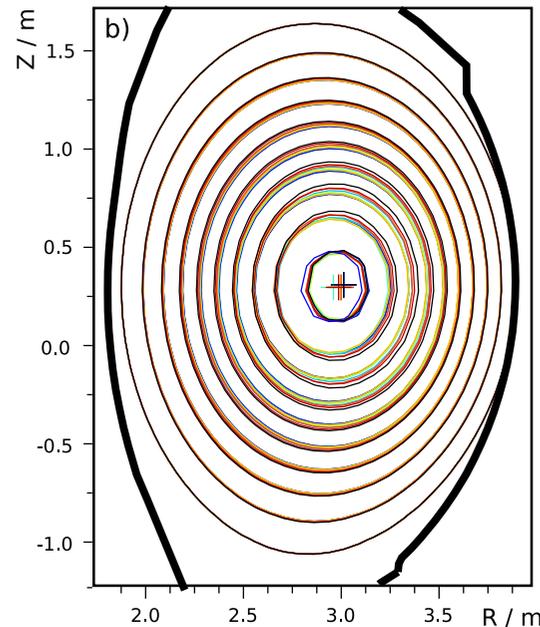
Pedestal parallel and perpendicular currents can be separated:



- Very good information on edge current, even from magnetics alone!

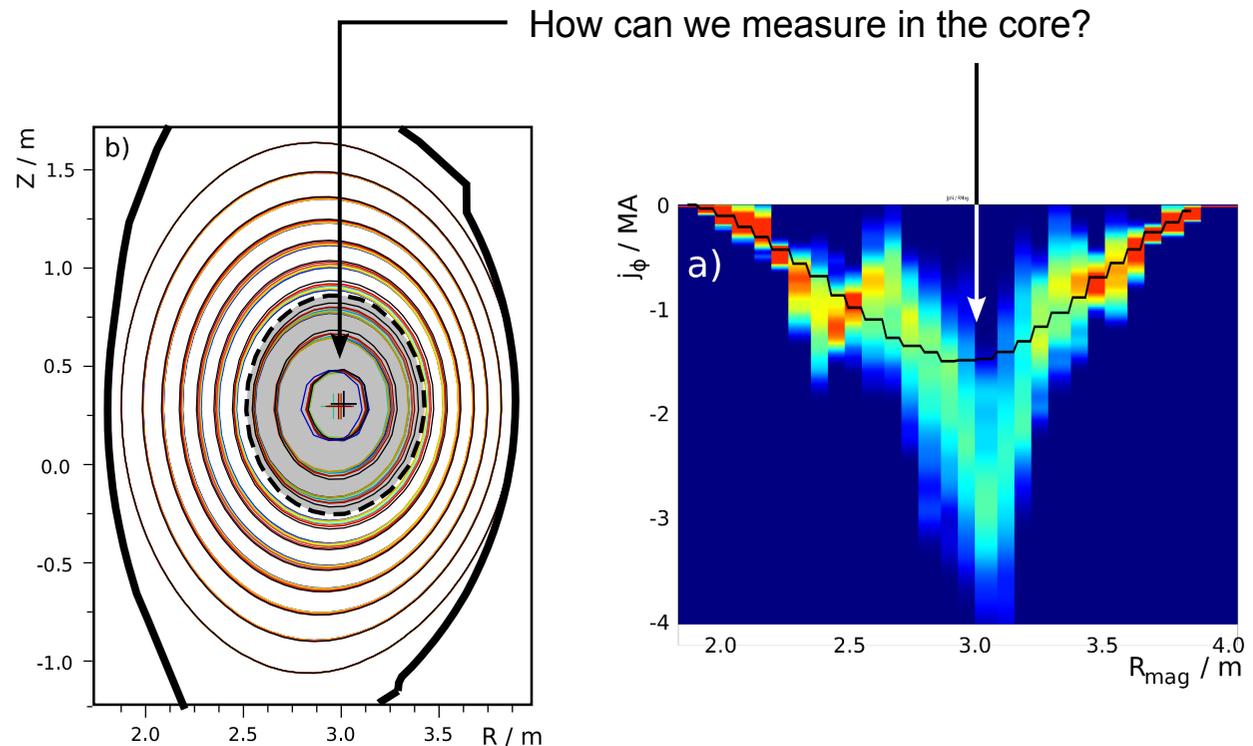
Outline

- Introduction
 - Flux surfaces and current profiles
 - Magnetic equilibrium
 - Bayesian analysis
- Bayesian equilibrium
 - Current-tomography
 - Current tomography + Grad-Shafranov
 - L-Mode reconstructions
 - H-Mode results
- Internal measurements
 - Motional Stark effect.
 - Coherence Imaging
 - Imaging MSE
 - Direct j_ϕ imaging.
- Integrated Data Analysis
 - Current diffusion
 - IMSE results in comparison
 - Sawtooth models
- Rigorous determination of uncertainty
- Too computationally intensive for H-mode
- Need internal measurements!



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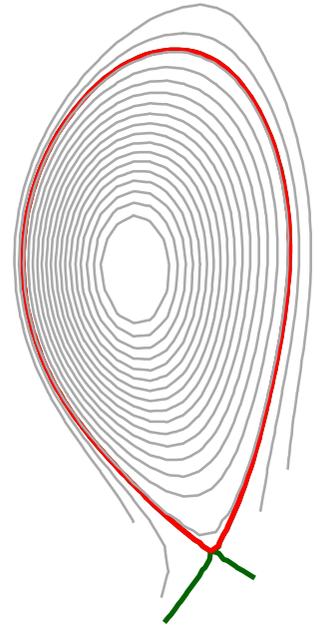
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Internal Measurements

How can we measure deep inside the plasma?

Magnetic Surfaces
Plasma Edge

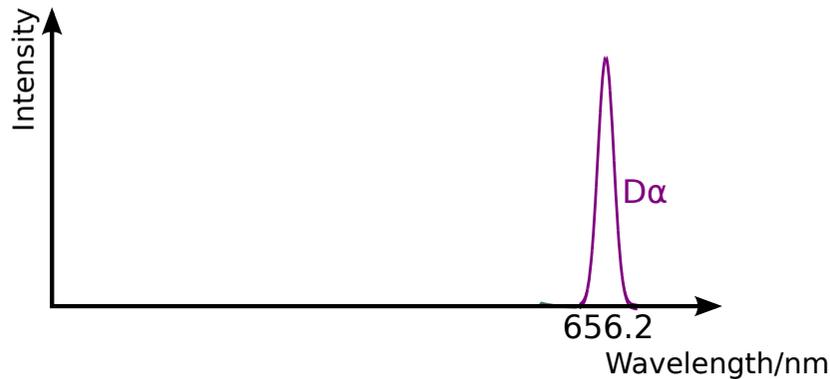
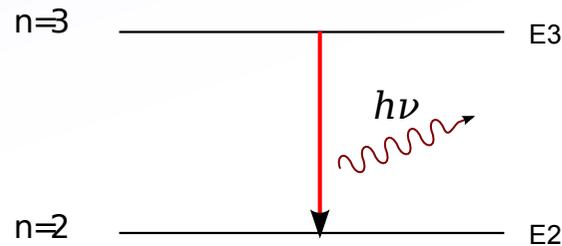


Internal Measurements

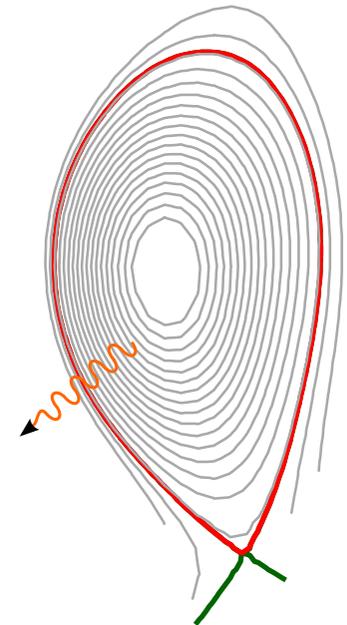
How can we measure deep inside the plasma?

Spectroscopy - observe the light emitted by atoms in the plasma:

e.g. Hydrogen Balmer- α line:



Magnetic Surfaces
Plasma Edge



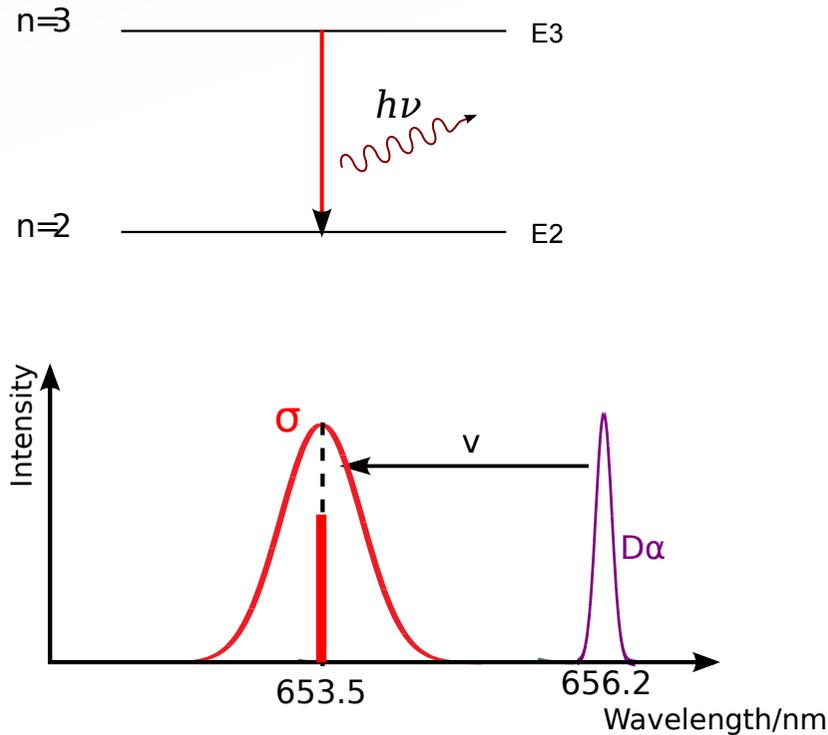
Inject high-energy neutral hydrogen into core of plasma (for heating/fueling)
Excitation by ion/electron impact excites the higher energy levels.
Spontaneous decay emits photon that can be measured by a spectrometer.

Internal Measurements

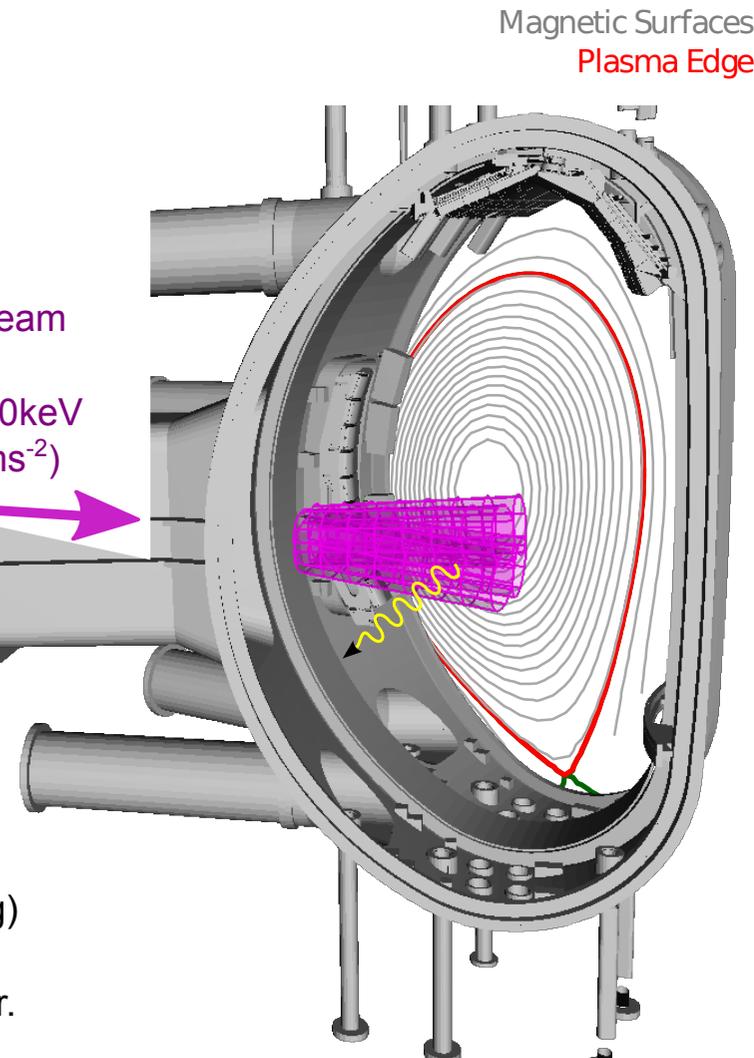
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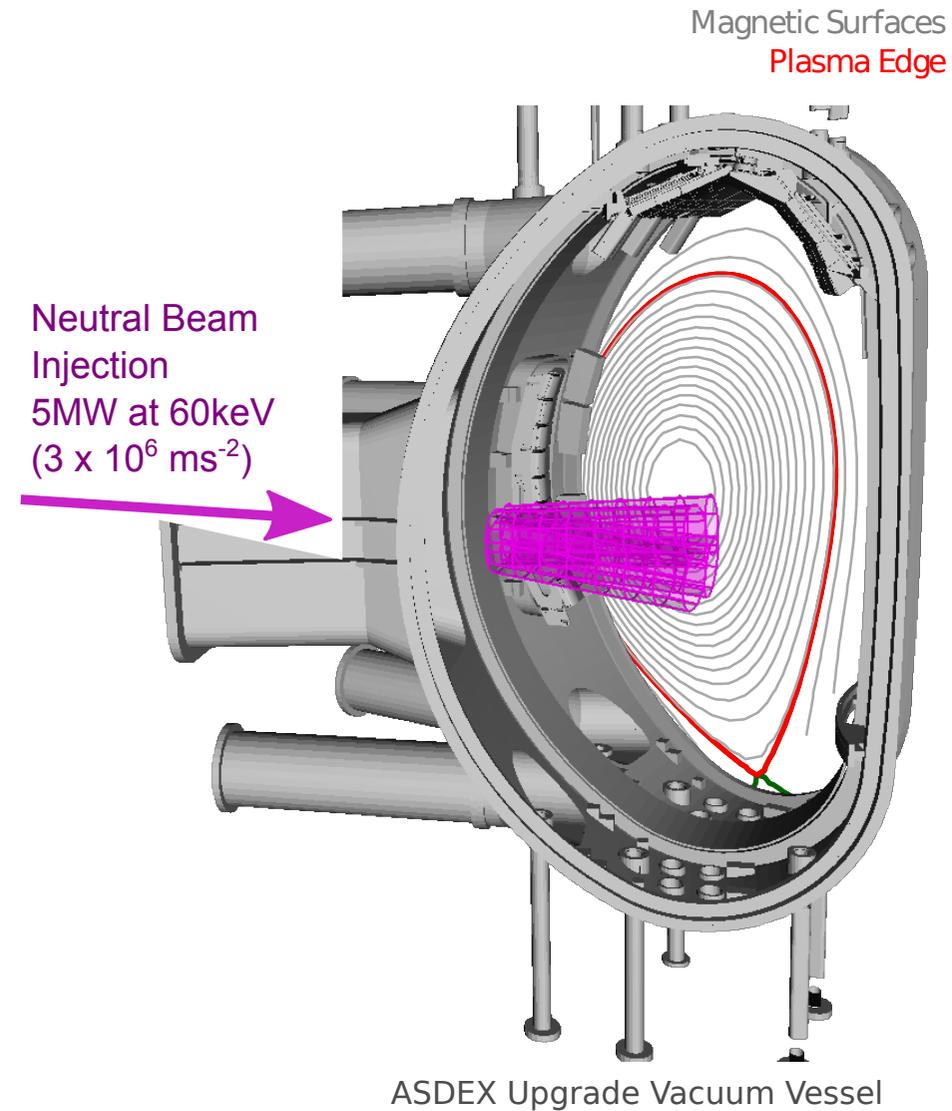
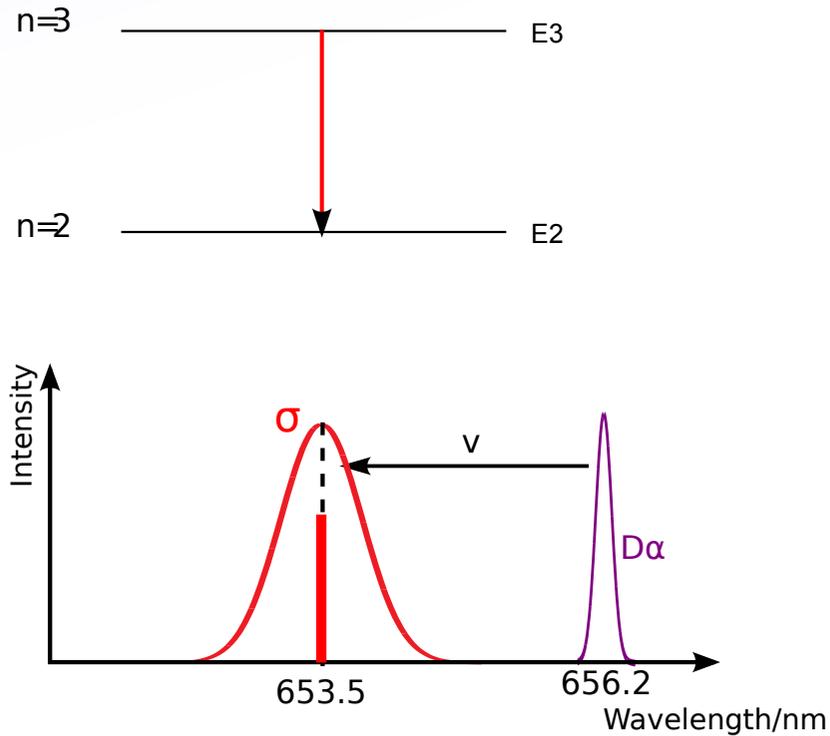
Neutral Beam
Injection
5MW at 60keV
($3 \times 10^6 \text{ ms}^{-2}$)



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ASDEX Upgrade Vacuum Vessel

Motional Stark Effect Polarimetry



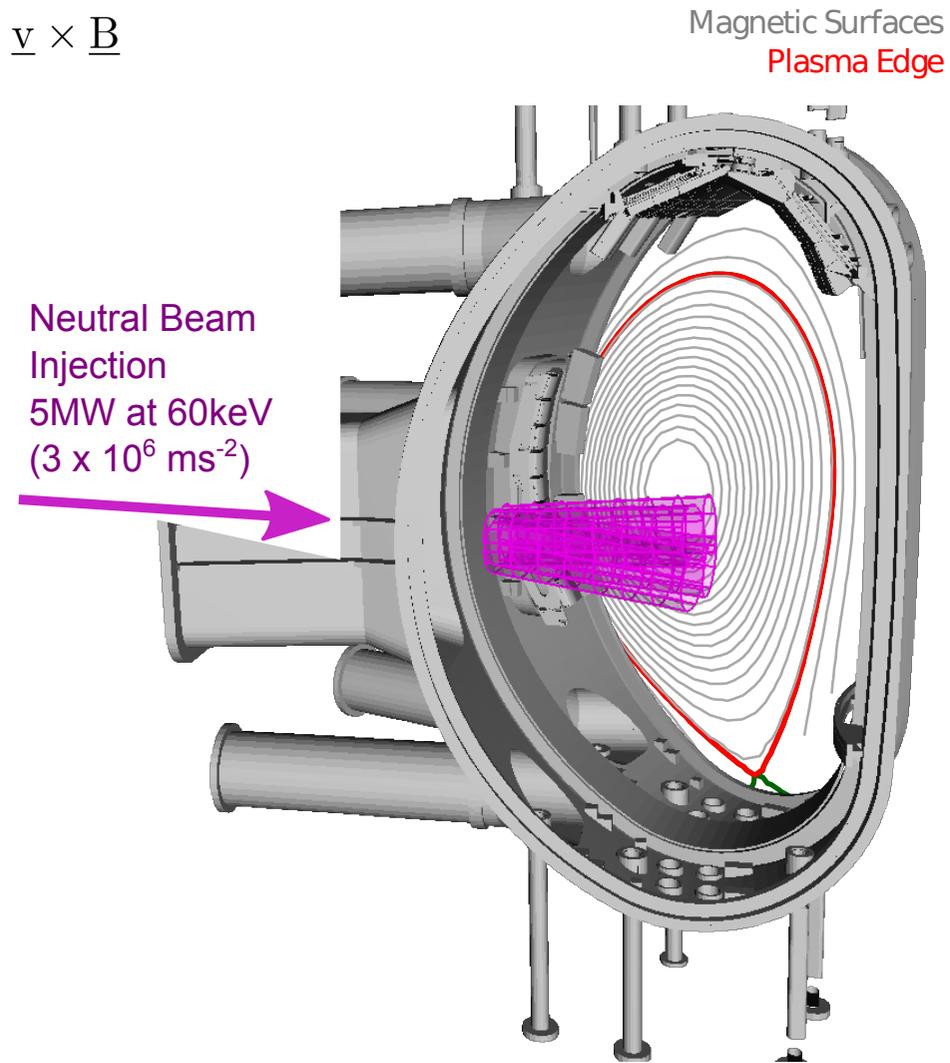
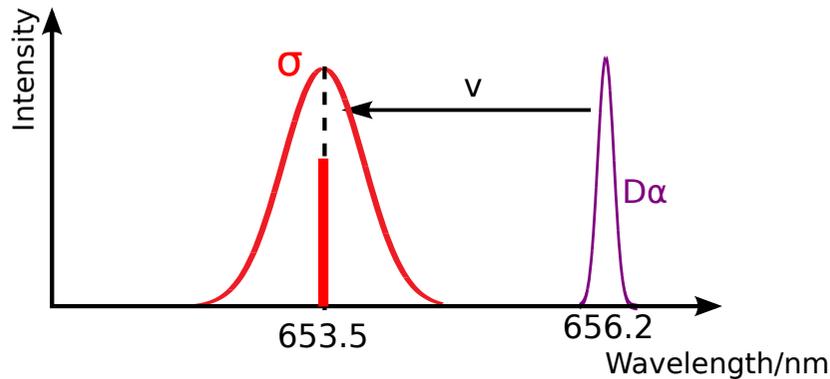
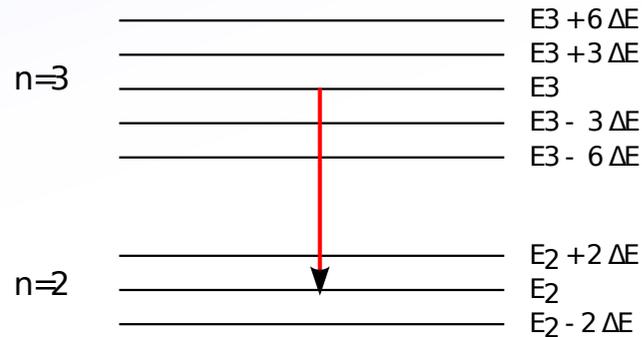
Motional Stark Effect Polarimetry

The atomic energy levels are modified by the local magnetic/electric fields:

- Zeeman splitting (magnetic field)
- Stark splitting (electric field):

Stark splitting by Lorenz-transformed magnetic field:

$$\underline{E} = \underline{v} \times \underline{B}$$



ASDEX Upgrade Vacuum Vessel

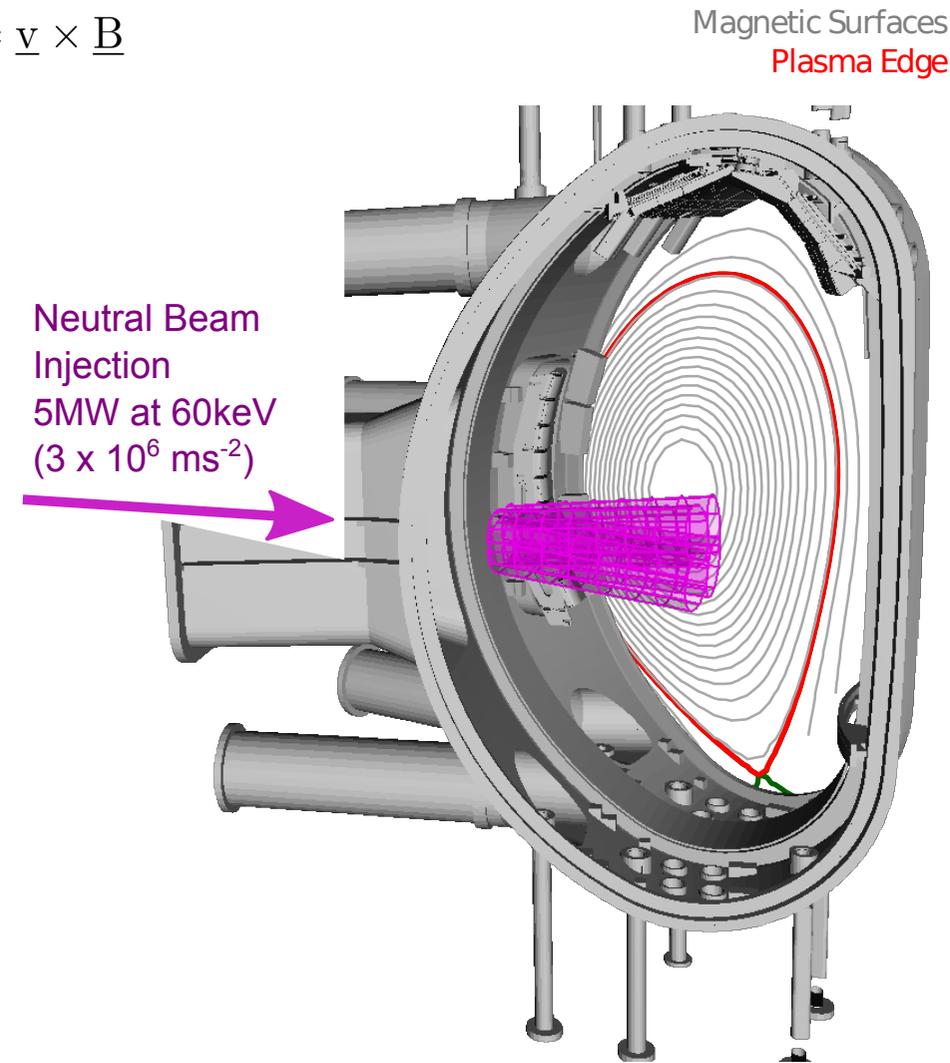
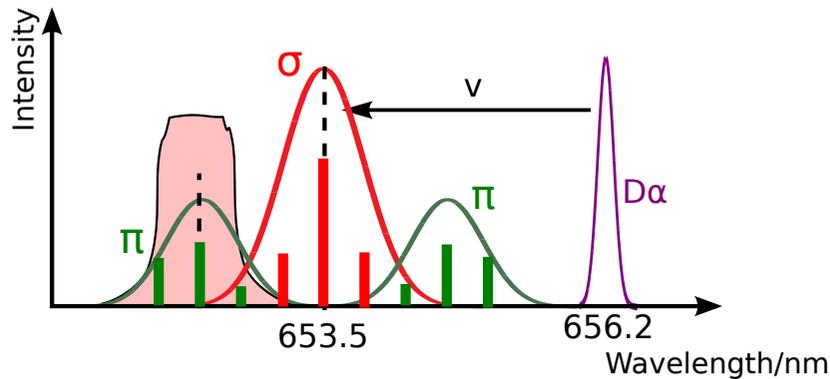
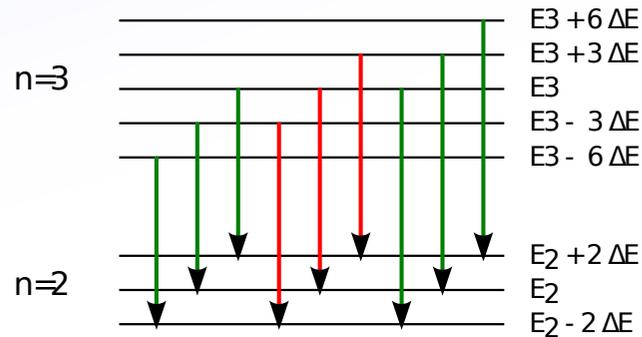
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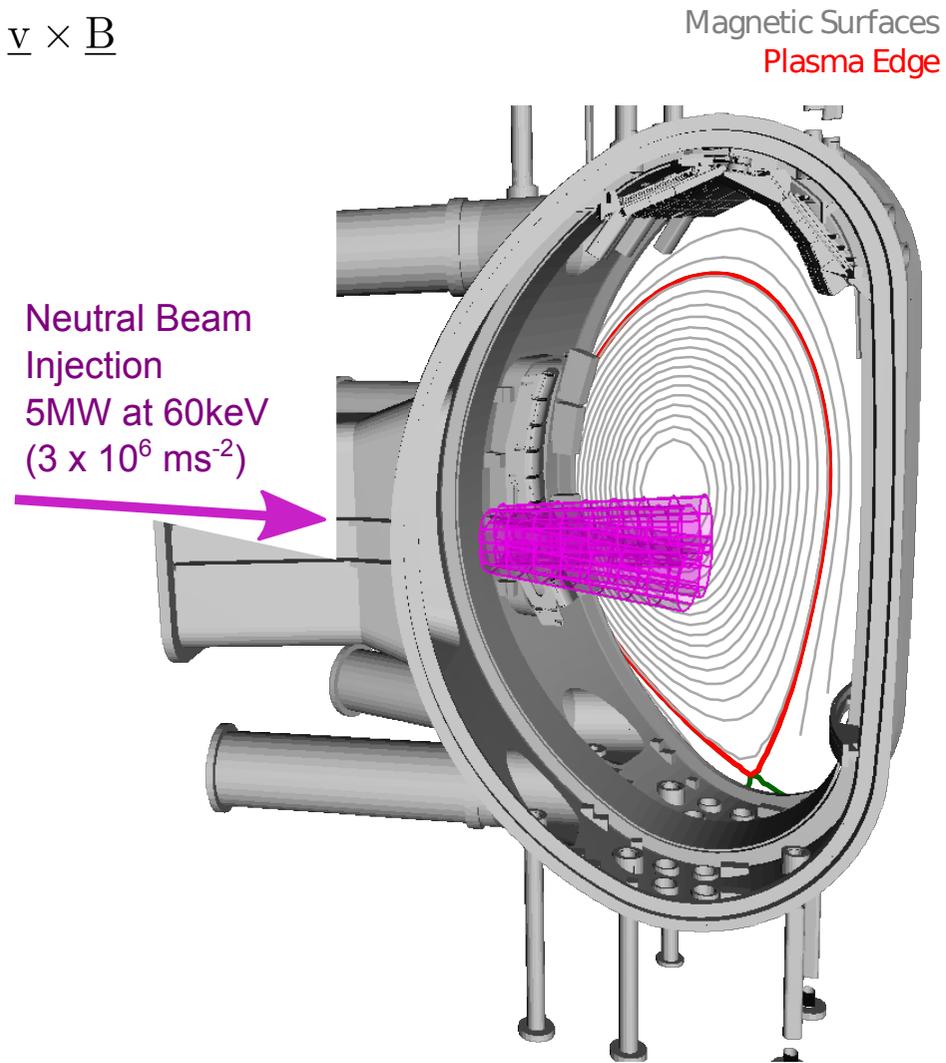
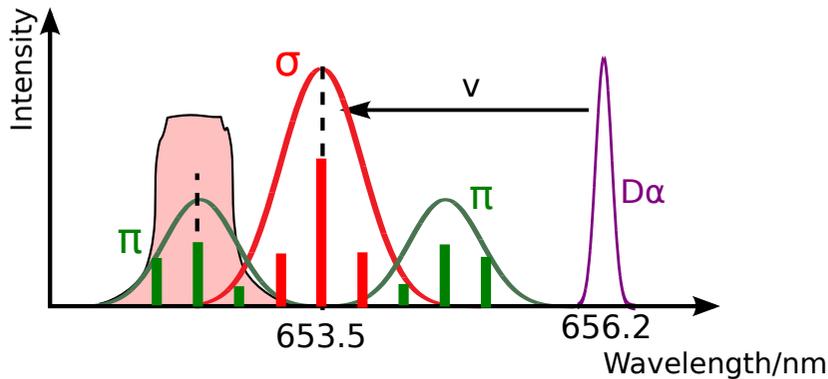
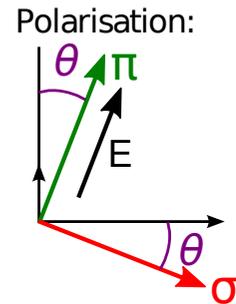
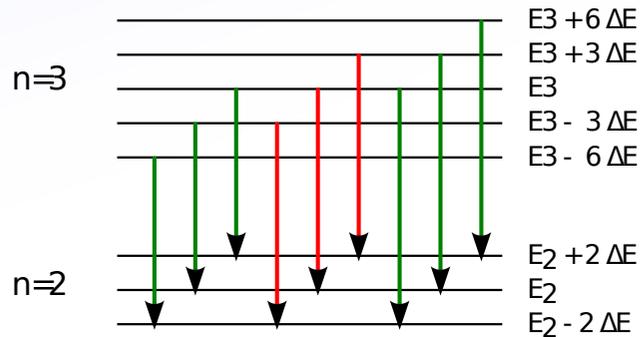
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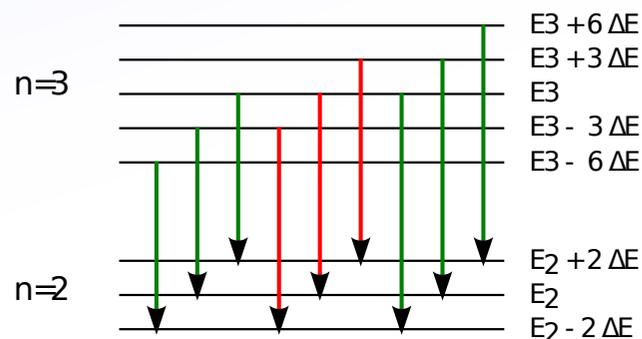
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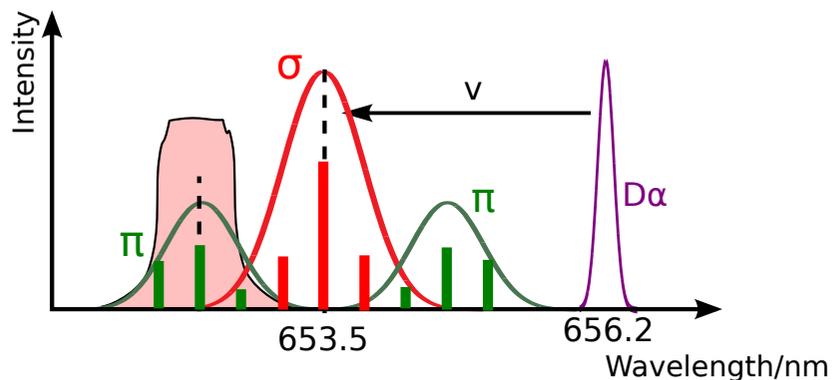
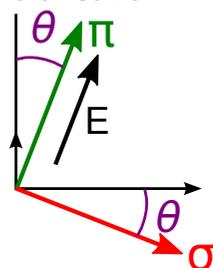
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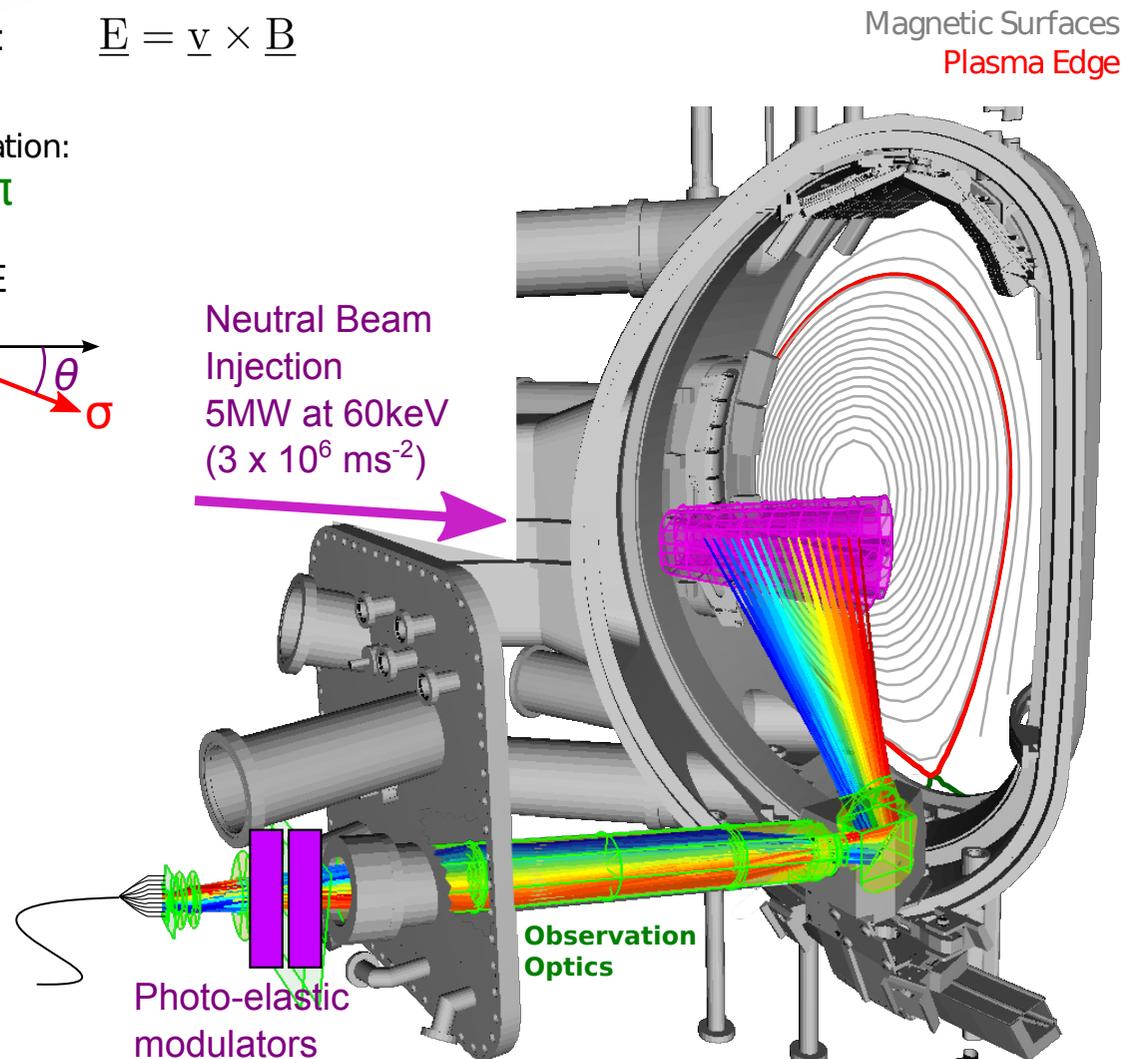
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Polarisation:



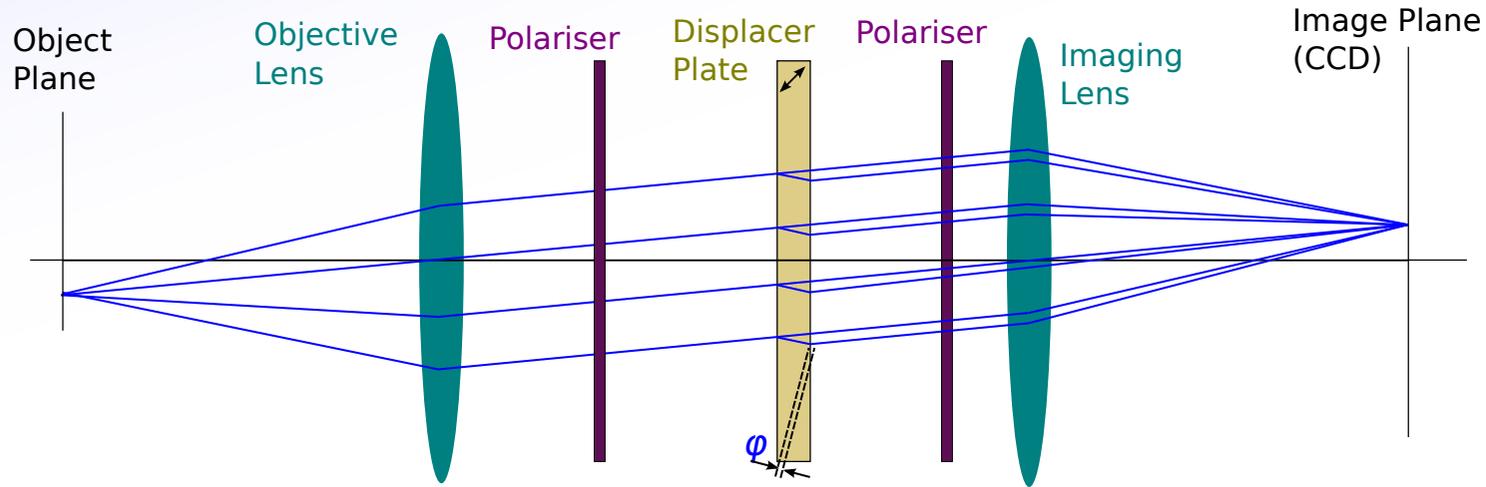
The MSE diagnostic:
Isolate a π component and measure the
polarisation --> magnetic field direction.



ASDEX Upgrade Vacuum Vessel

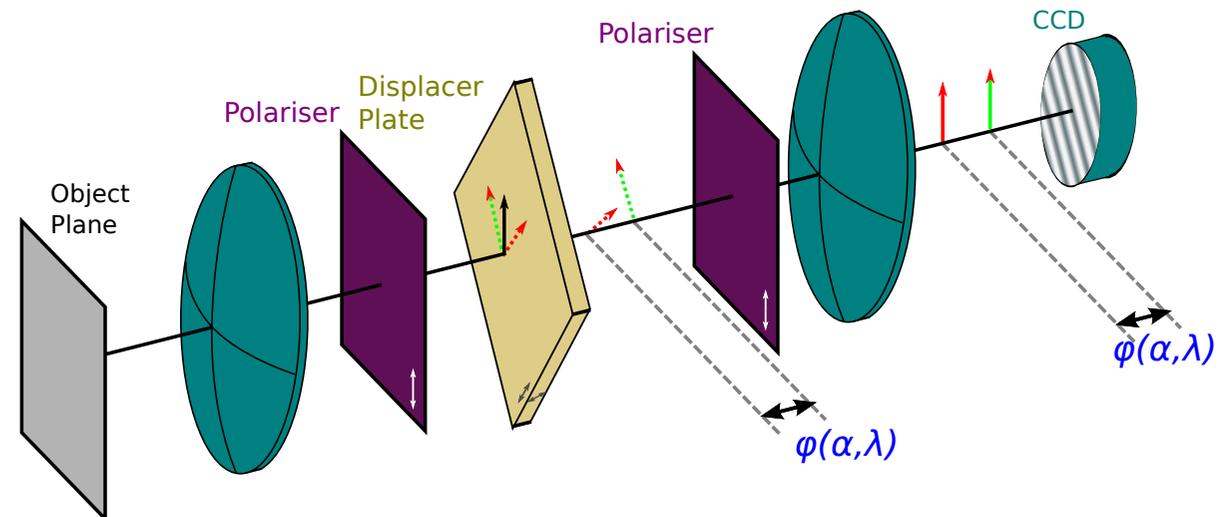
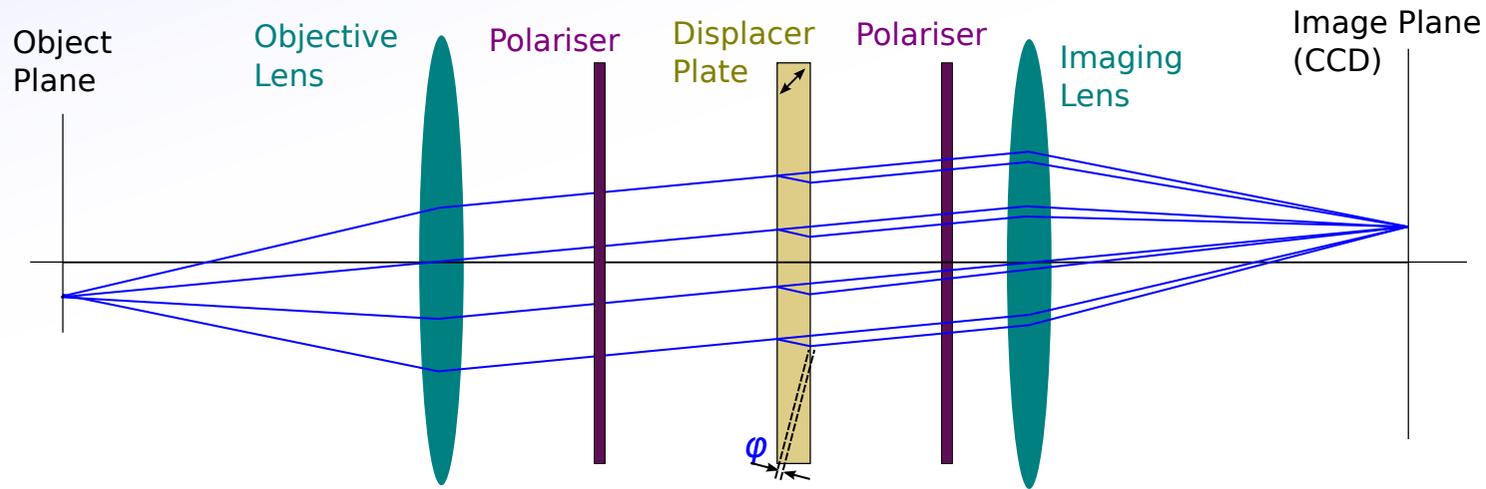
Coherence Imaging

Coherence Imaging: Spectroscopic technique, modulated in space and imaged with a CMOS camera.
Also used for imaging spectroscopic moments



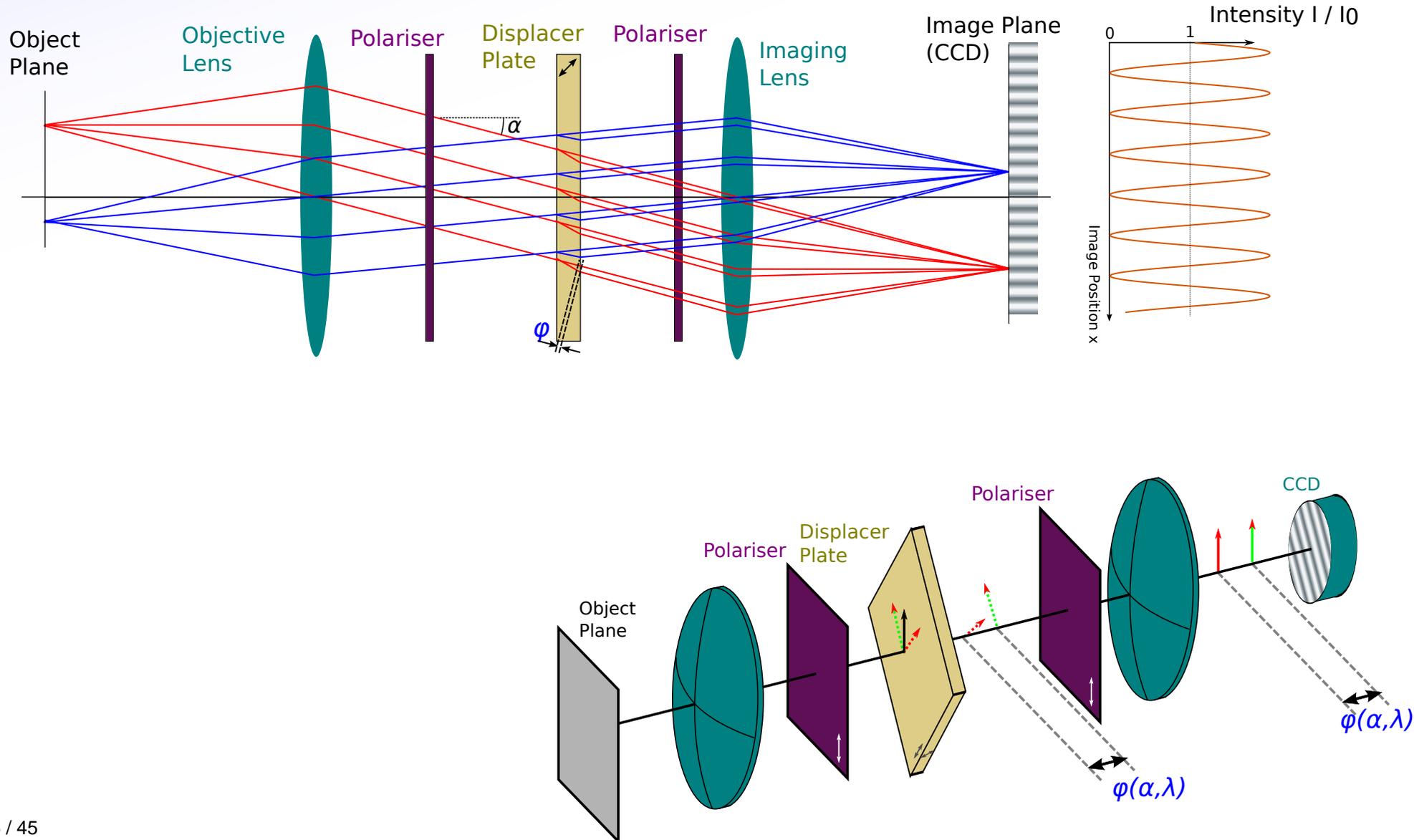
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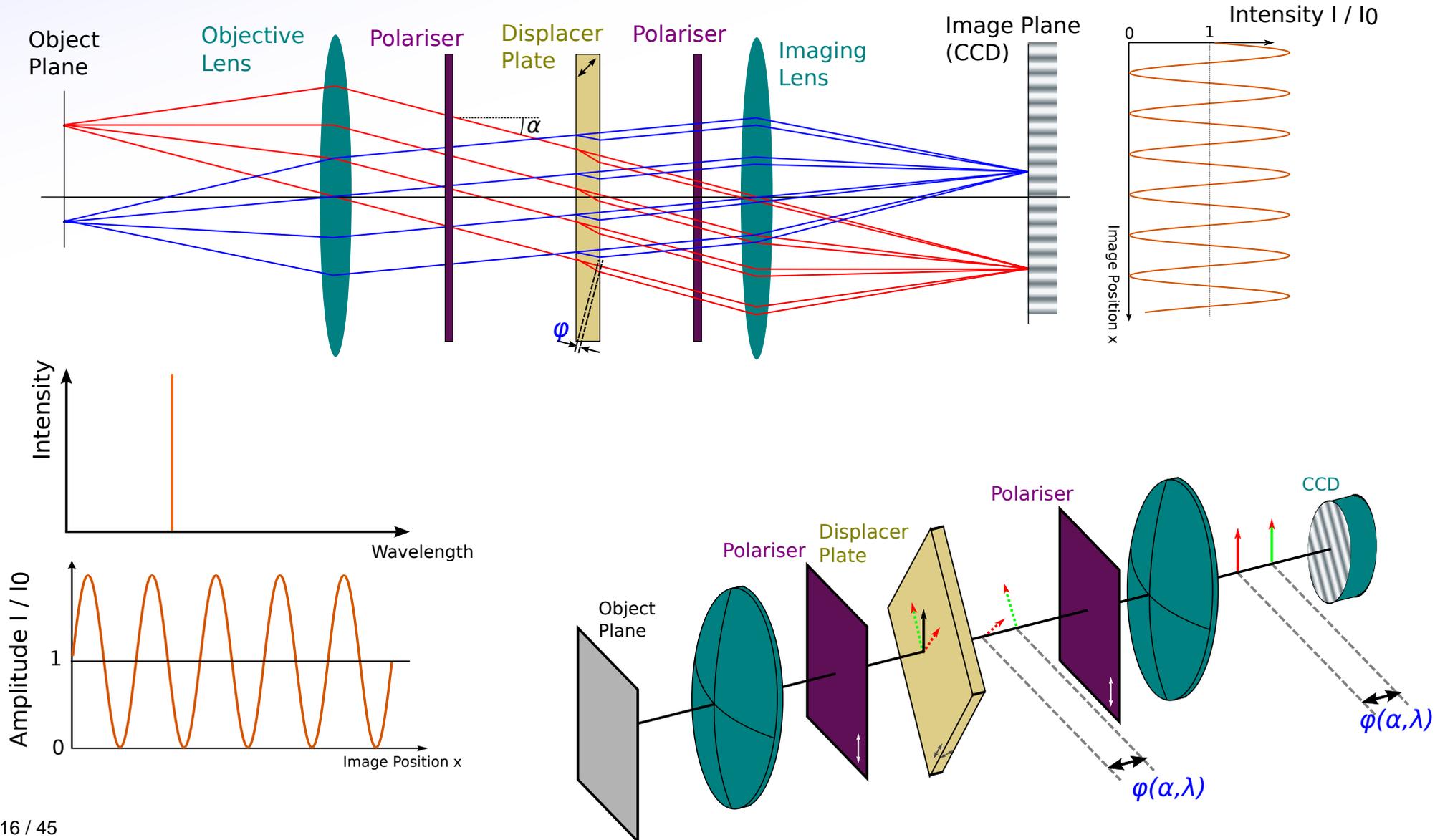
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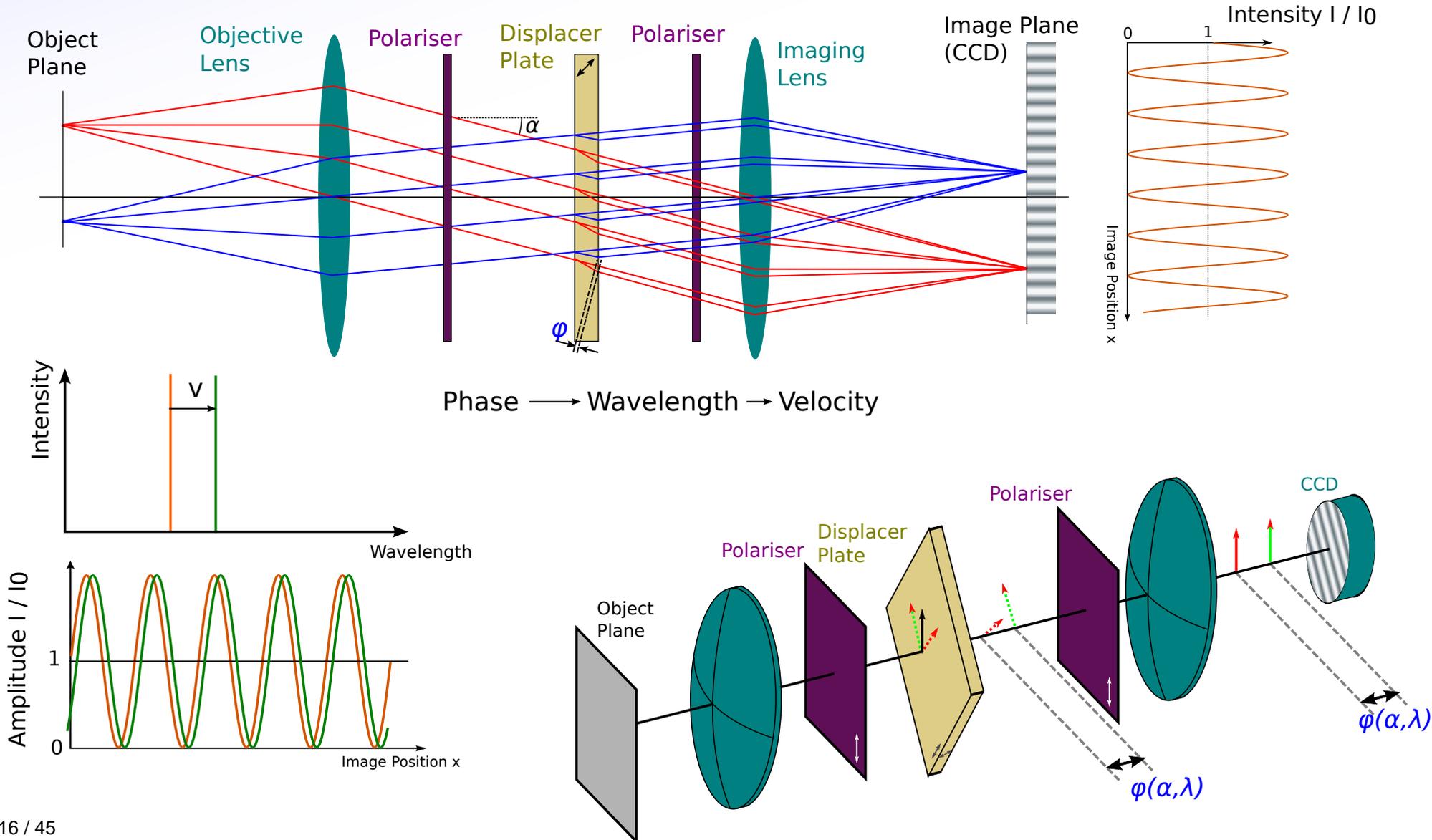
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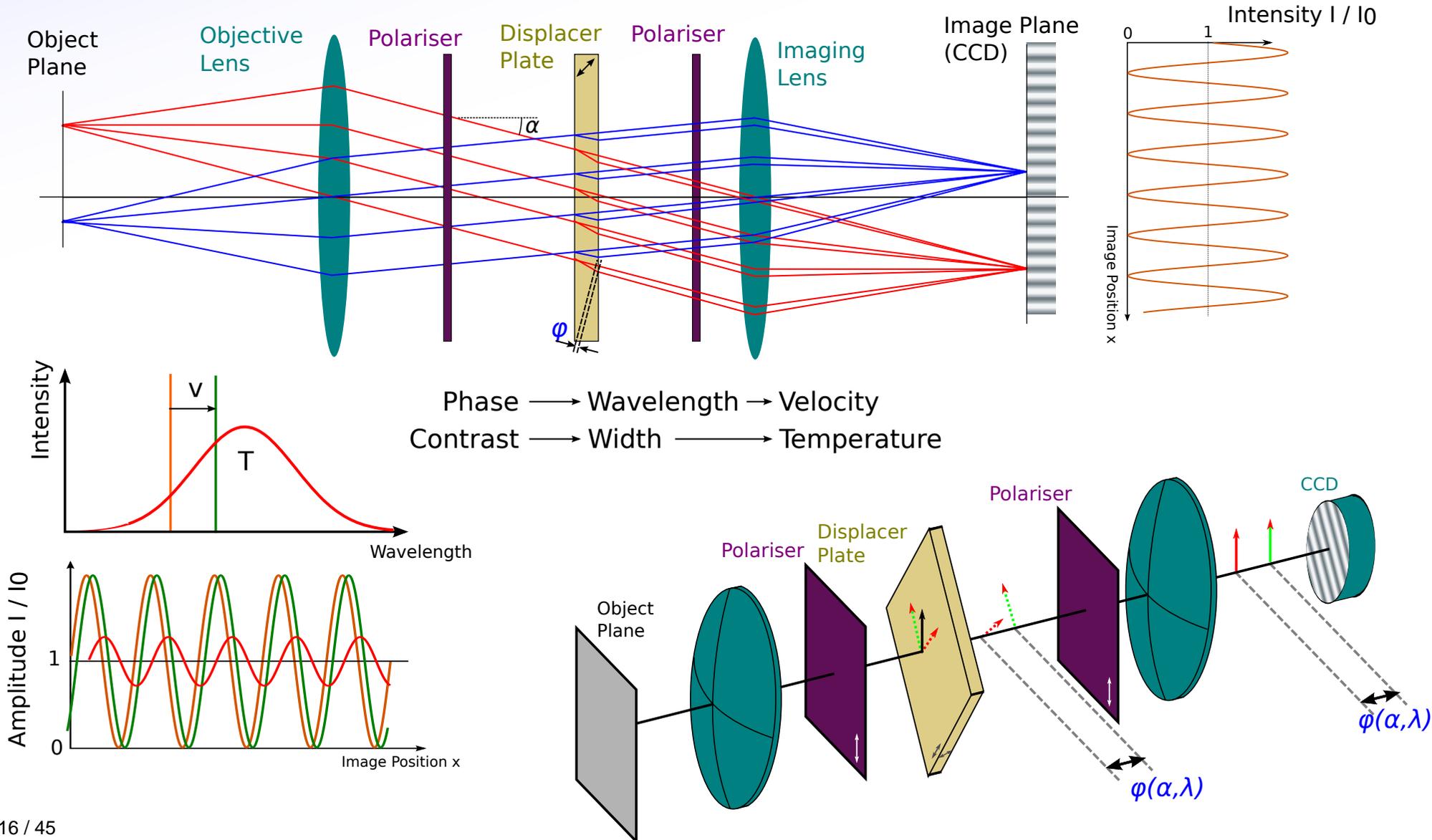
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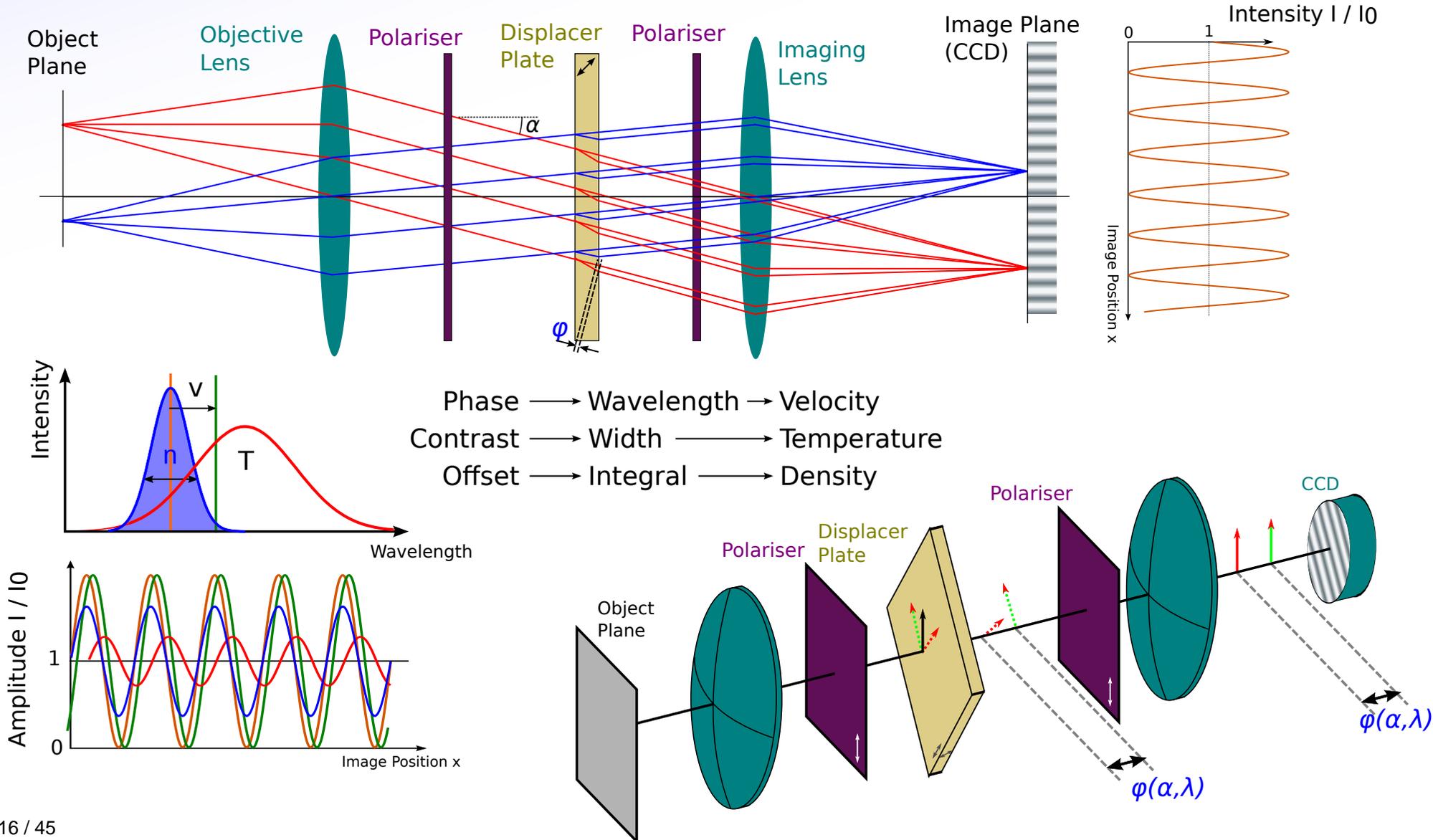
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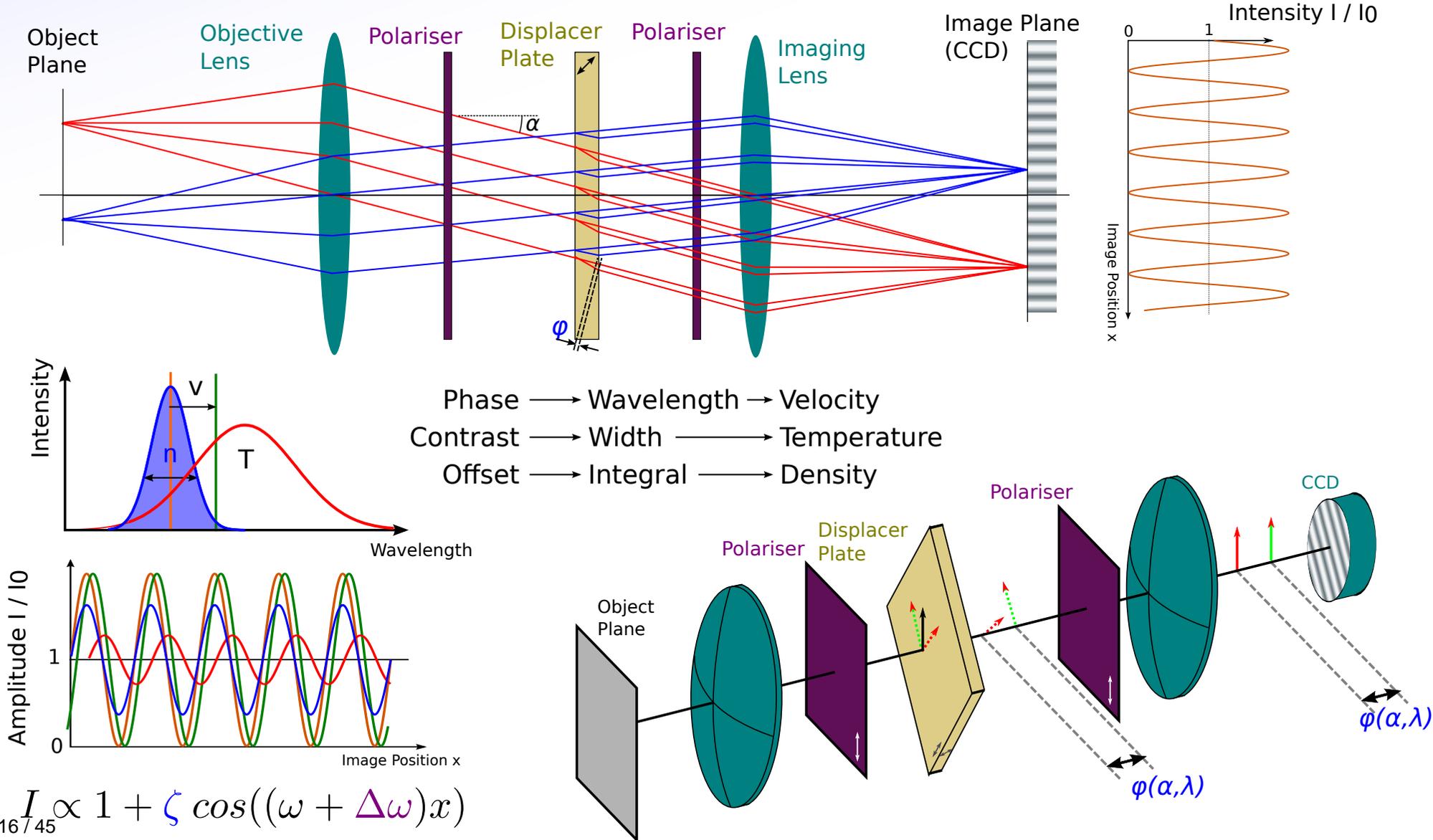
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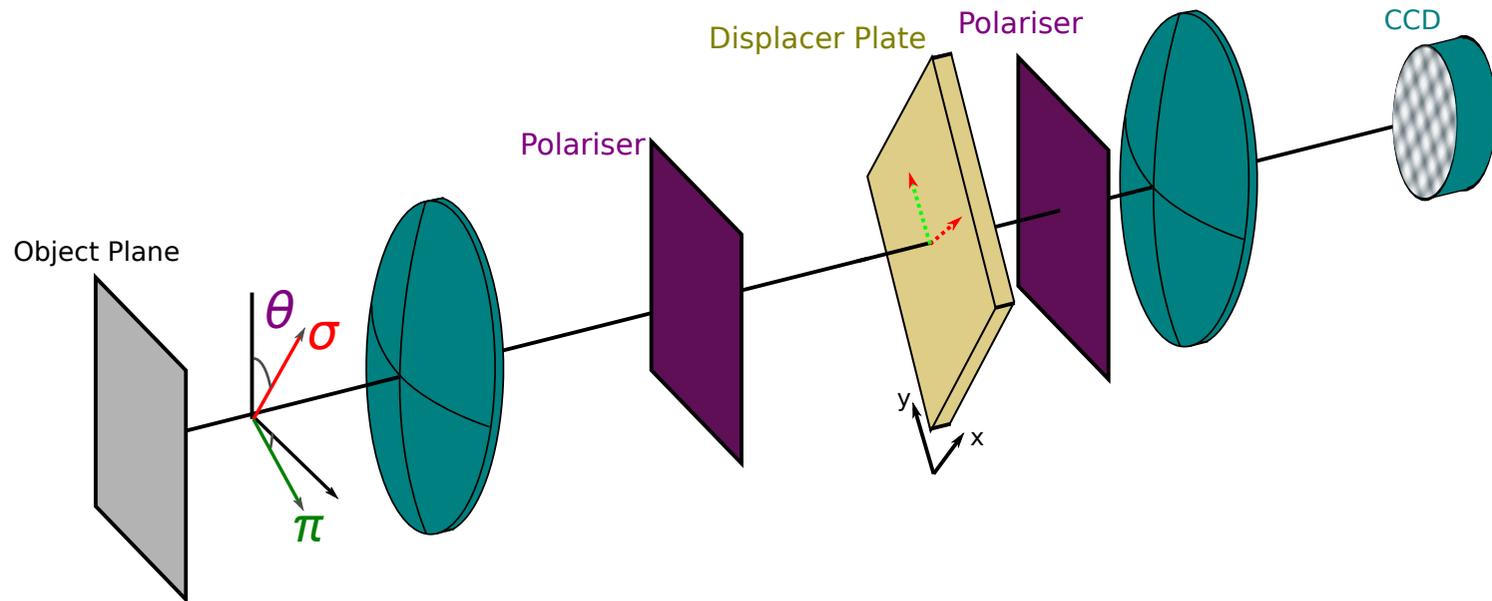
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Multiplet Polarisation Coherence Imaging

Removing the first polariser gives a
dependence on the initial polarisation:

$$I \propto 1 + \zeta \cos 2\theta \cos(x)$$

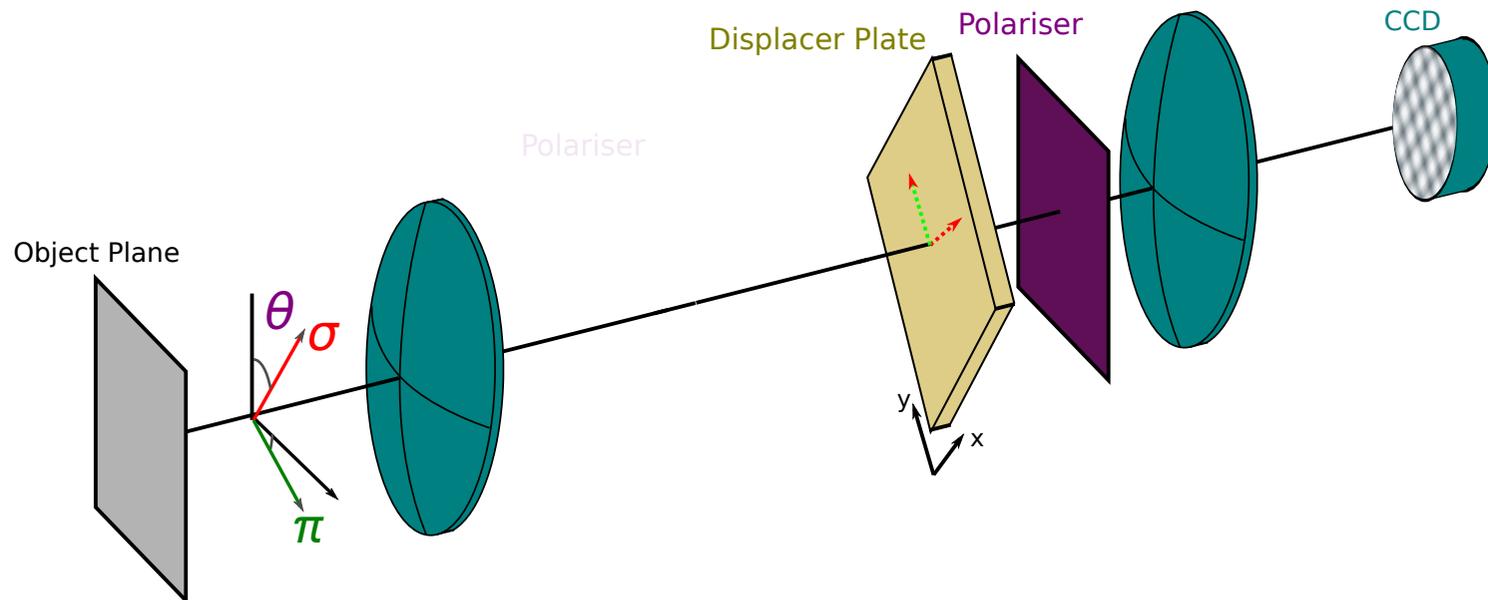
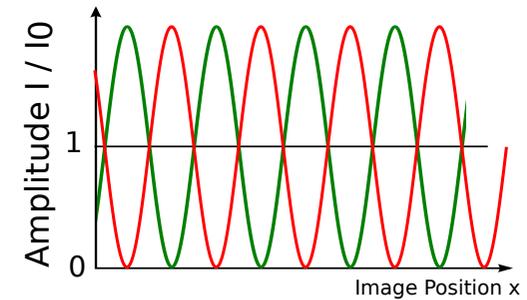
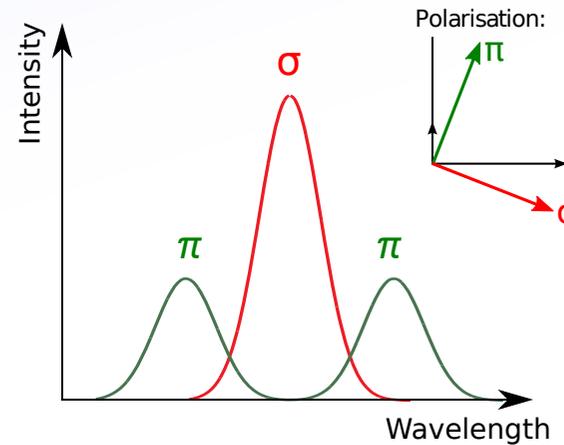


Multiplet Polarisation Coherence Imaging

Removing the first polariser gives a dependence on the initial polarisation:

$$I \propto 1 + \zeta \cos 2\theta \cos(x)$$

For the Stark/Zeeman spectrum, the π component is at 90° to σ , introducing a 180° phase shift, so they would cancel.



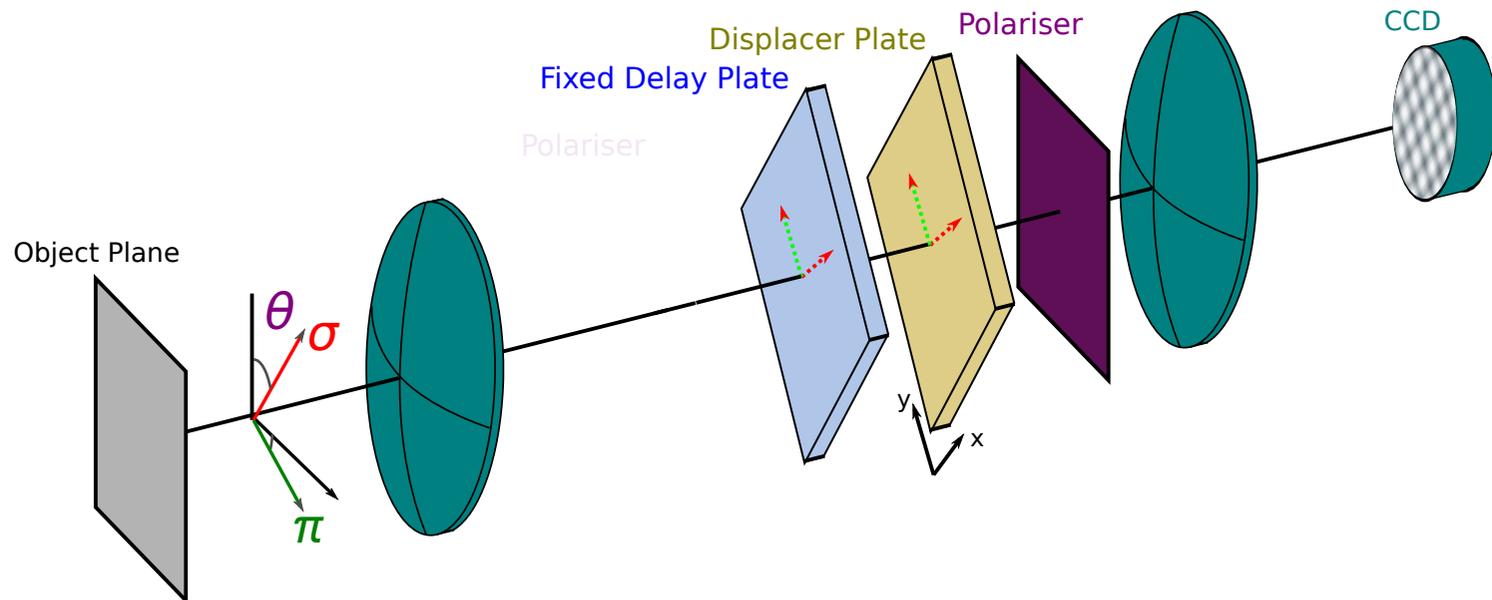
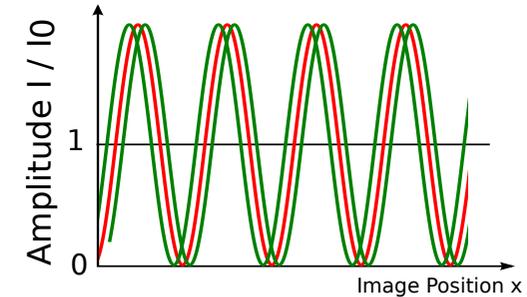
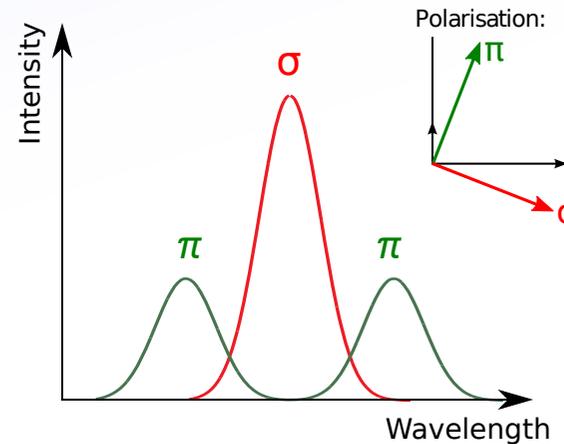
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At some specific plate thickness t , the phase of the π wings is 180° from σ . This cancels the 180° from the opposite polarisation, and the patterns add. We add a delay plate with the optimal τ_0 .



Multiplet Polarisation Coherence Imaging

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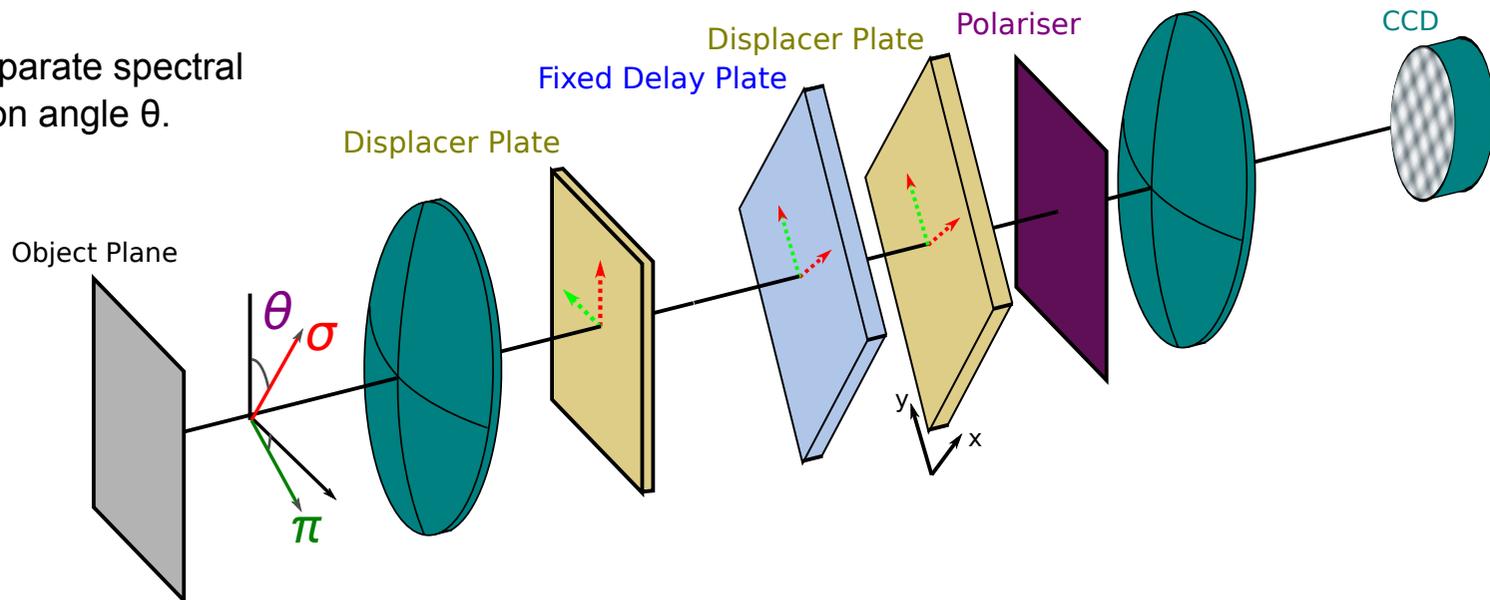
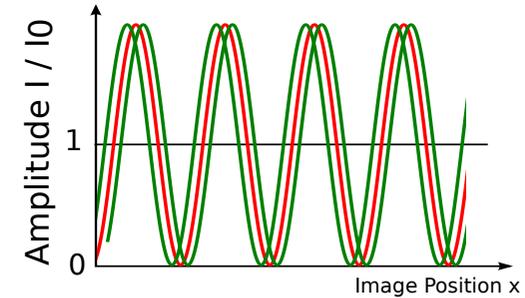
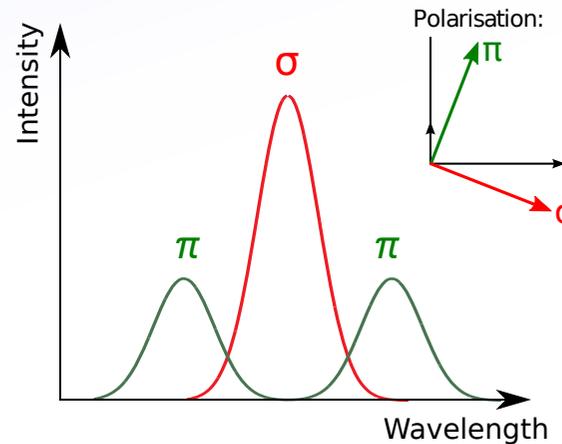
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At some specific plate thickness t , the phase of the π wings is 180° from σ . This cancels the 180° from the opposite polarisation, and the patterns add. We add a delay plate with the optimal τ_0 .

However, we now need to separate spectral contrast ζ from the polarisation angle θ .

add another displacer
at 45° . Combined effect
adds 2 extra terms:



$$I \propto 1 + \zeta \cos 2\theta \cos(x) + \zeta \sin 2\theta \cos(x - y) - \zeta \sin 2\theta \cos(x + y)$$

Imaging Motional Stark Effect results

Raw image (without neutral beam)

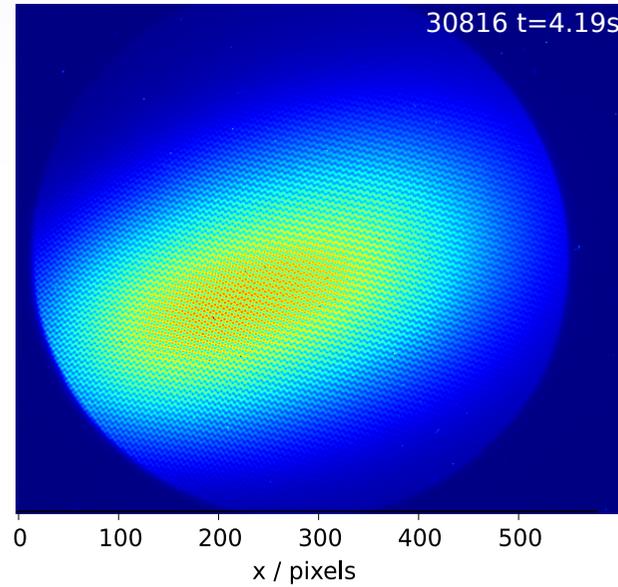


Imaging Motional Stark Effect results

Raw image (without neutral beam)



Raw image (with neutral beam)

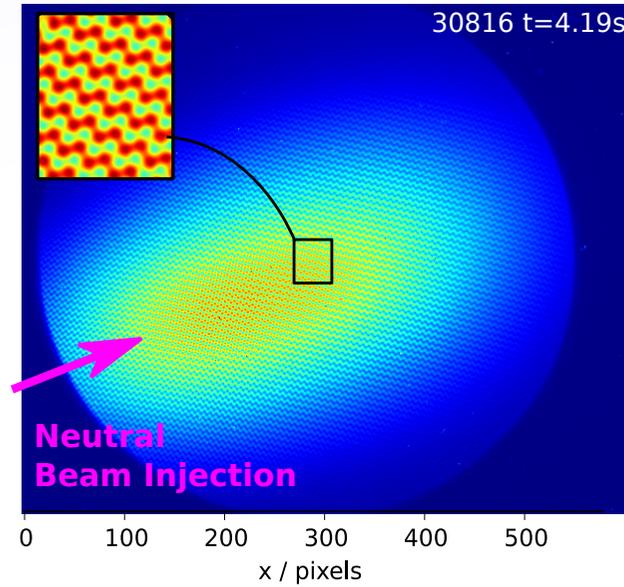


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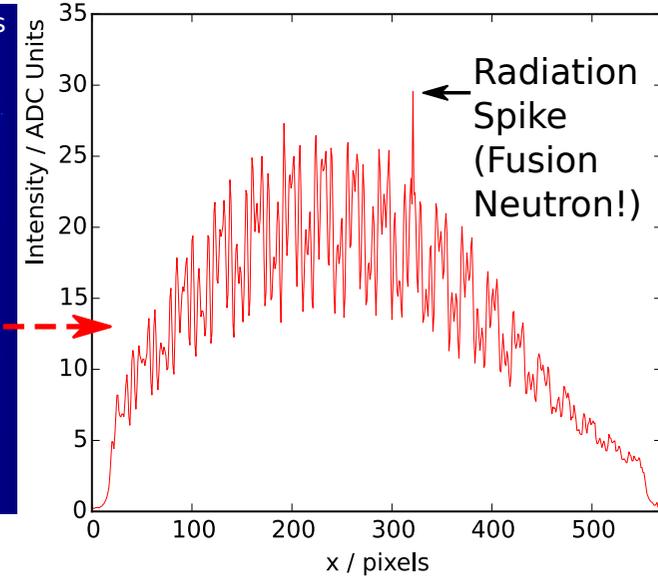
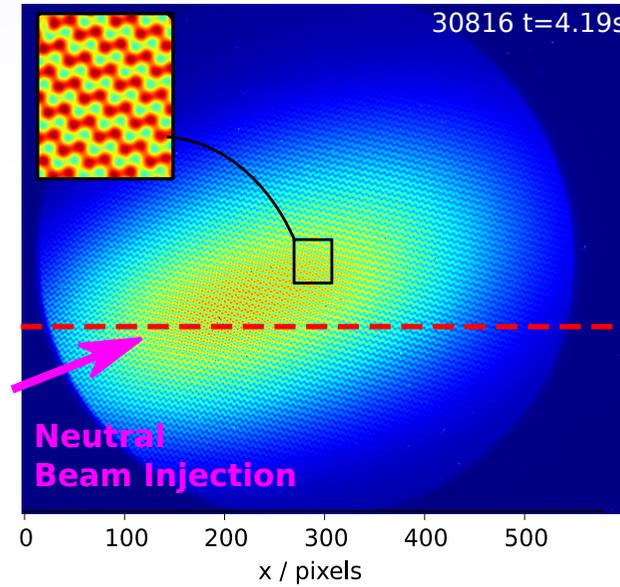


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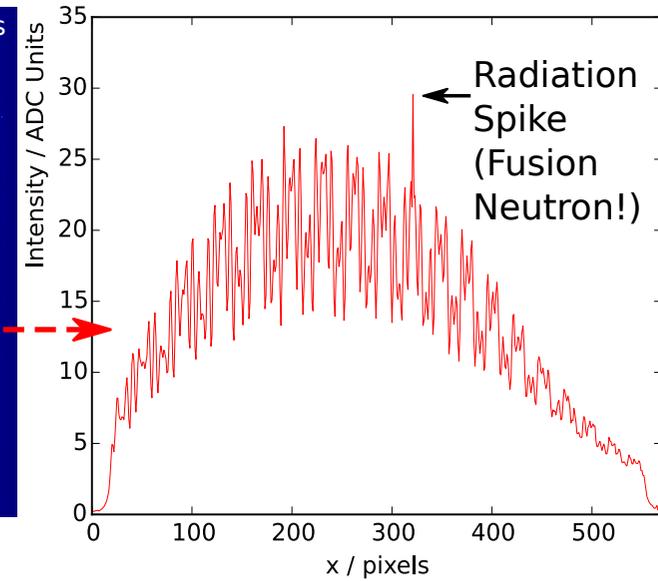
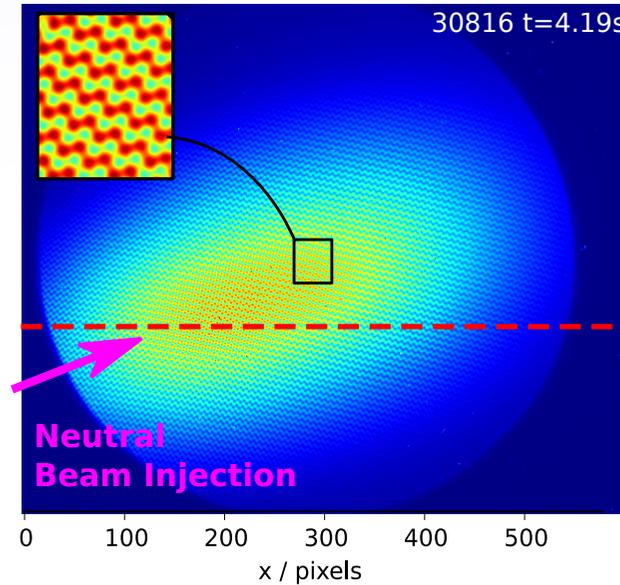


Imaging Motional Stark Effect results

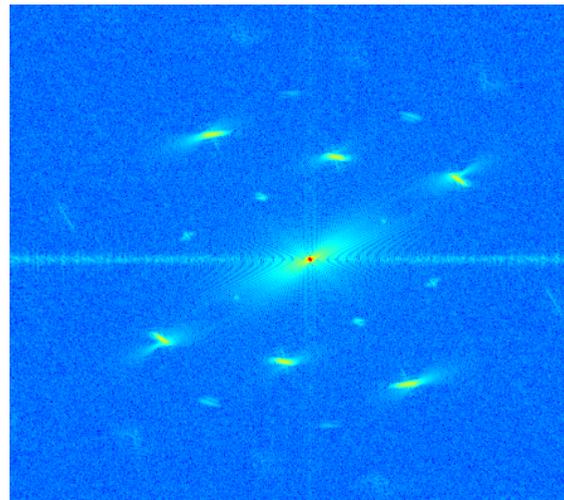
Raw image (without neutral beam)



Raw image (with neutral beam)



Fourier transform



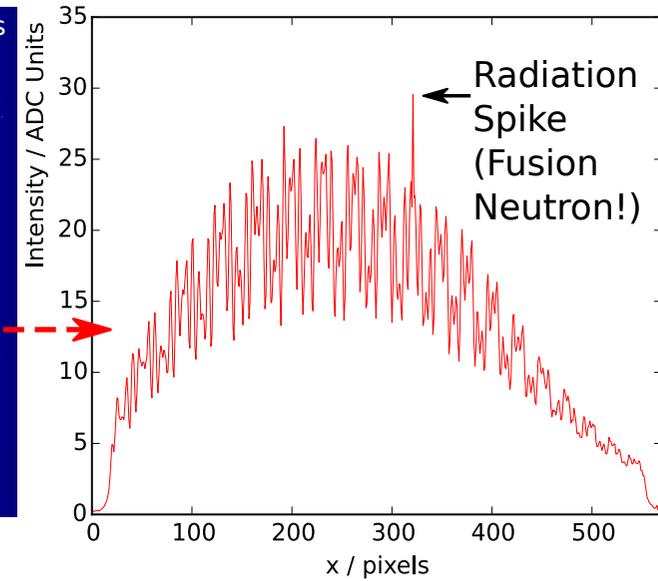
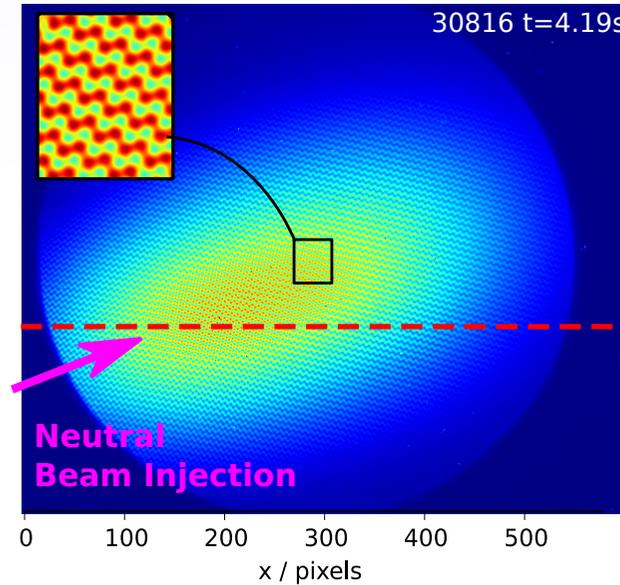
$$\begin{aligned}
 I \propto & 1 + \zeta \cos 2\theta \cos(x) \\
 & + \zeta \sin 2\theta \cos(x - y) \\
 & - \zeta \sin 2\theta \cos(x + y)
 \end{aligned}$$

Imaging Motional Stark Effect results

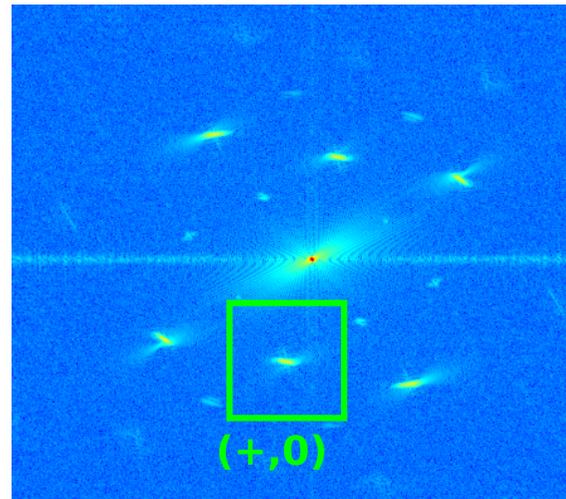
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Fourier transform



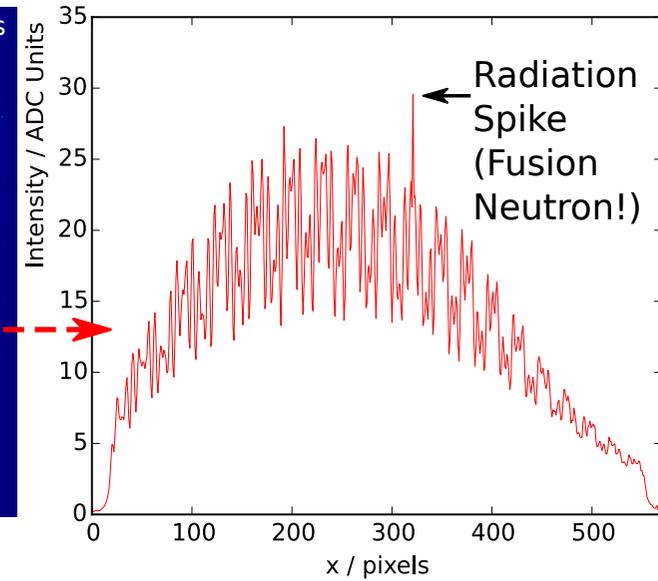
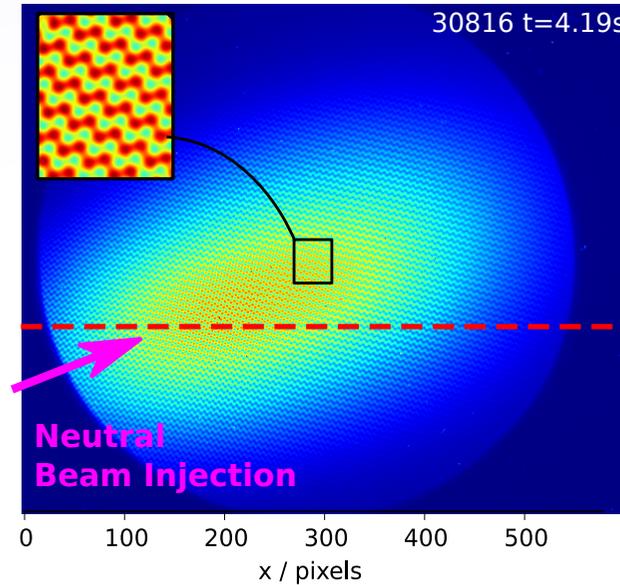
$$I \propto 1 + \zeta \cos 2\theta \cos(x) + \zeta \sin 2\theta \cos(x - y) - \zeta \sin 2\theta \cos(x + y)$$

Imaging Motional Stark Effect results

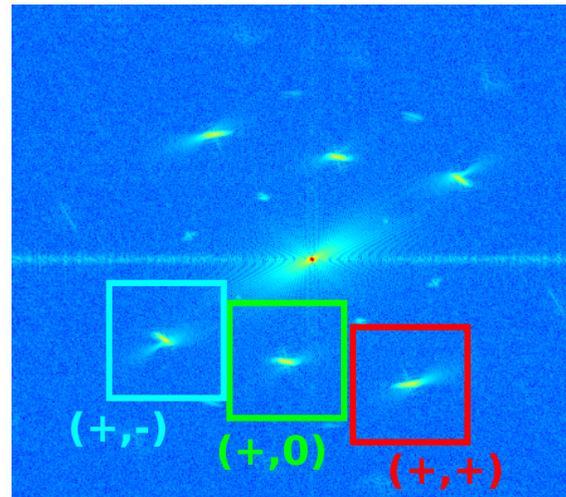
Raw image (without neutral beam)



Raw image (with neutral beam)



Fourier transform



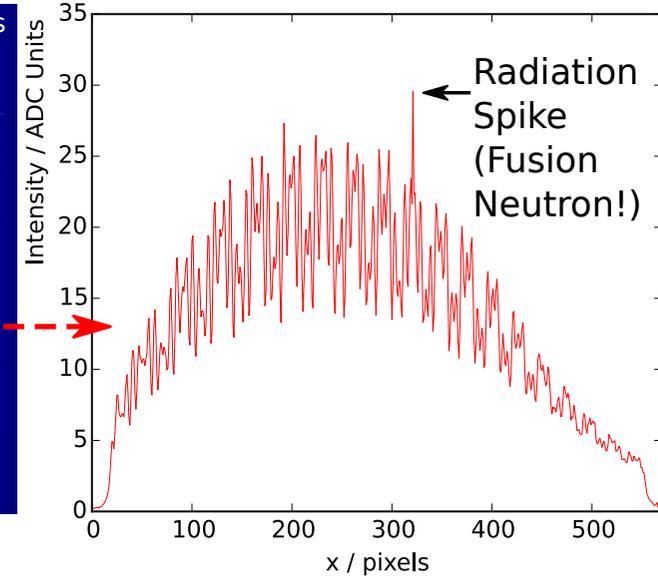
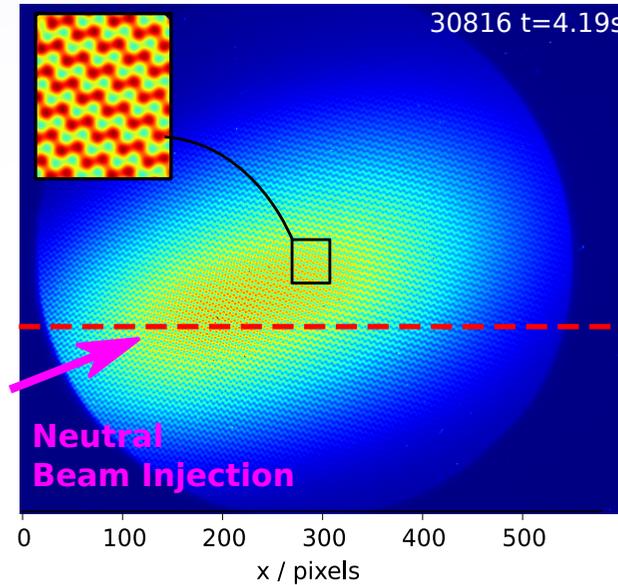
$$I \propto 1 + \zeta \cos 2\theta \cos(x) + \zeta \sin 2\theta \cos(x - y) - \zeta \sin 2\theta \cos(x + y)$$

Imaging Motional Stark Effect results

Raw image (without neutral beam)

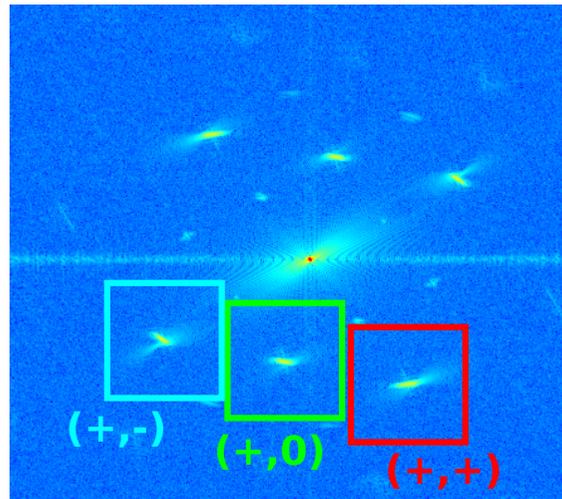


Raw image (with neutral beam)

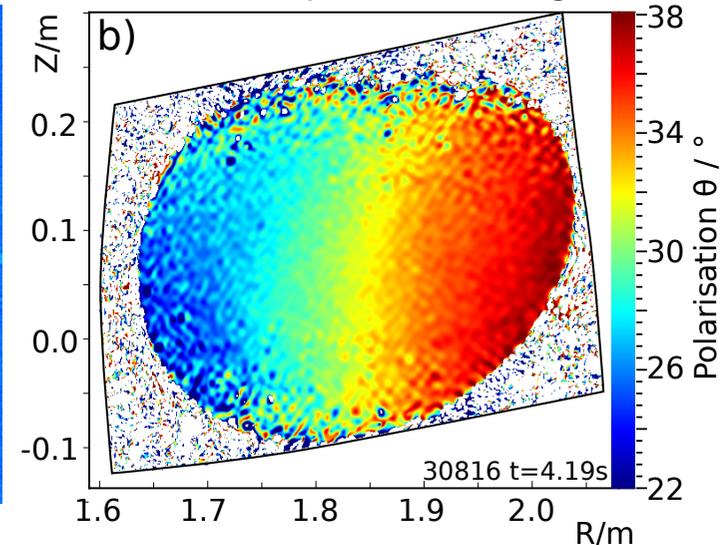


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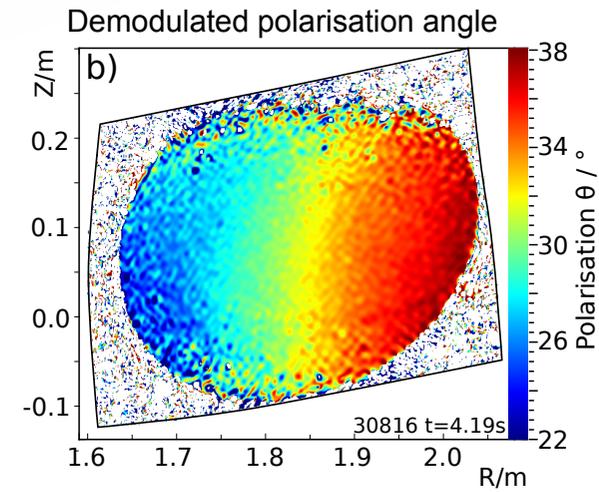
Demodulated polarisation angle



Imaging Motional Stark Effect results

What does θ tell us about j ?

Maybe we have enough data now to image j directly?



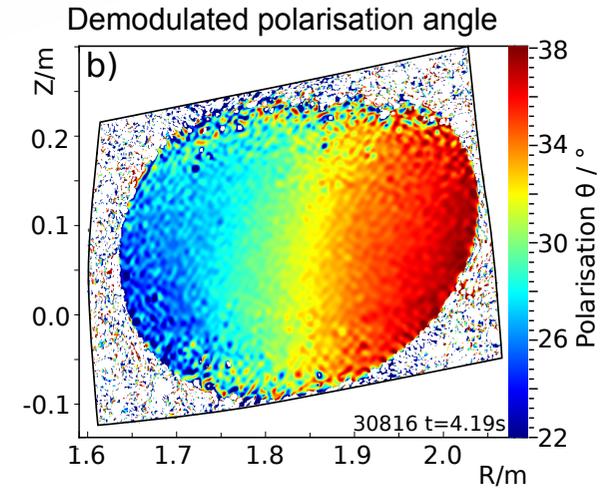
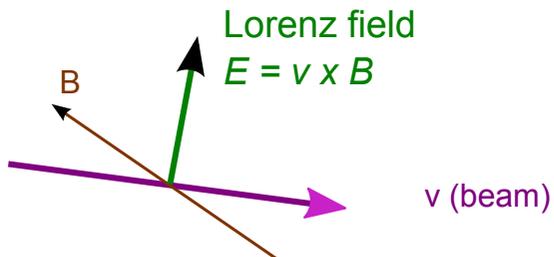
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Imaging Motional Stark Effect results

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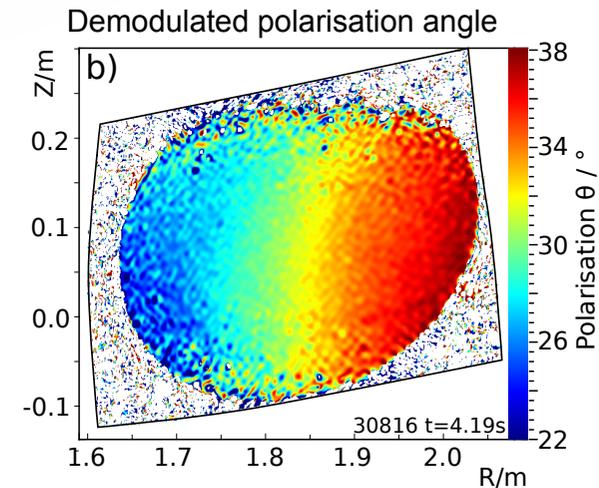
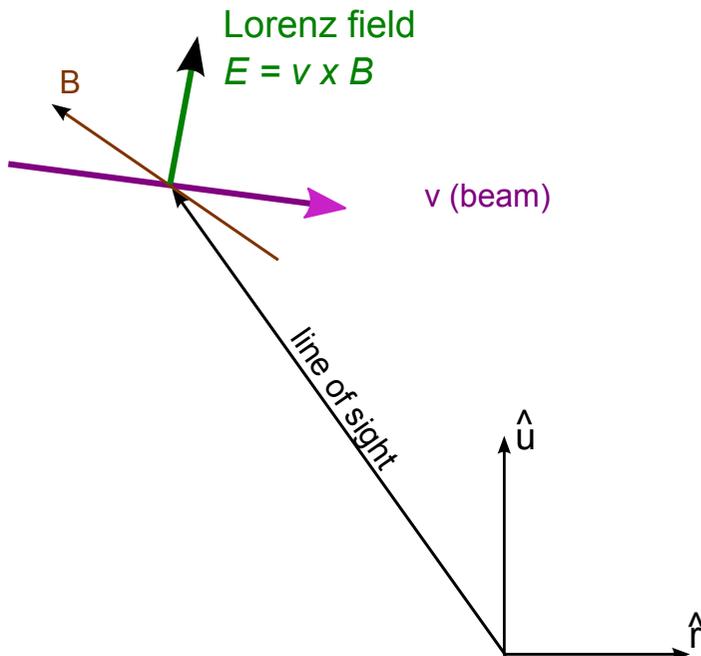
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$$\tan \theta = \frac{(\underline{v} \times \underline{B}) \cdot \hat{r}}{(\underline{v} \times \underline{B}) \cdot \hat{u}}$$



Imaging Motional Stark Effect results

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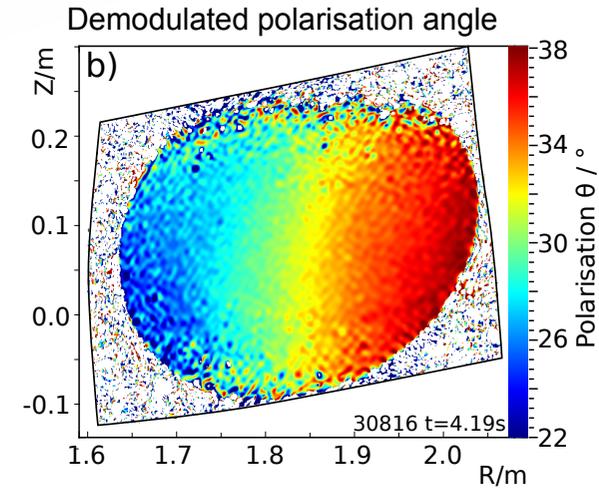
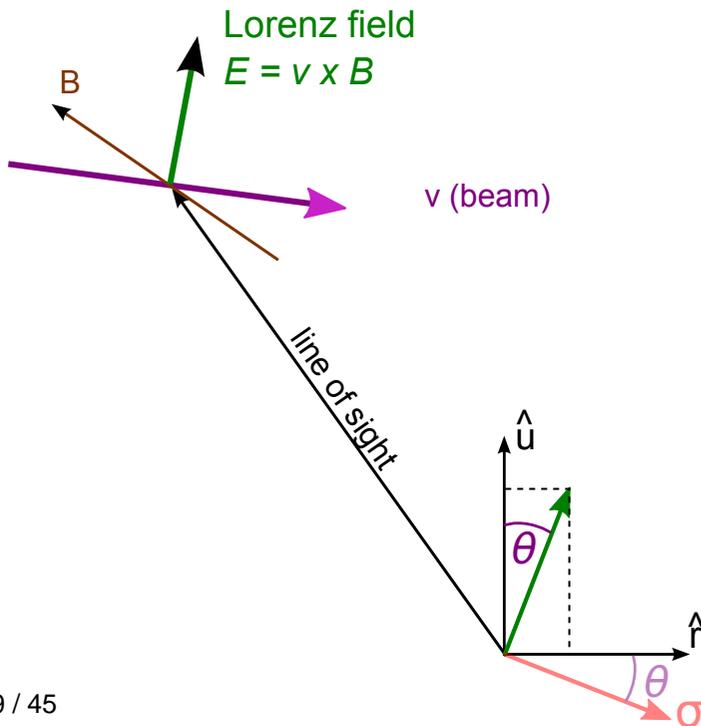
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Imaging Motional Stark Effect results

What does θ tell us about j ?

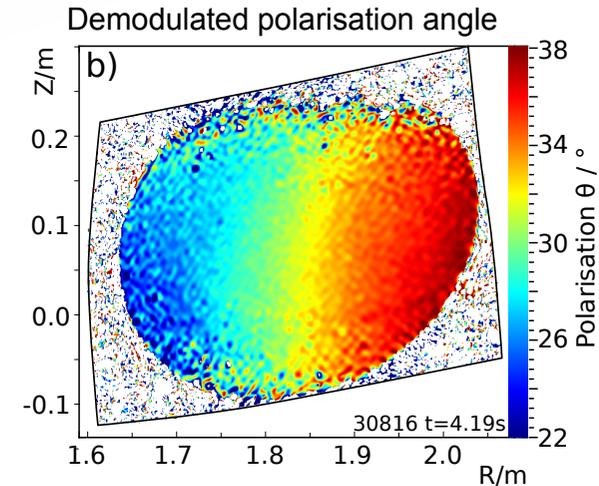
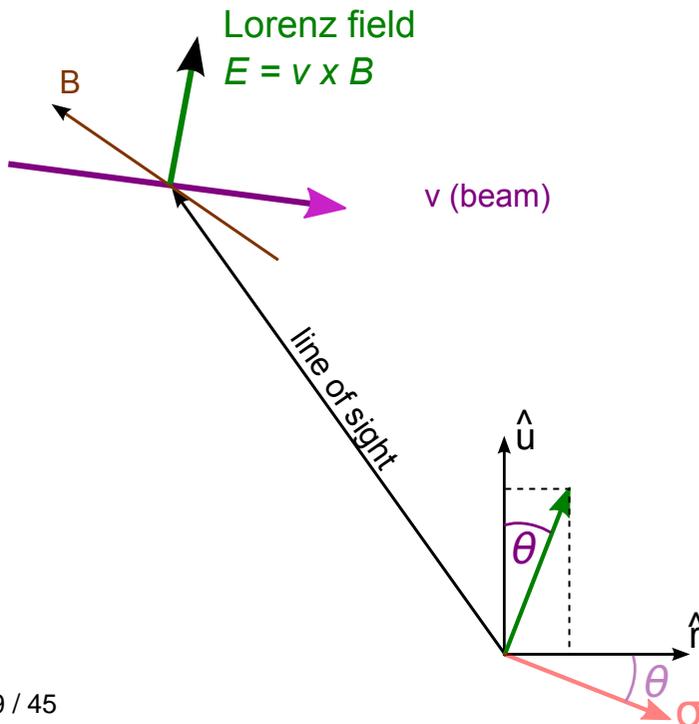
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$$\tan \theta = \frac{(\underline{v} \times \underline{B}) \cdot \hat{r}}{(\underline{v} \times \underline{B}) \cdot \hat{u}}$$



Choose offset to θ , such that $\theta = 0^\circ$ at magnetic axis ($u = v \times \phi$)

$$\tan \theta = \frac{(\underline{v} \times \underline{\phi} \cdot \hat{r})}{(\underline{v} \times \underline{\phi} \cdot \hat{u})} + \left[\frac{(\underline{v} \times \underline{\hat{R}} \cdot \hat{r})}{(\underline{v} \times \underline{\phi} \cdot \hat{u})} - \frac{(\underline{v} \times \underline{\phi} \cdot \hat{r})}{(\underline{v} \times \underline{\phi} \cdot \hat{u})} \frac{(\underline{v} \times \underline{\hat{R}} \cdot \hat{u})}{(\underline{v} \times \underline{\phi} \cdot \hat{u})} \right] \frac{B_R}{B_\phi} + \left[\frac{(\underline{v} \times \underline{\hat{Z}} \cdot \hat{r})}{(\underline{v} \times \underline{\phi} \cdot \hat{u})} - \frac{(\underline{v} \times \underline{\phi} \cdot \hat{r})}{(\underline{v} \times \underline{\phi} \cdot \hat{u})} \frac{(\underline{v} \times \underline{\hat{Z}} \cdot \hat{u})}{(\underline{v} \times \underline{\phi} \cdot \hat{u})} \right] \frac{B_Z}{B_\phi}$$

Imaging Motional Stark Effect results

What does θ tell us about j ?

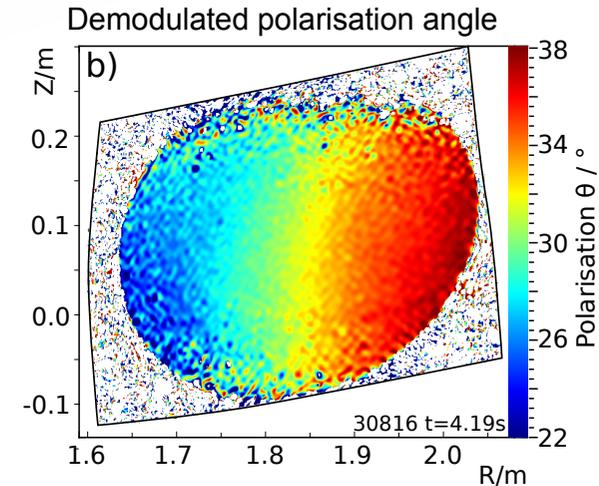
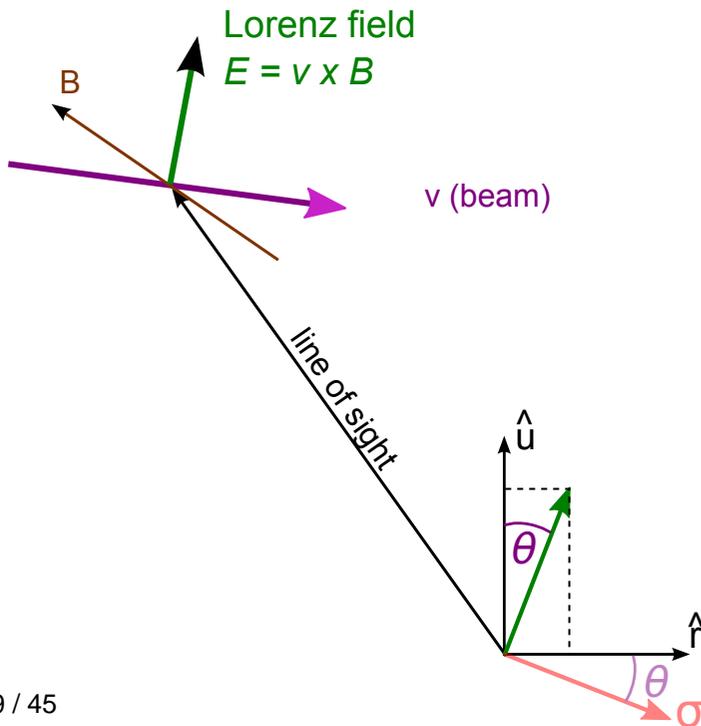
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$$\tan \theta = 0.6 \frac{B_Z}{B_\phi}$$

Imaging Motional Stark Effect results

What does θ tell us about j ?

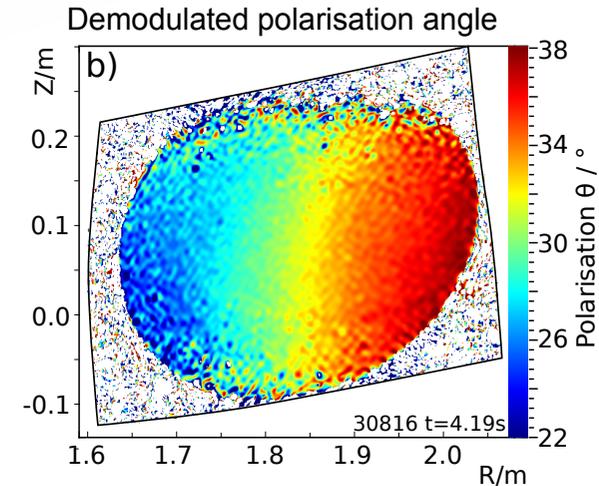
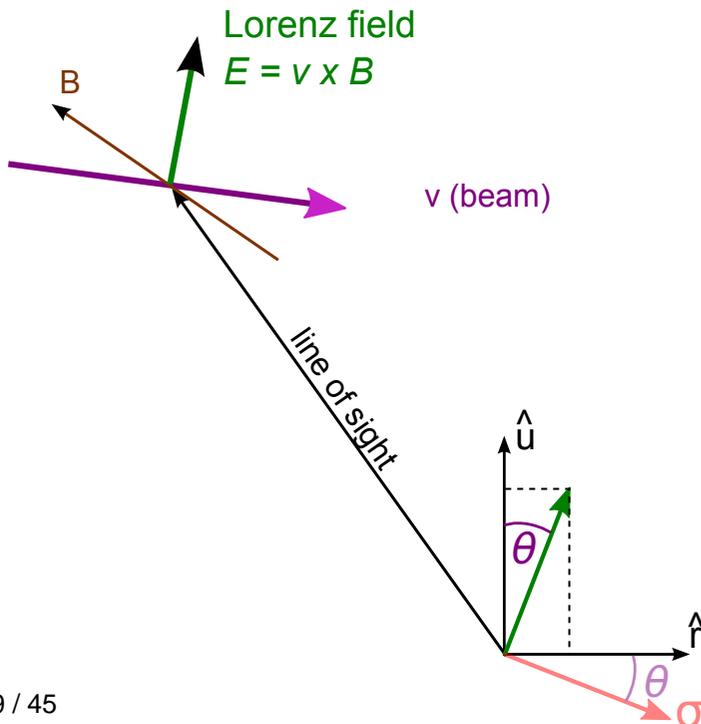
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Magnetic field pitch angle.
Gives B_z since B_ϕ is well known.

Imaging Motional Stark Effect results

What does θ tell us about j ?

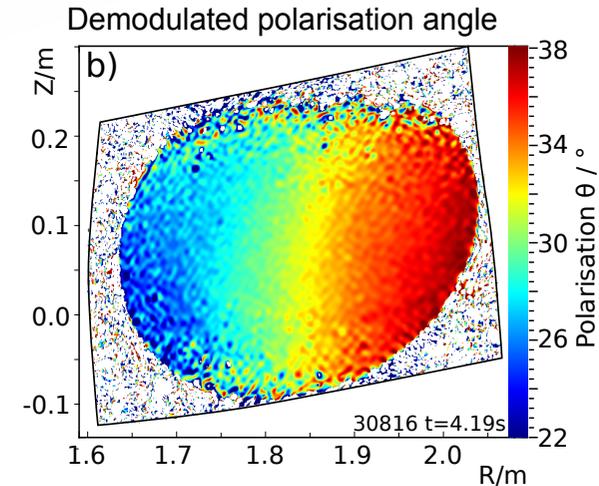
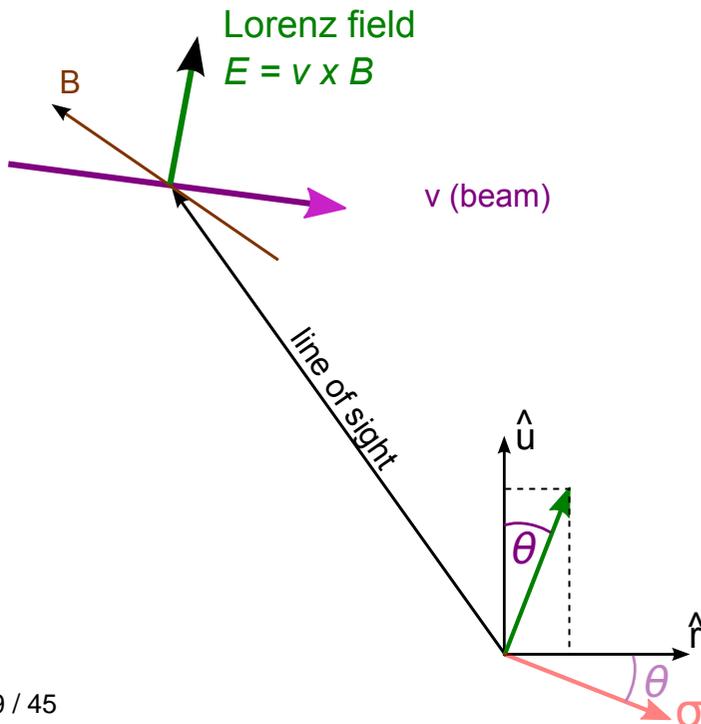
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$\tan \theta = 0.6 \frac{B_Z}{B_\phi}$ — Magnetic field pitch angle.
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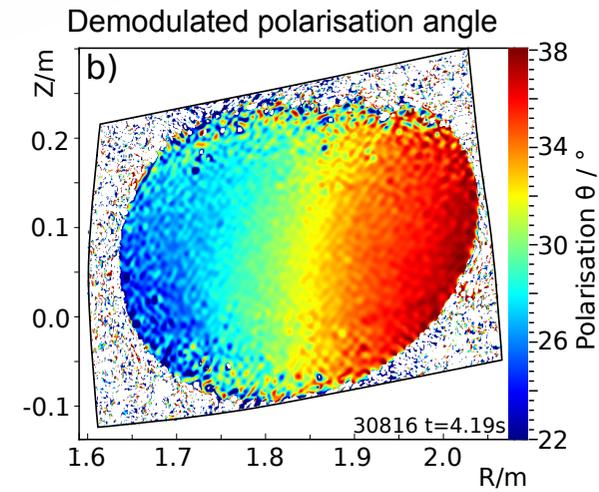
We can now make a B_z image:

$$B_Z \approx \frac{1.7 B_0}{R} \theta$$

$$B_0 = B_\phi \text{ at } R=1\text{m}$$

Imaging Motional Stark Effect results

What does B_z tell us about j in the core?

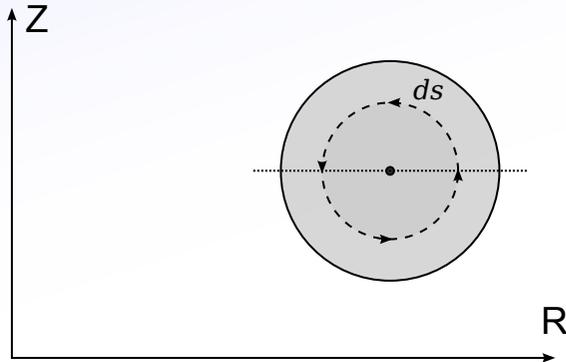


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Imaging Motional Stark Effect results

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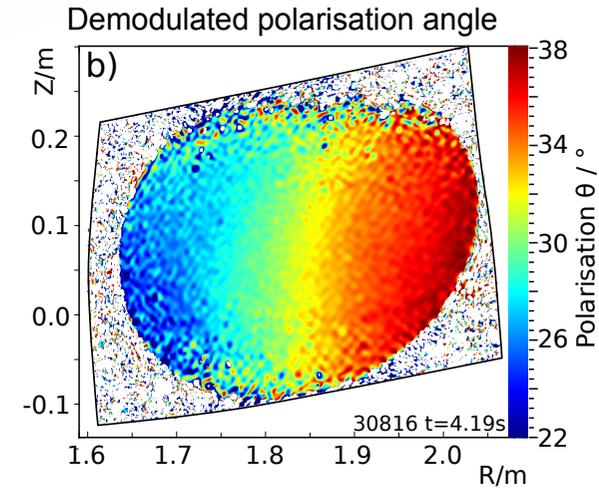
Large aspect ratio approximation (assume core is a cylinder)



$$\int B \cdot ds = \mu_0 \int j dA$$

$$B = \frac{1}{2} \mu_0 j r$$

$$\frac{dB_z}{dR} = \frac{1}{2} \mu_0 j$$

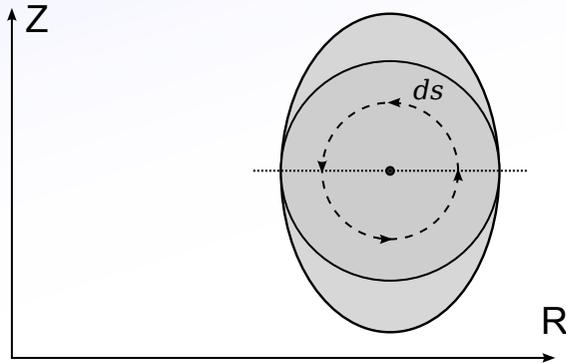


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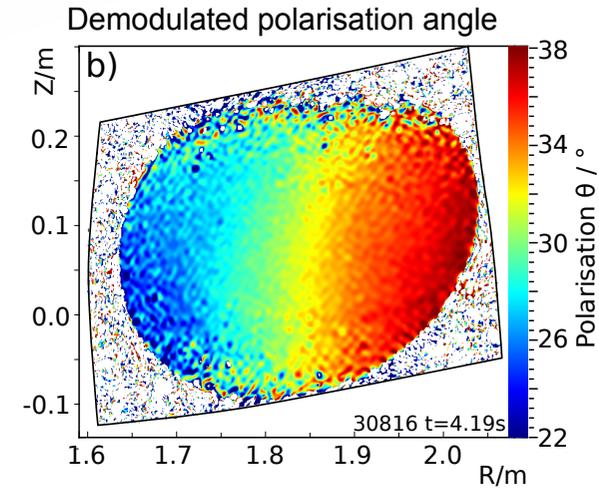
[CC.Petty Nucl. Fus. 2002]

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$$\mu_0 j \approx - \left(1 + \frac{1}{\kappa^2} \right) \frac{dB_z}{dR}$$

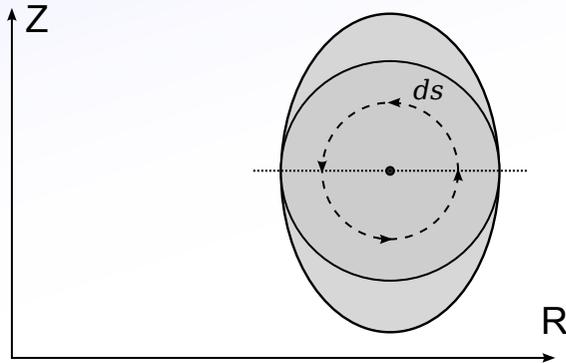


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Imaging Motional Stark Effect results

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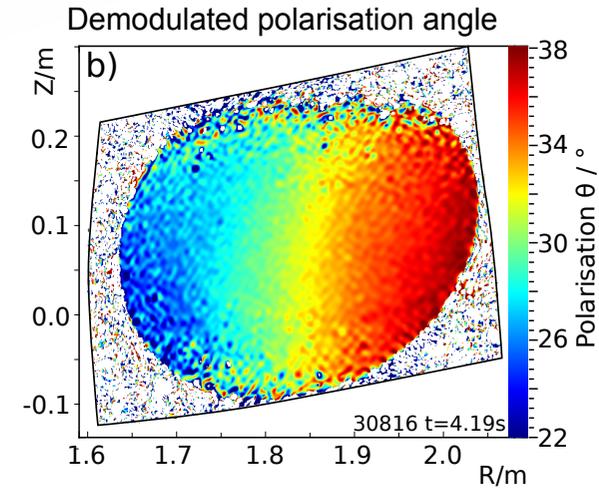
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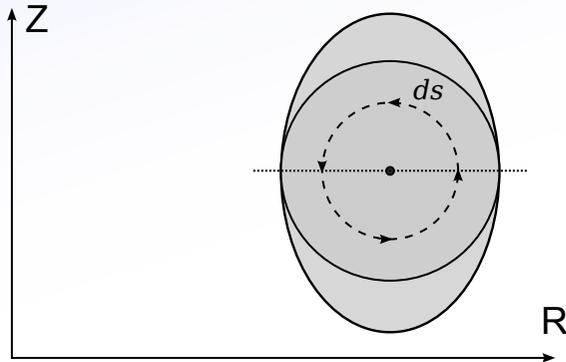


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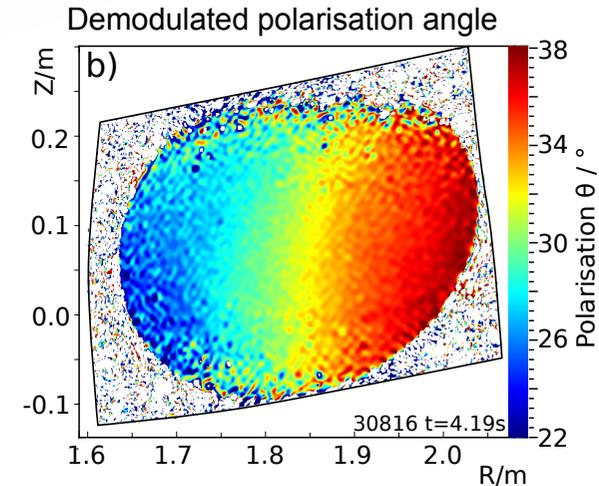
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To first order, local j relates to local derivative of measurement

This is only an approximation! ... but we now understand that $d\theta/dR$ holds the information about j .

What can we see in $d\theta/dR$ at the axis?



$$B_z \approx \frac{1.7B_0}{R} \theta$$

Typically, in core:

$$R \sim 1.6m$$

$$\theta < 5^\circ$$

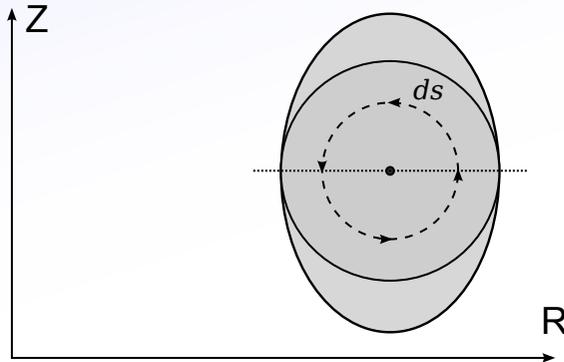
$$d\theta/dR \sim 35^\circ m^{-1}$$

$$d\theta/dR \gg (\theta/R)$$

Imaging Motional Stark Effect results

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$$\int B \cdot ds = \mu_0 \int j dA$$

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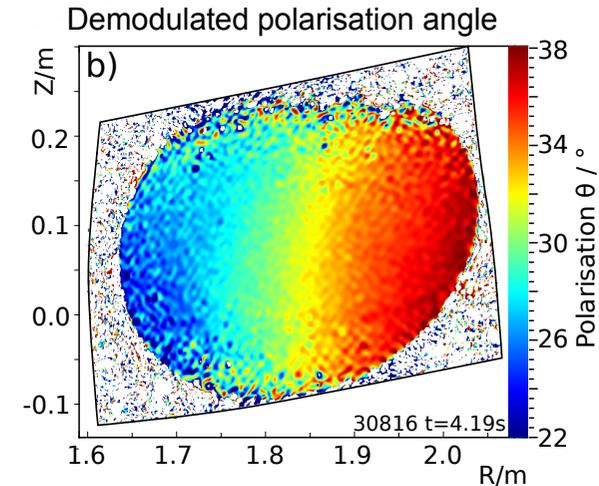
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[CC.Petty Nucl. Fus. 2002]



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Typically, in core:

$$\begin{aligned} R &\sim 1.6m \\ \theta &< 5^\circ \\ d\theta/dR &\sim 35^\circ m^{-1} \\ d\theta/dR &\gg (\theta/R) \end{aligned}$$

This is only an approximation! ... but we now understand that $d\theta/dR$ holds the information about j .

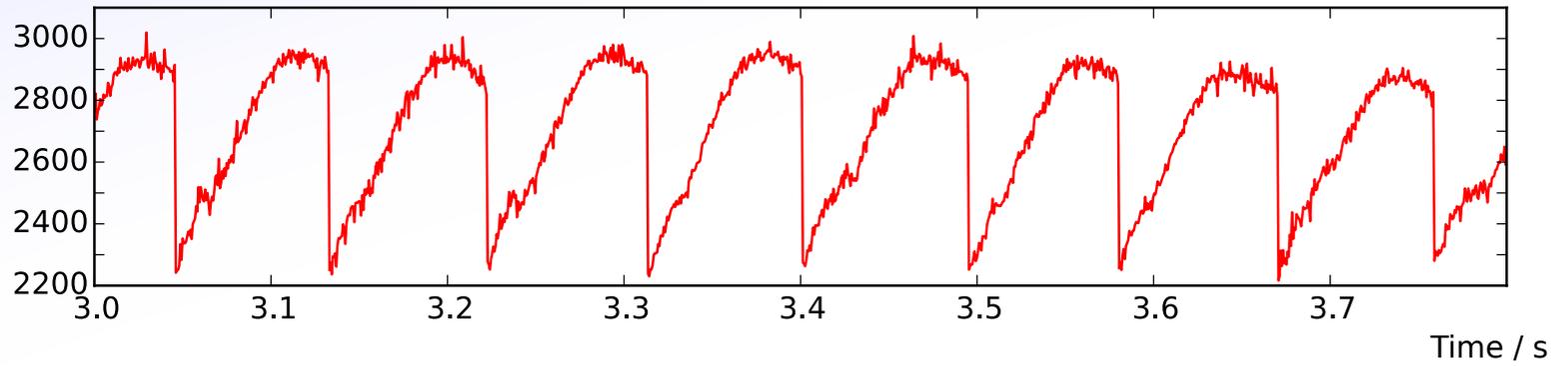
What can we see in $d\theta/dR$ at the axis?

Central safety factor also requires location of centre:

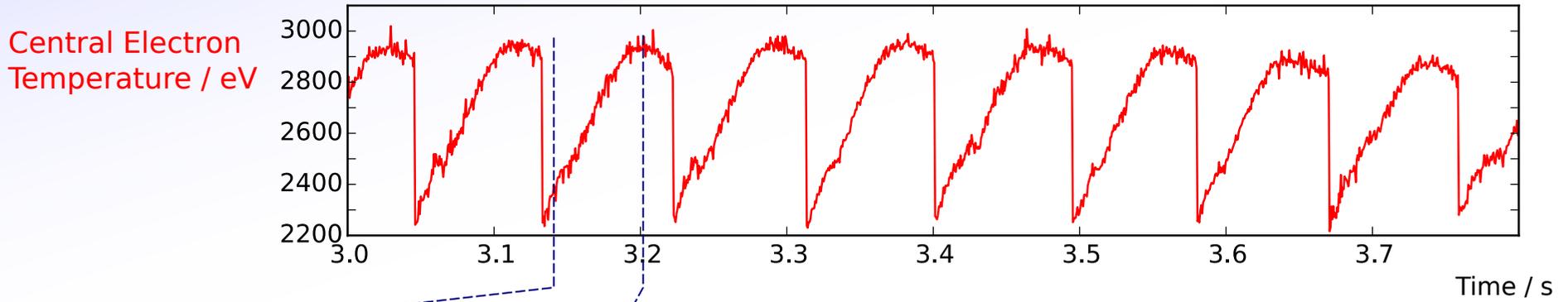
$$q_0 \approx \frac{2B_\phi}{\mu_0 j_0 R} \approx \frac{2B_\phi^{1m}}{\mu_0 j_0 R_0^2}$$

Sawteeth - Magnetic Reconnection

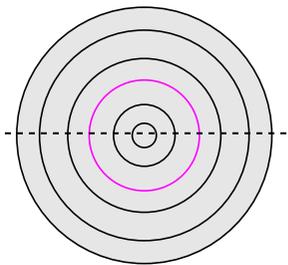
Central Electron
Temperature / eV



Sawteeth - Magnetic Reconnection



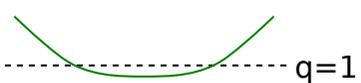
Slow build-up of
pressure and current.



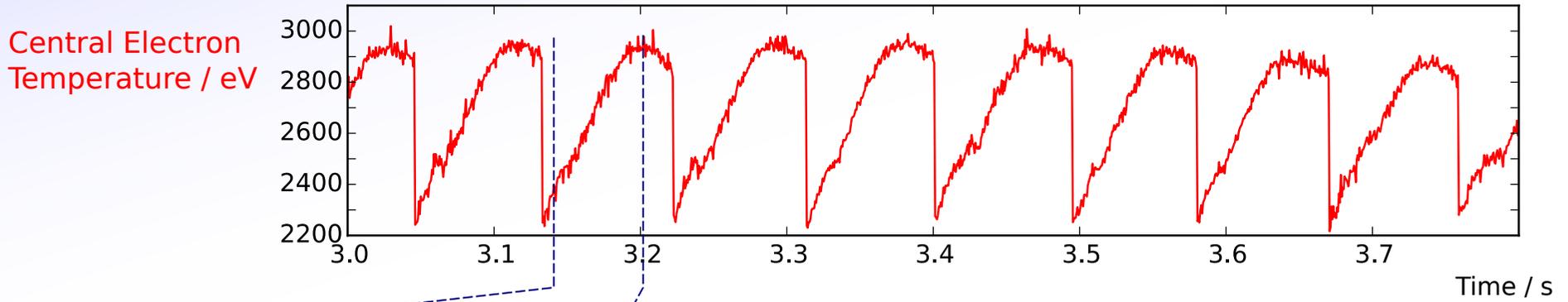
Pressure



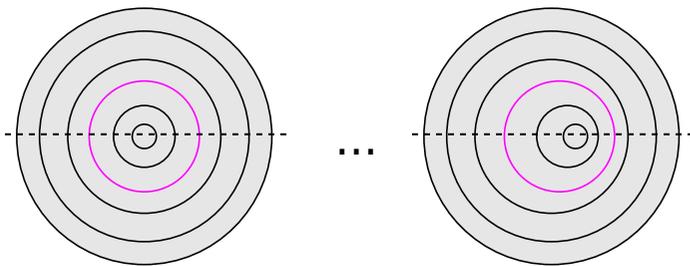
Safety Factor q



Sawteeth - Magnetic Reconnection



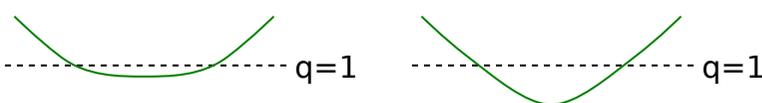
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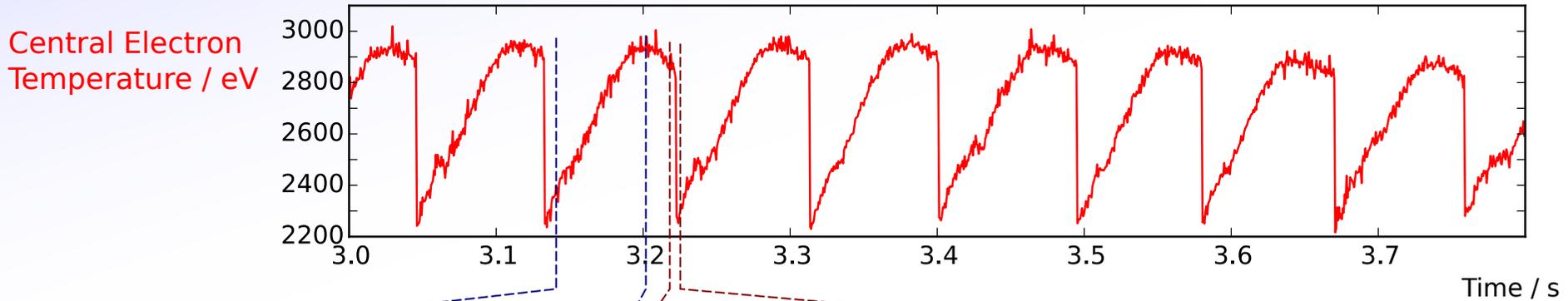
Pressure



Safety Factor q



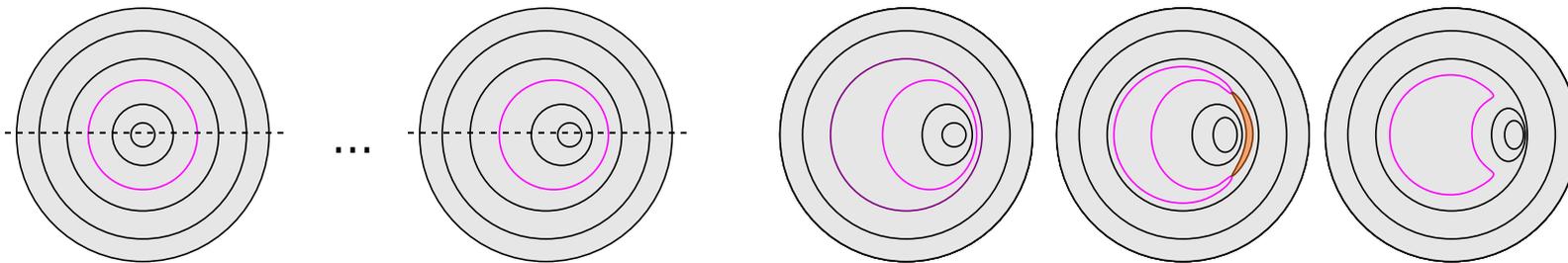
Sawteeth - Magnetic Reconnection



Slow build-up of pressure and current.

Fast magnetic reconnection event redistributes energy and particles outside $q=1$ surface.

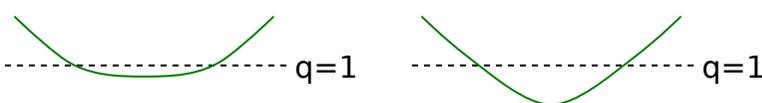
Cycle Repeats



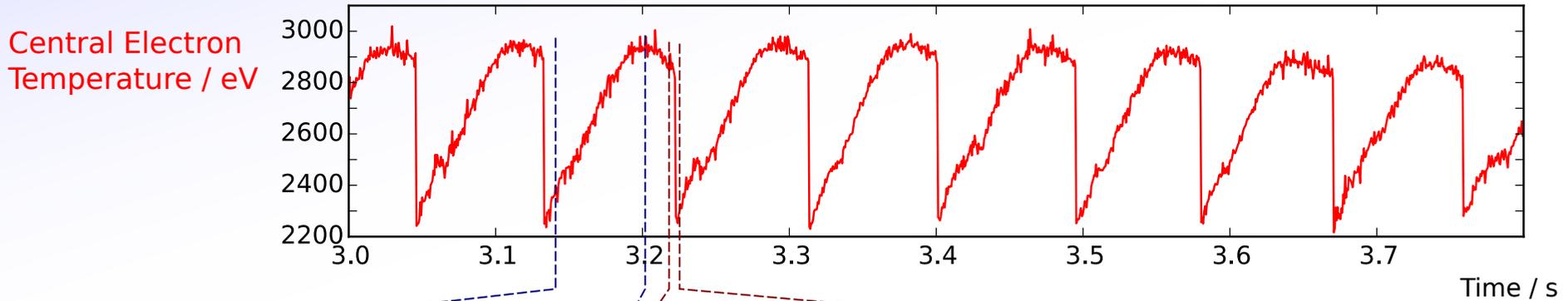
Pressure



Safety Factor q



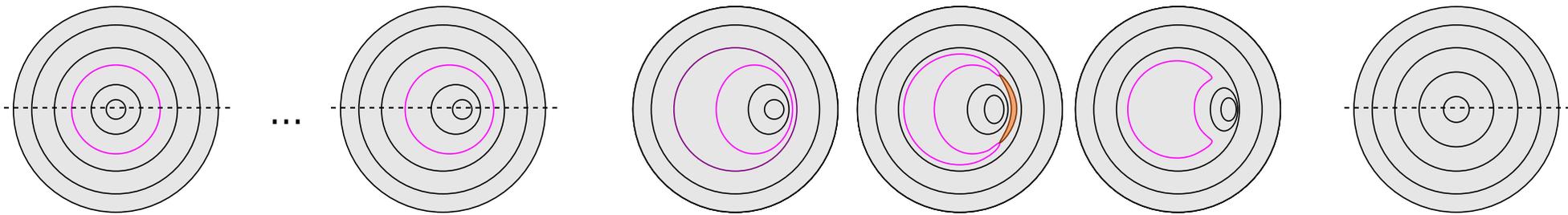
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Cycle Repeats



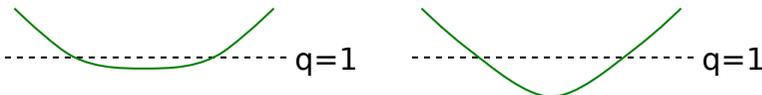
Pressure



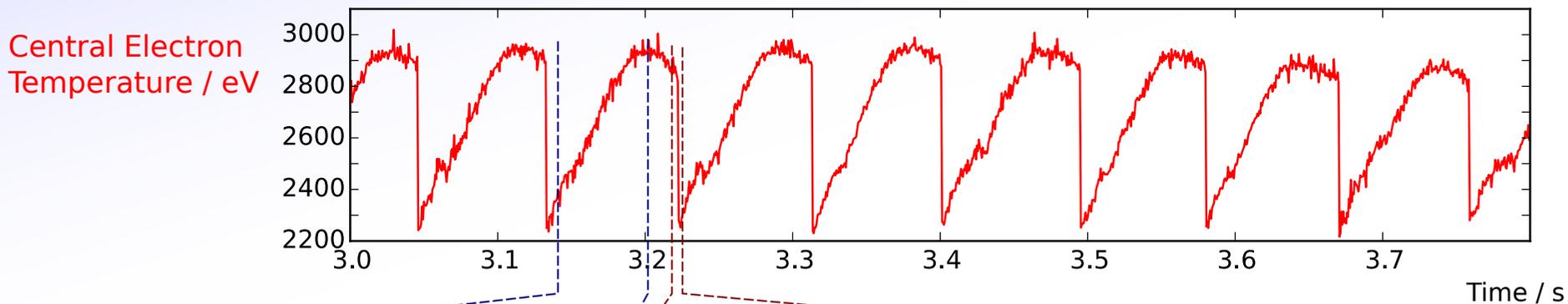
Pressure



Safety Factor q



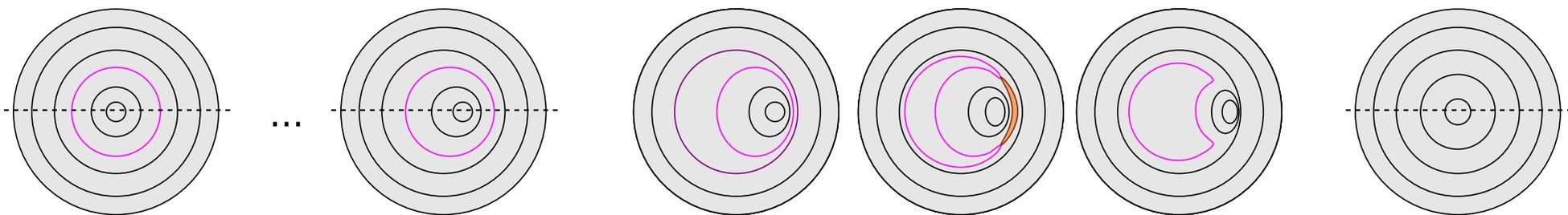
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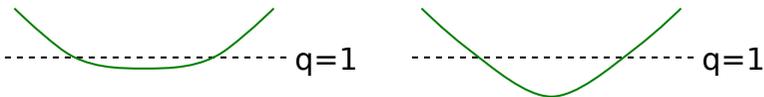
Pressure



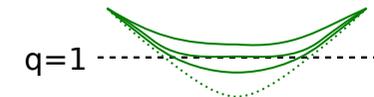
Reconnection observed is much faster than normal models allow (single-fluid MHD).



Safety Factor q



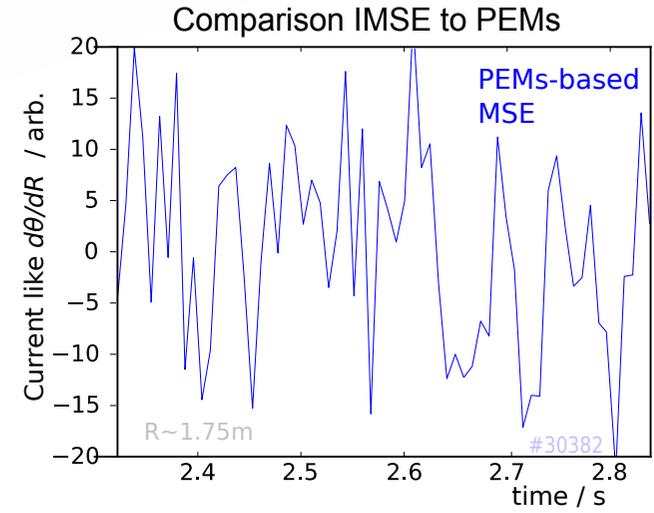
Many new models proposed, e.g. stochastic reconnection.





Sawteeth - Magnetic Reconnection

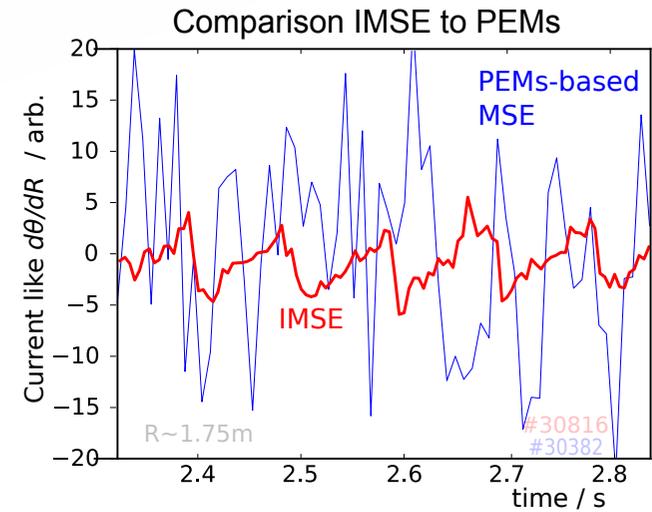
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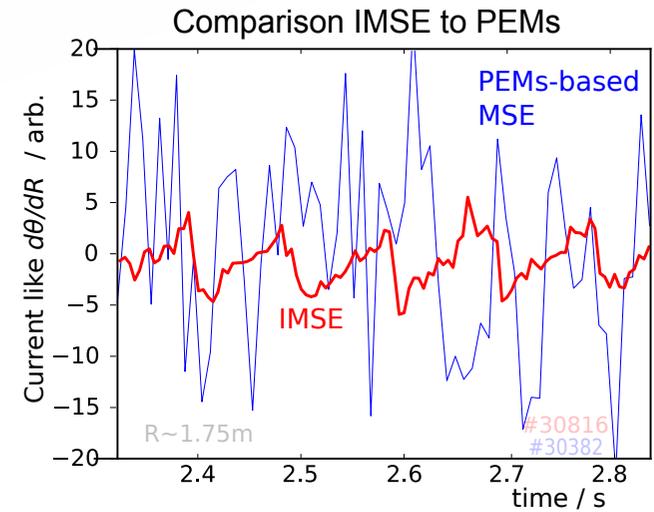
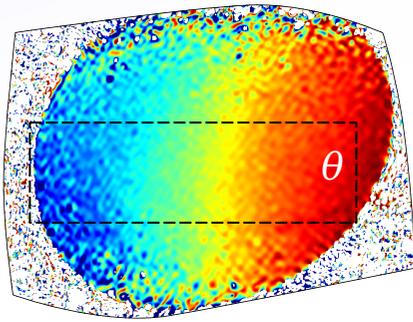
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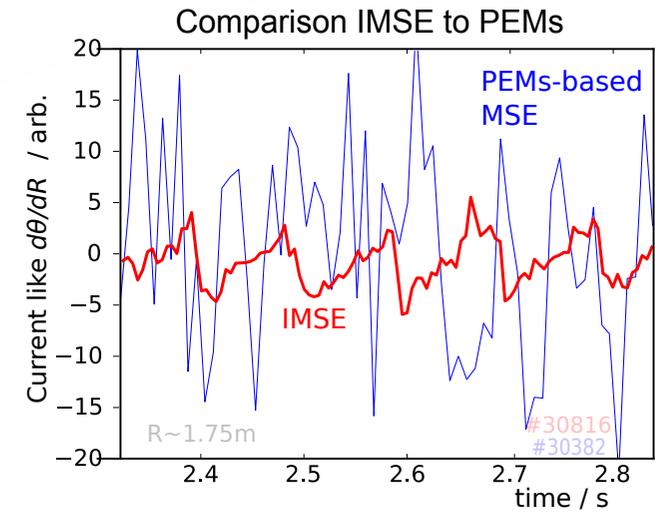
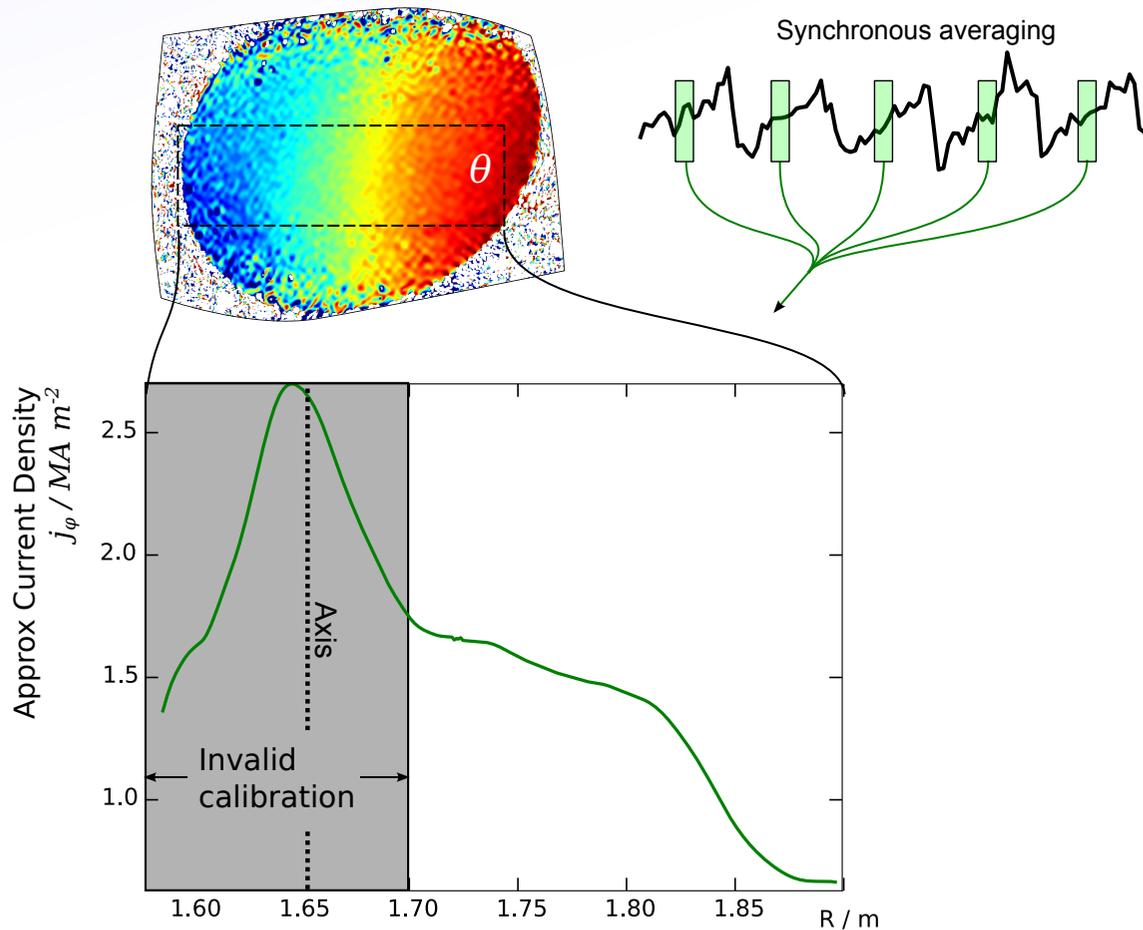
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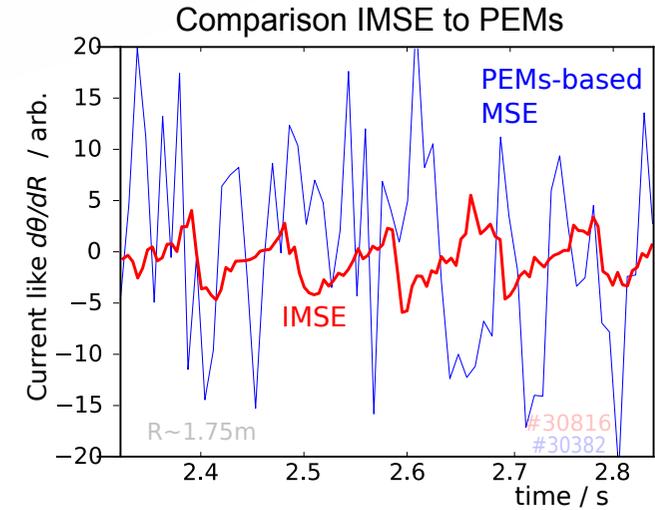
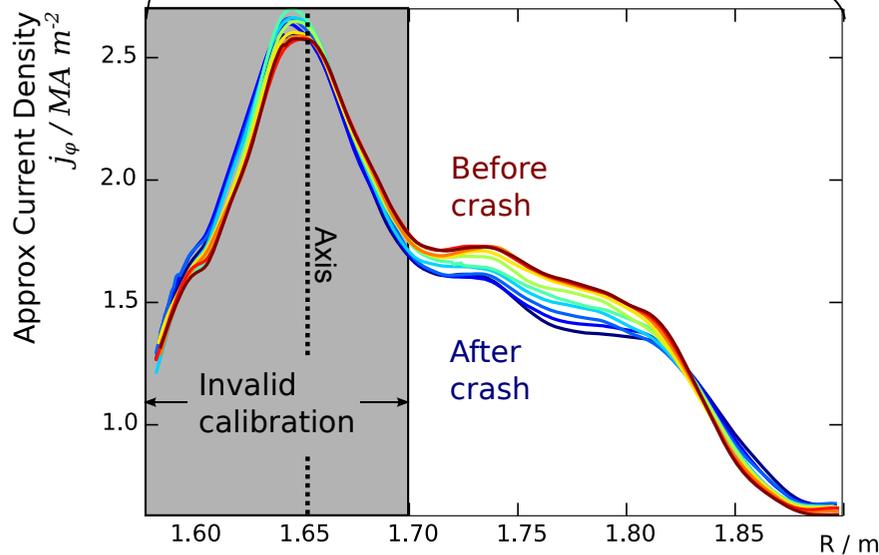
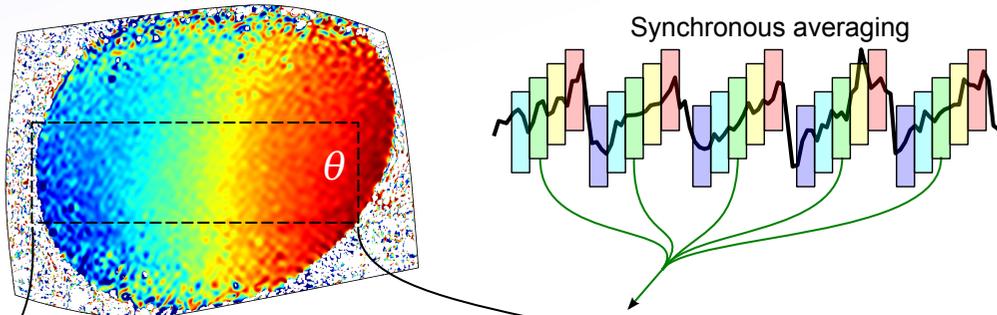
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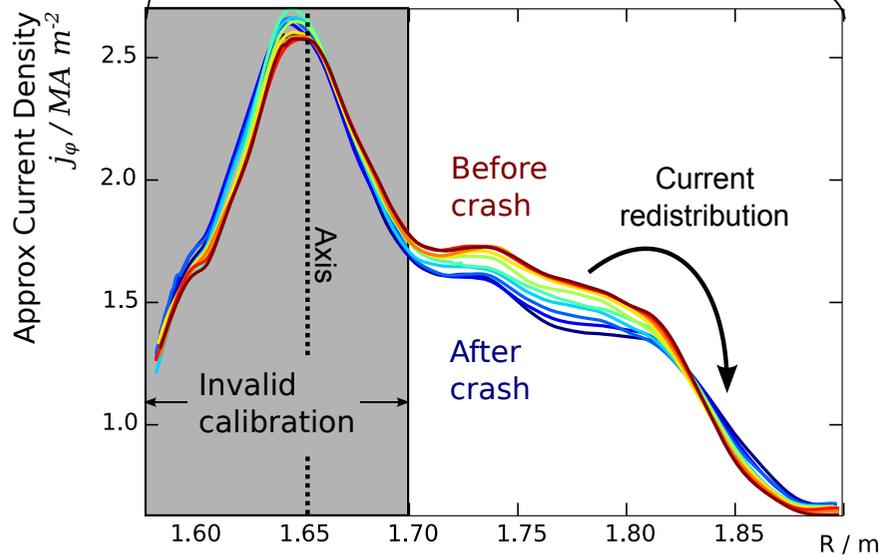
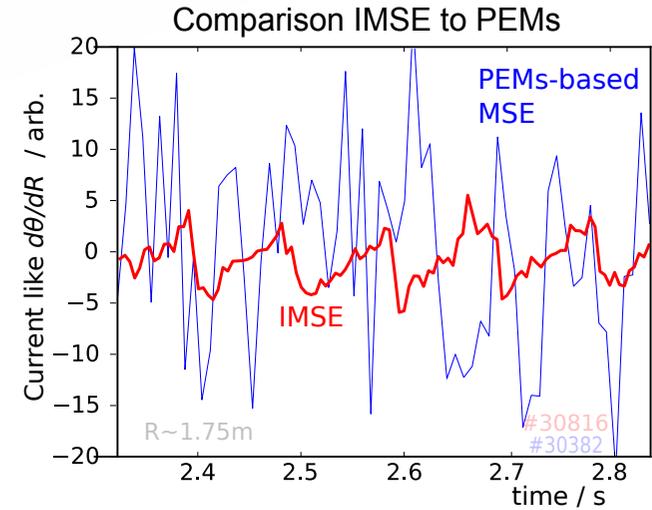
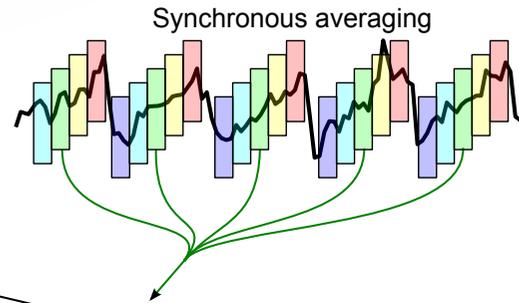
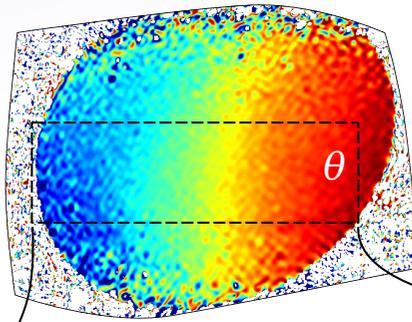
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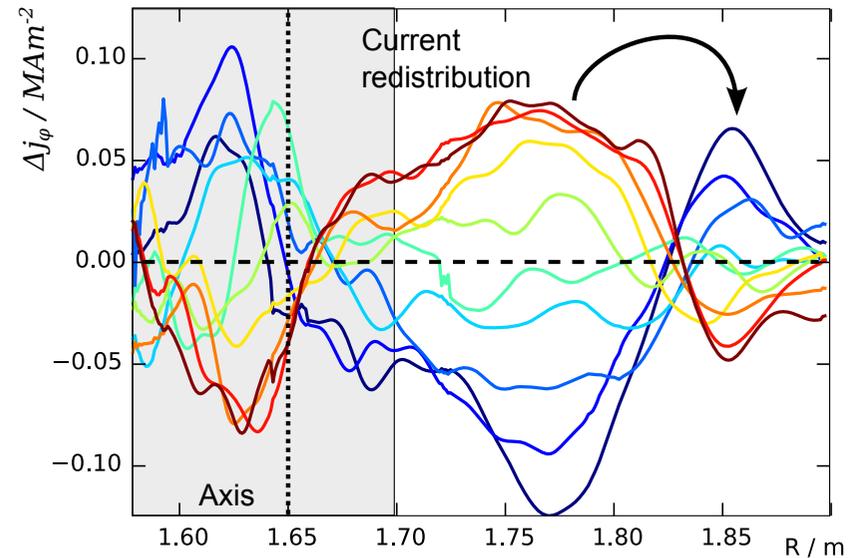
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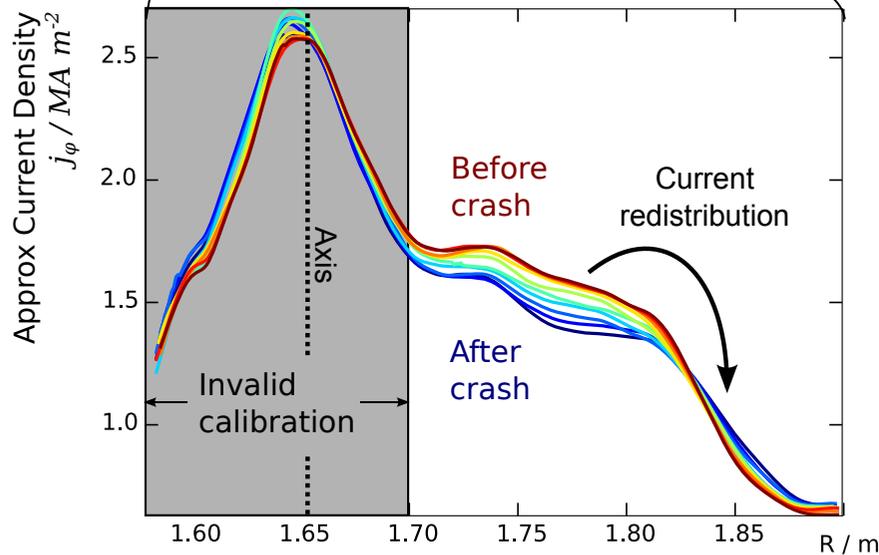
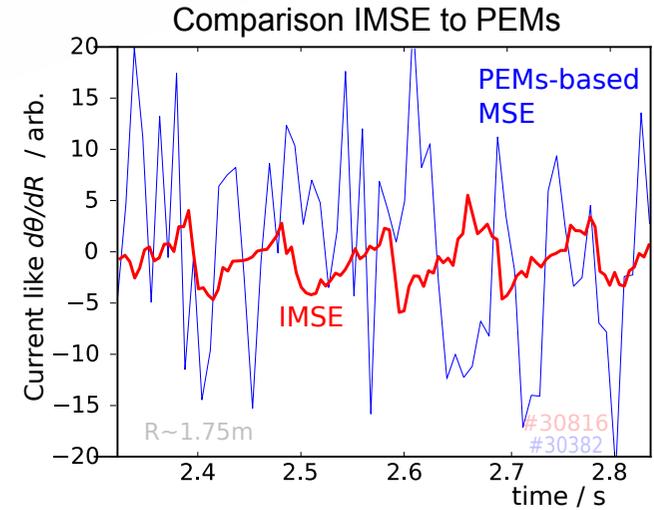
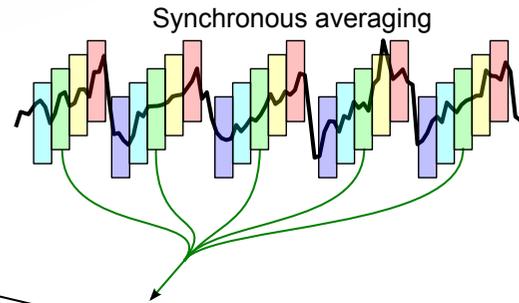
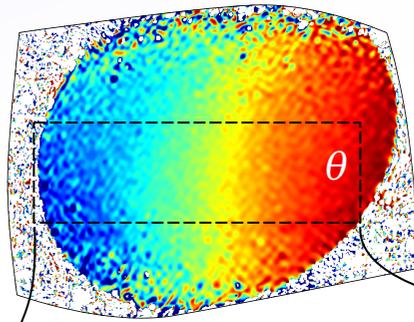
Difference from average profile:



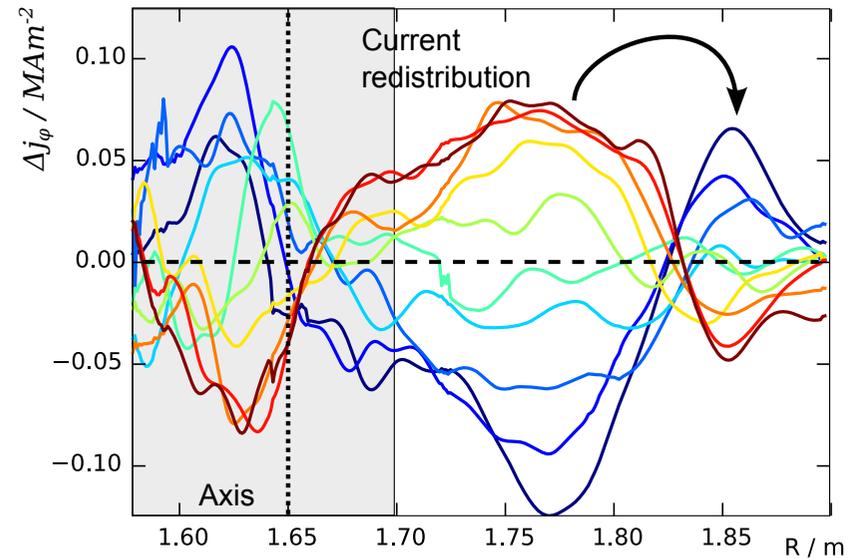
Sawteeth - Magnetic Reconnection

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Difference from average profile:



Current redistribution: $\Delta j \sim 0.050 \text{ MA m}^{-2}$

Measurements every ~3cm (resolution): $\Delta(d\theta/dR) \sim 0.7^\circ \text{ m}^{-1}$

--> $\Delta\theta \pm 0.02^\circ$ required for $\Delta R = 3\text{cm}$



Outline

- Introduction
 - Flux surfaces and current profiles
 - Magnetic equilibrium
 - Bayesian analysis

- Bayesian equilibrium
 - Current-tomography
 - Current tomography + Grad-Shafranov
 - L-Mode reconstructions
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- Internal measurements
 - Motional Stark effect.
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 - Direct j_φ imaging.

- Rigorous determination of uncertainty
 - Too computationally intensive for H-mode
 - Need internal measurements!

- Excellent internal measurements.
 - Good dynamics from very approximate derivation of Δj_φ
 - Calibration very difficult to required accuracy.
 - Need to include in equilibrium

- Integrated Data Analysis
 - Current diffusion
 - IMSE results in comparison
 - Sawtooth models



Integrated Data Analysis - Equilibrium

New approach to equilibrium at ASDEX Upgrade:

- Grad-Shafranov solver, but with rigorous treatment of errors
- Try to mitigate effect of nonphysical regularisation with as much realistic information as possible:
 - Pressure constraints: n_e , T_e , T_i , Z_{eff} , fast-ions (from modeling)
 - Geometric information (Inboard/outboard agreement of diagnostics)

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Current Diffusion Equation (CDE):

$$\sigma_{\parallel} \frac{\partial \psi}{\partial t} = \frac{R_0 J^2}{\mu_0 \rho} \frac{\partial}{\partial \rho} \left(\frac{G_2}{J} \frac{\partial \psi}{\partial \rho} \right) - \frac{V'}{2\pi\rho} (j_{bs} + j_{cd})$$

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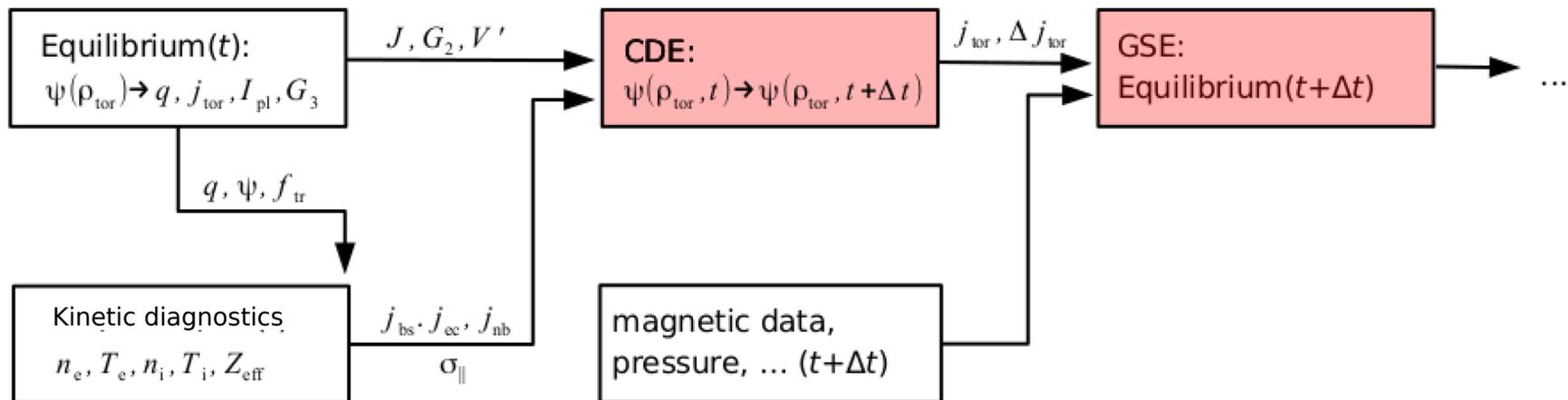
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Bootstrap current

Current drive (ECCD, NBI etc)

Provides a weak constraint on j_{φ} from expected evolution from previous time-points.

- i.e. physically realistic (and informative) prior information.





Integrated Data Analysis Equilibrium

Example: Counter-current Electron Cyclotron Current Drive (ECCD)

ECCD drives localised on-axis current

- Not seen by magnetics (small due to low area of centre)
- No effect on pressure profile = not seen by kinetic inputs

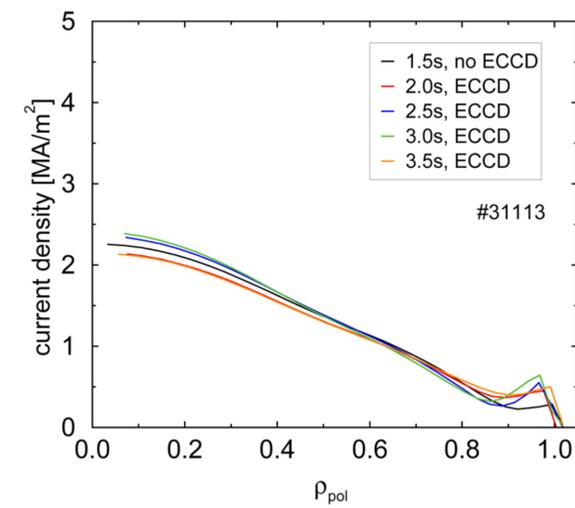
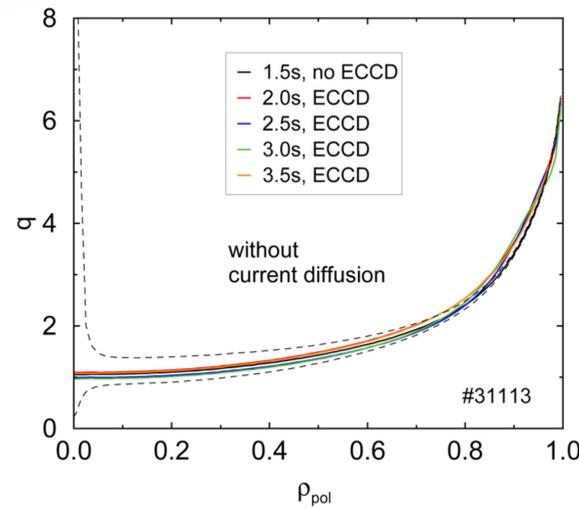
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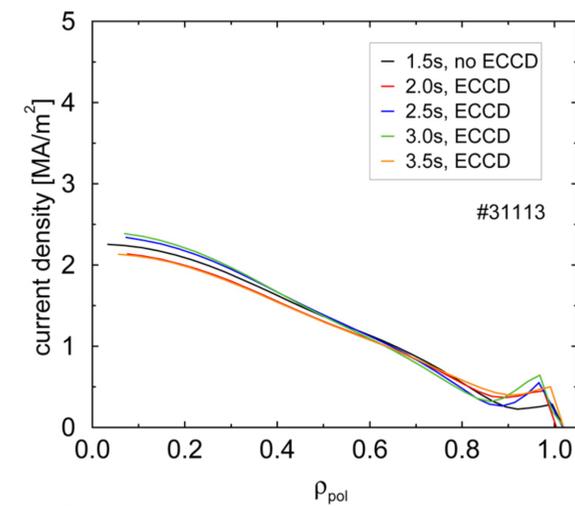
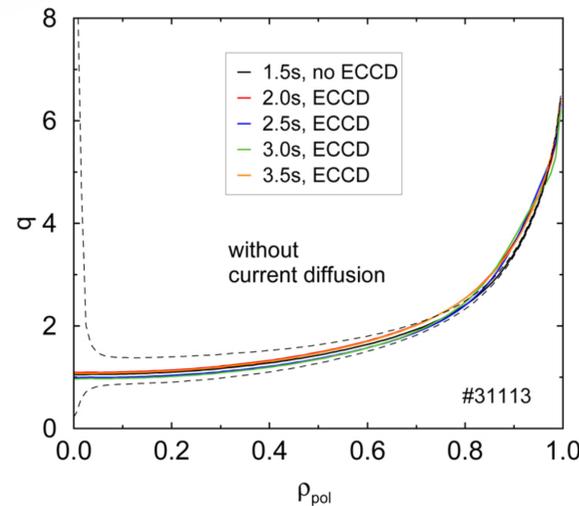
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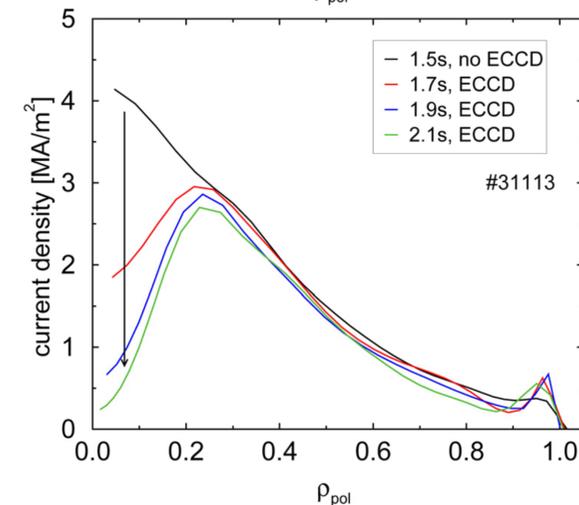
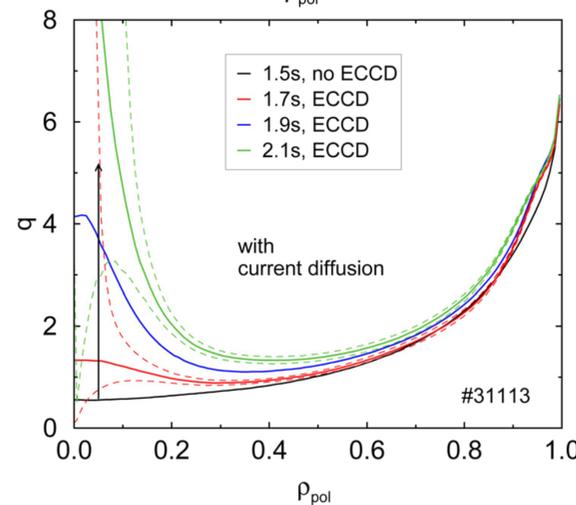
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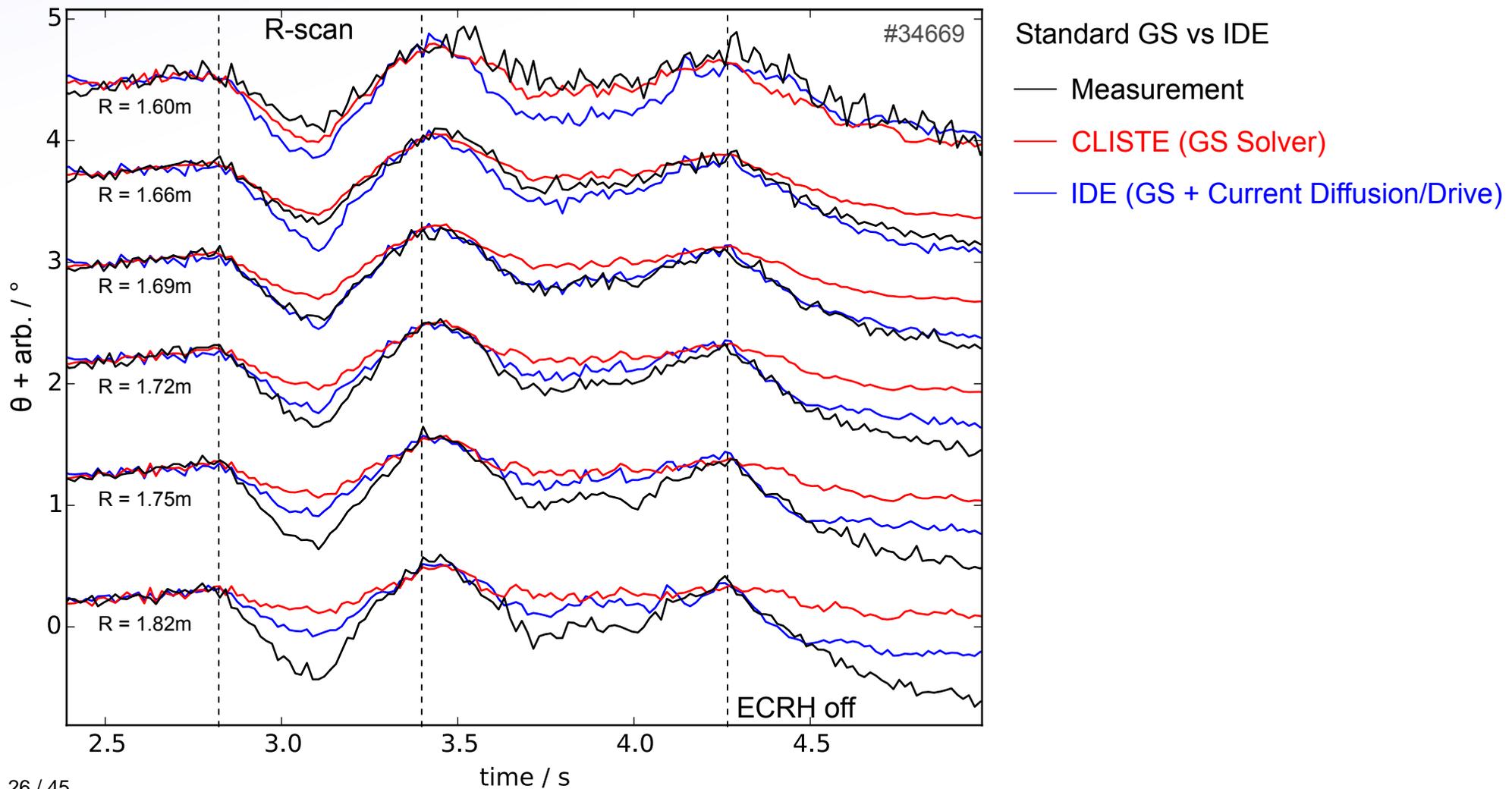
With CDE + Current Drive:



Integrated Equilibrium vs IMSE

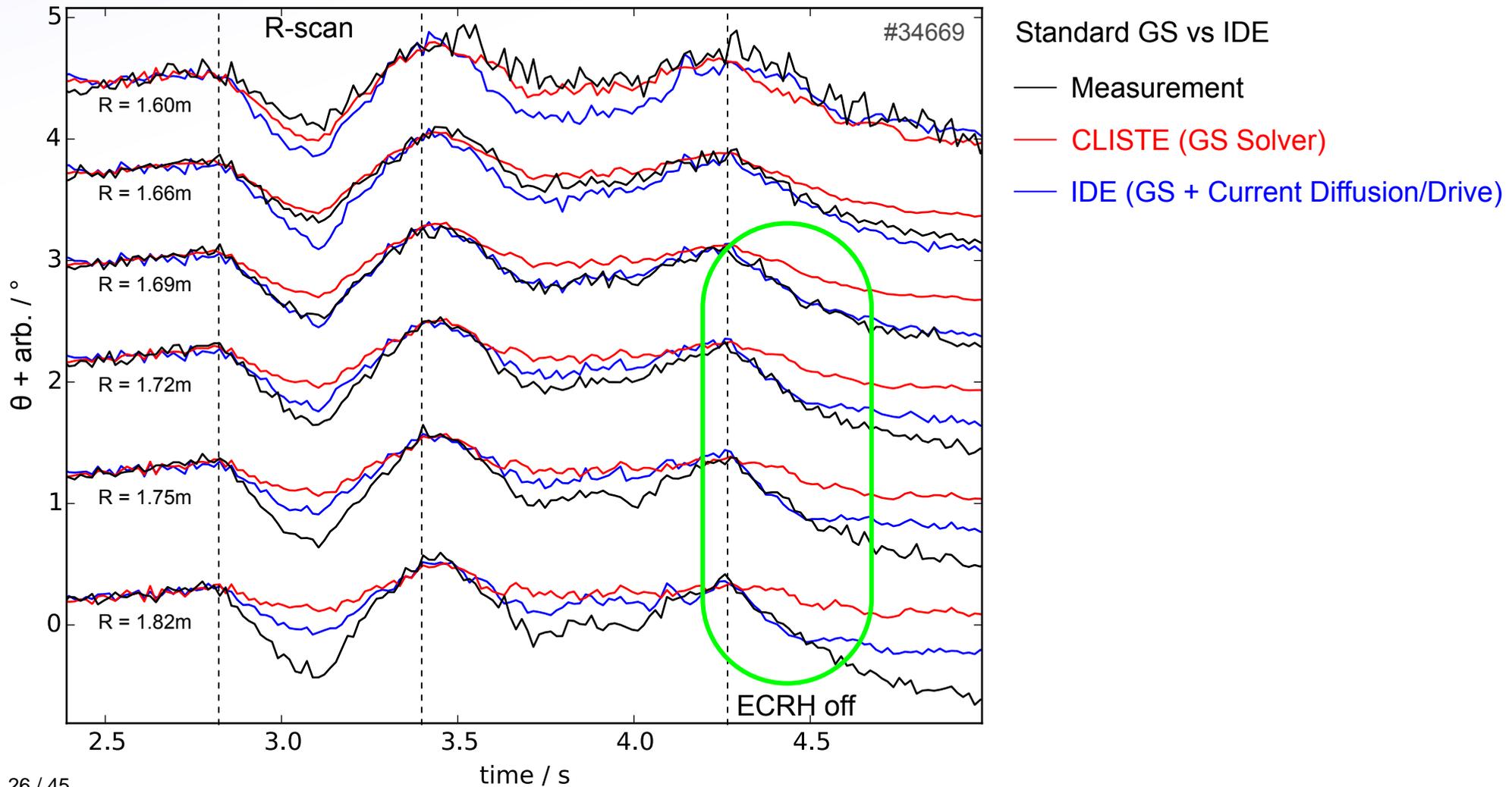
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- 1) During R-scan
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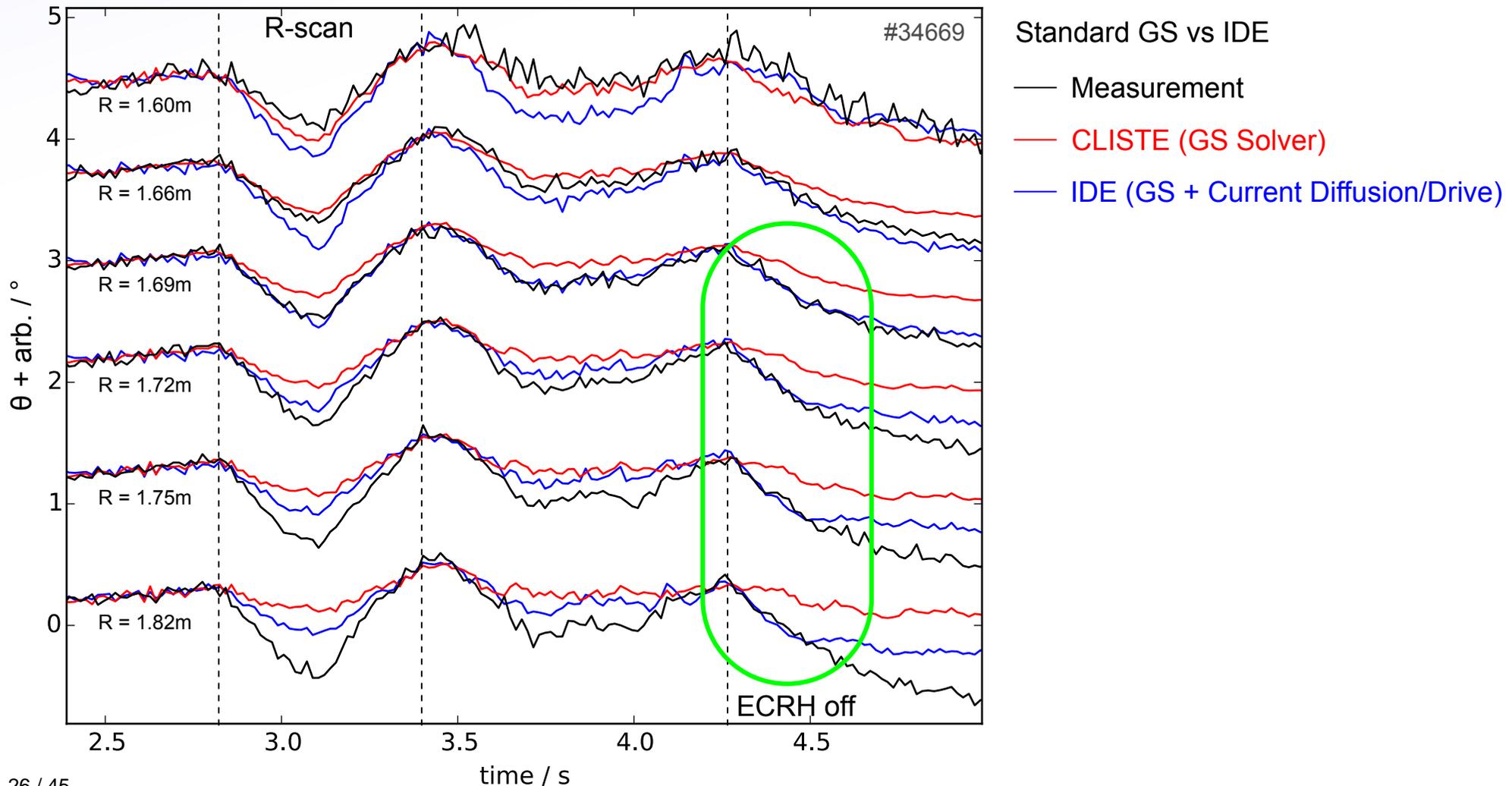


Integrated Equilibrium vs IMSE

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However, there is still physics only seen by diagnostic!



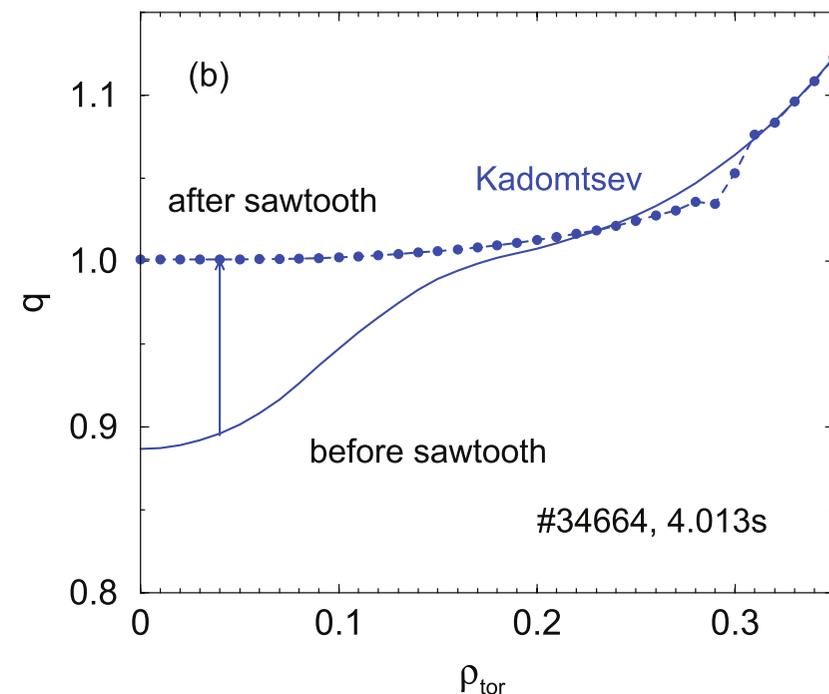
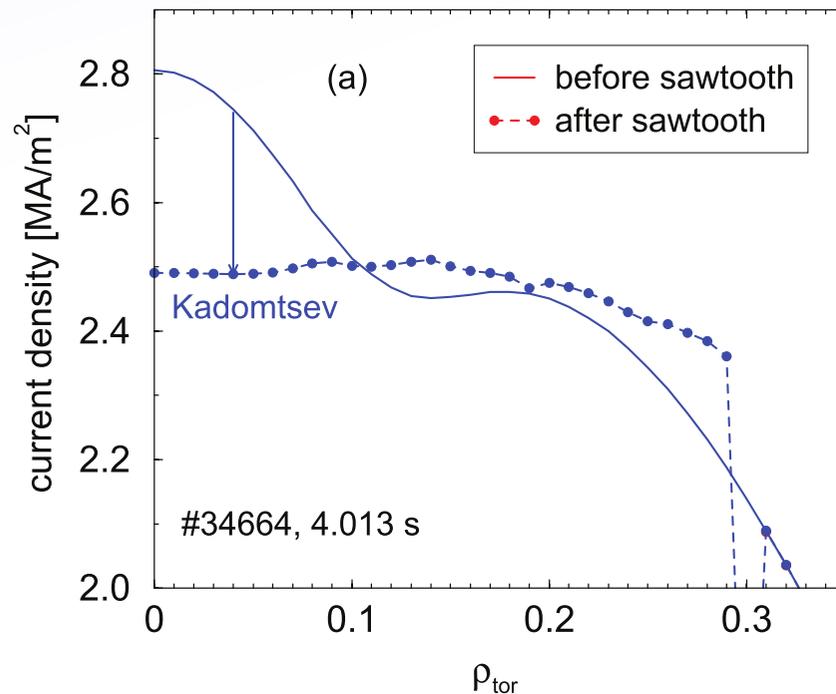


Integrated Equilibrium - Sawteeth

- During sawteeth (reconnection), current diffusion not applicable.

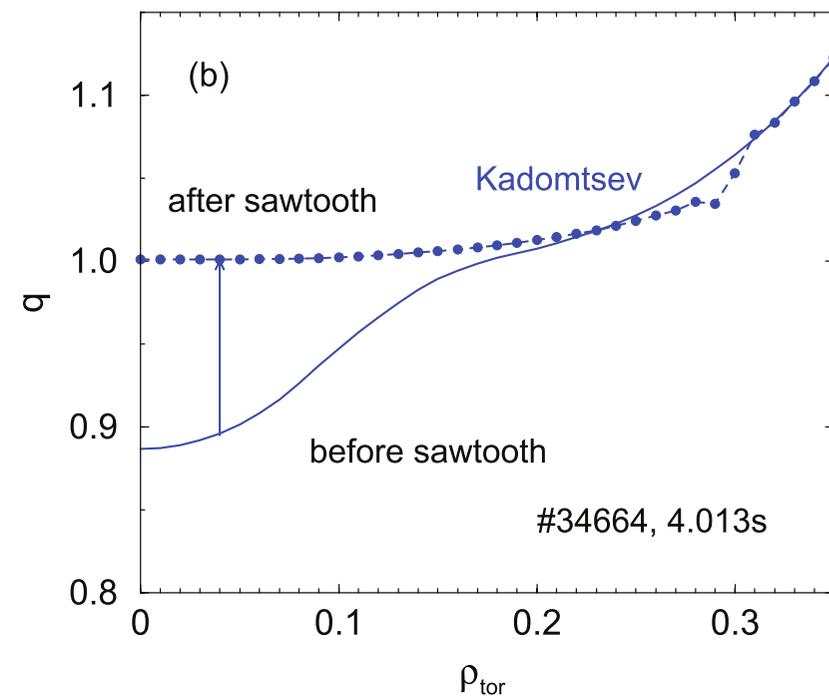
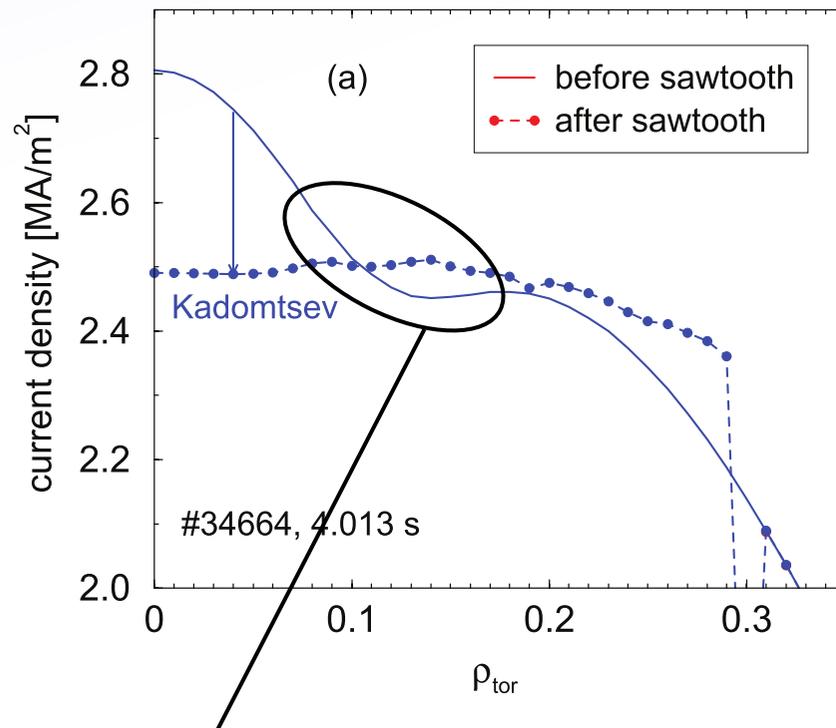
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- Include different sawtooth models in equilibrium code and compare IMSE predictions to measurements.
 - Kadomtsev: Complete reconnection. $q_0 \rightarrow 1$. Current outside $q=1$ surface.
 - Flat-current model (FCM): Current conserved outside $q=1$, flat current density inside.



Integrated Equilibrium - Sawteeth

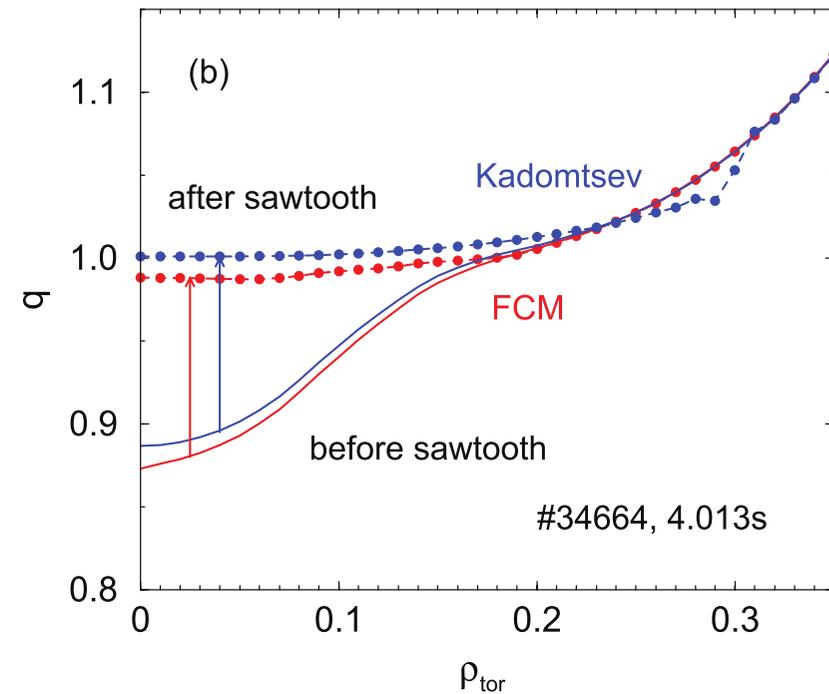
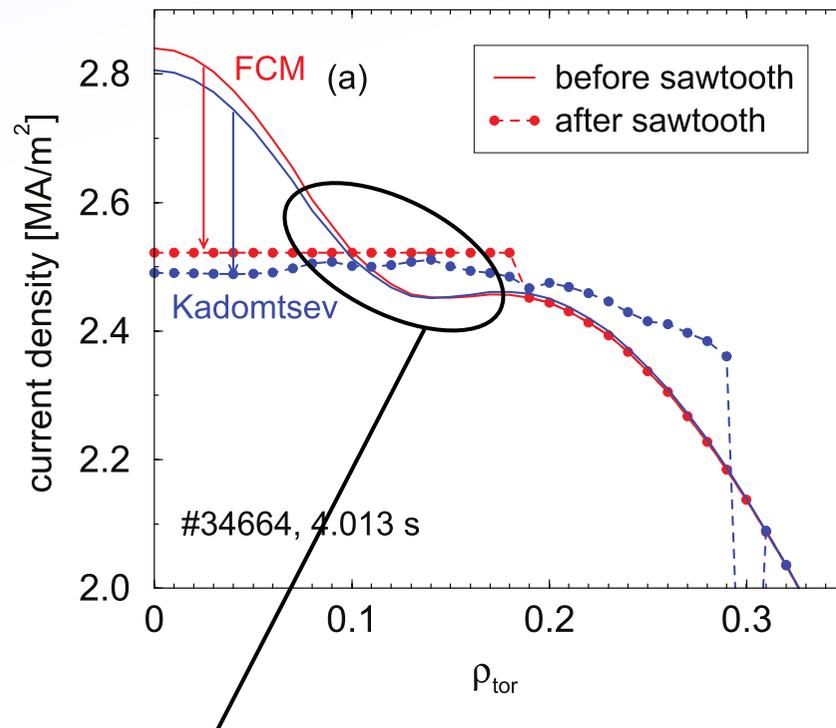
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Current redistribution similar to seen in Δj_ϕ images.

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Current redistribution similar to seen in Δj_ϕ images.

Difference between models requires absolute $j_\phi \sim 0.02 \text{ MA m}^{-2} \rightarrow d\theta/dR \sim 0.01^\circ (3 \text{ cm}^{-3})$



Integrated Equilibrium vs IMSE - Sawteeth

Required precision is so high, many other factors become important:

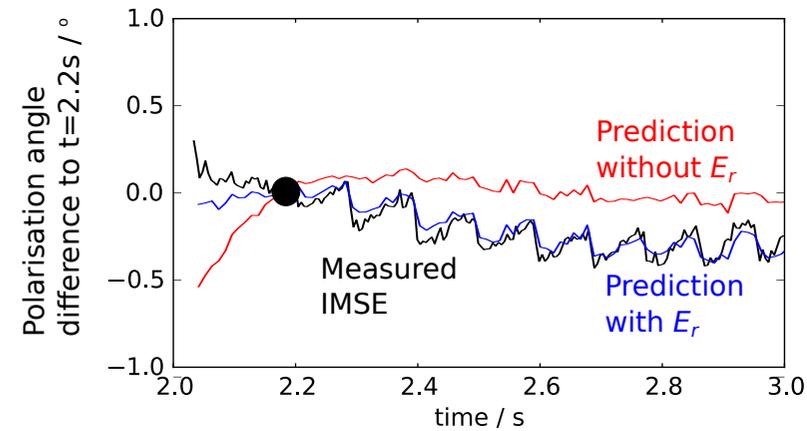
Integrated Equilibrium vs IMSE - Sawteeth

Required precision is so high, many other factors become important:

Plasma radial electric field:

$$E = v \times B + E_r$$

At some locations, ΔE_r during sawtooth dominates measurement:



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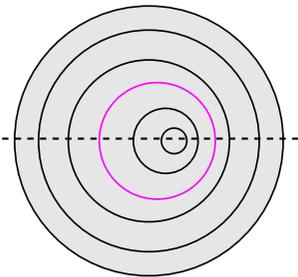
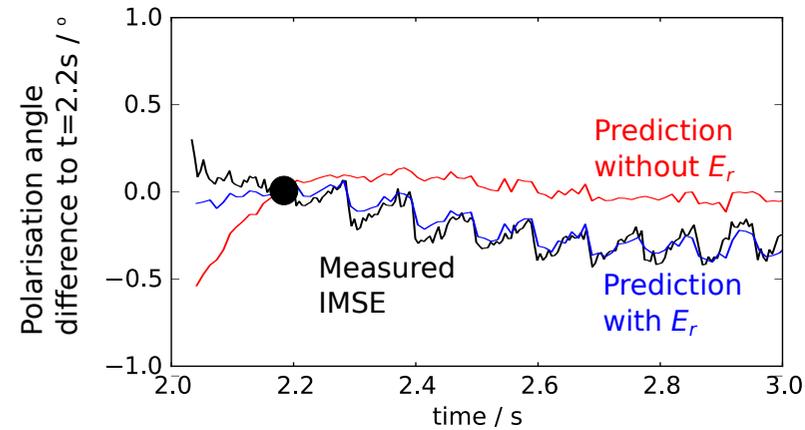
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Movement of plasma axis with pressure.



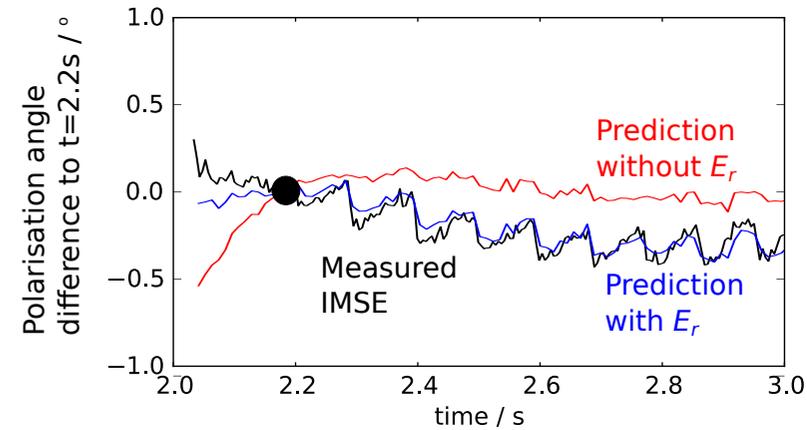
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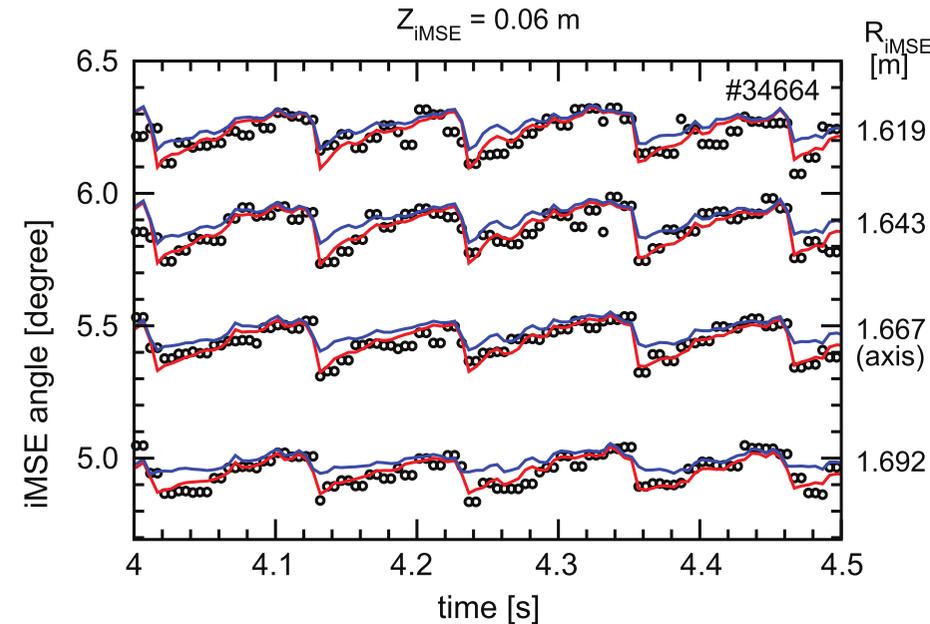
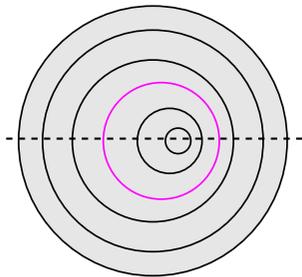
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Movement of plasma axis with pressure.

(including redistribution of fast-ions from neutral beam)

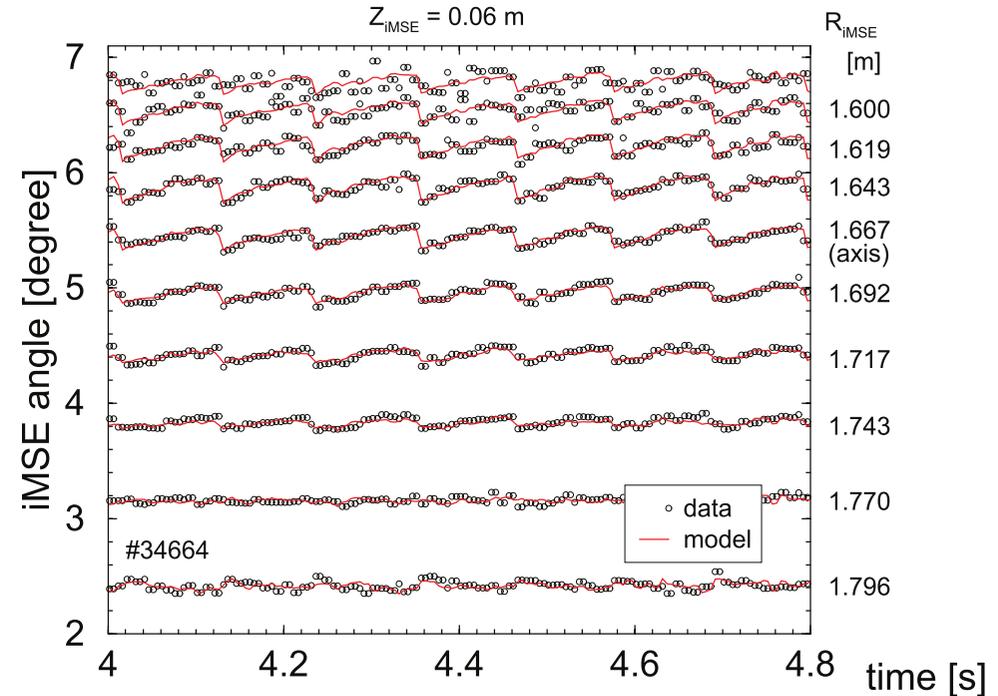
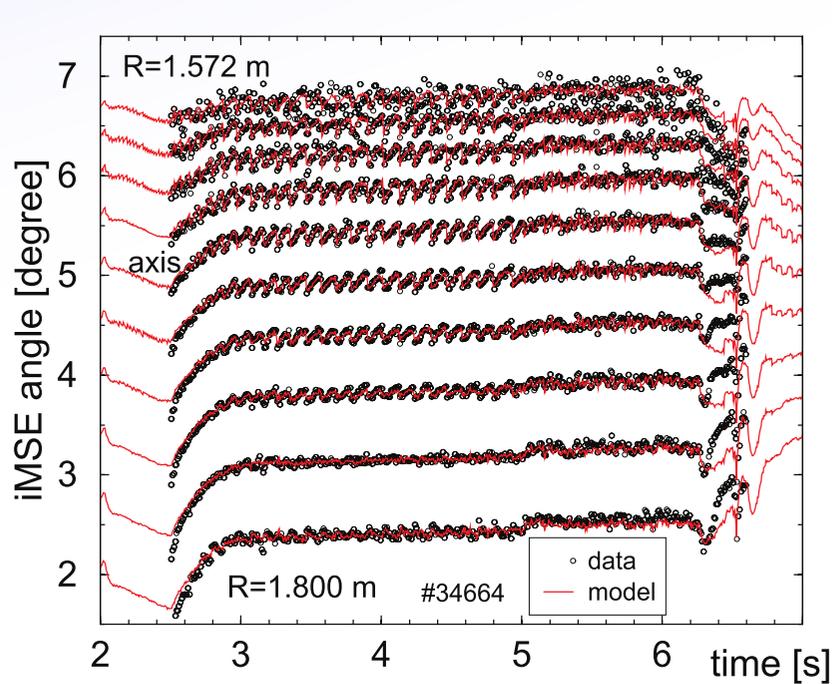


Integrated Equilibrium vs IMSE - Sawteeth

Required precision is so high, many other factors become important:

but...

we now have good agreement between full integrated model and IMSE measurements for sawtooth evolution in θ .

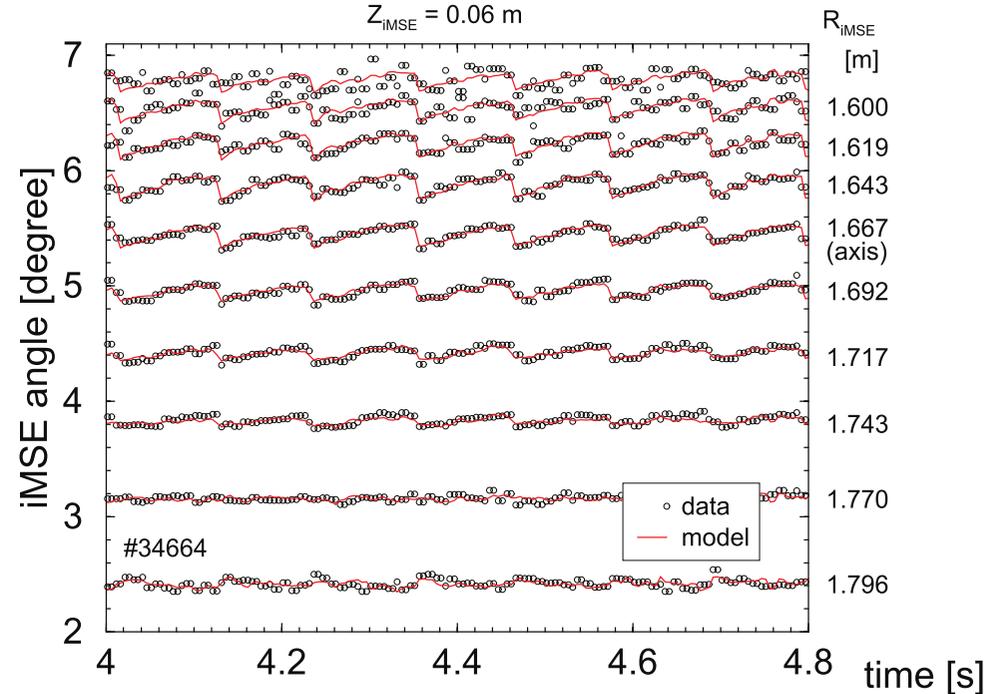
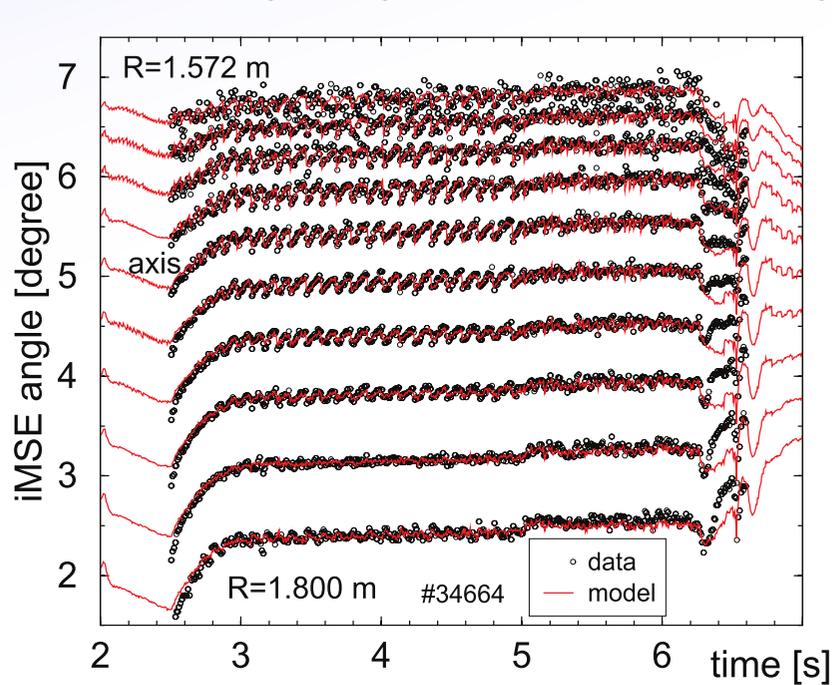


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- This is where we are - 'the state of the art ... science'
What next?

IMSE:
- Improve calibration systematics,.

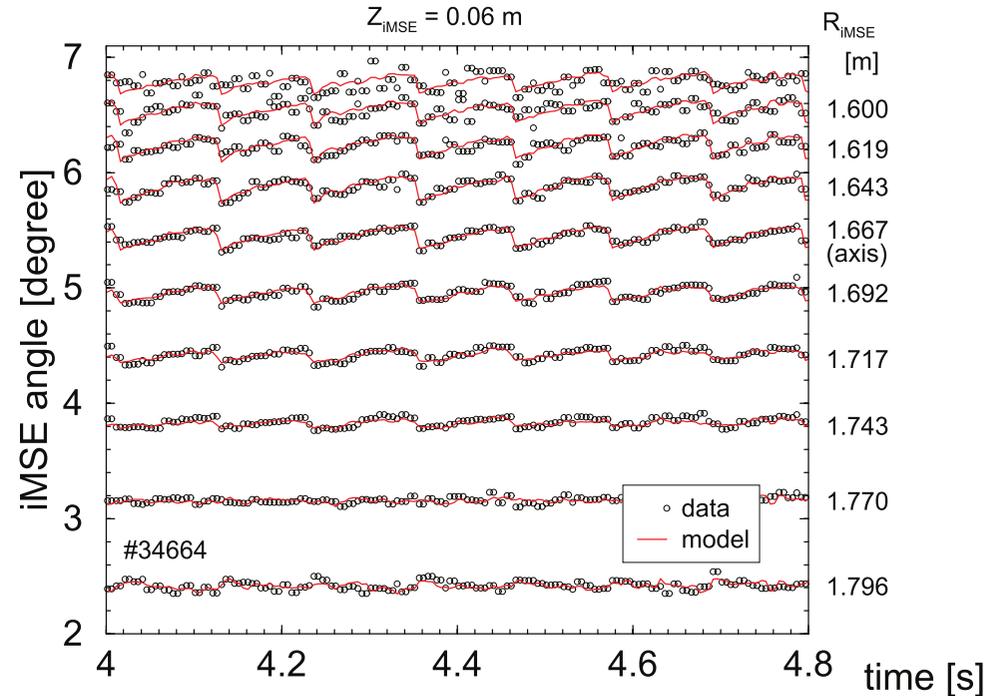
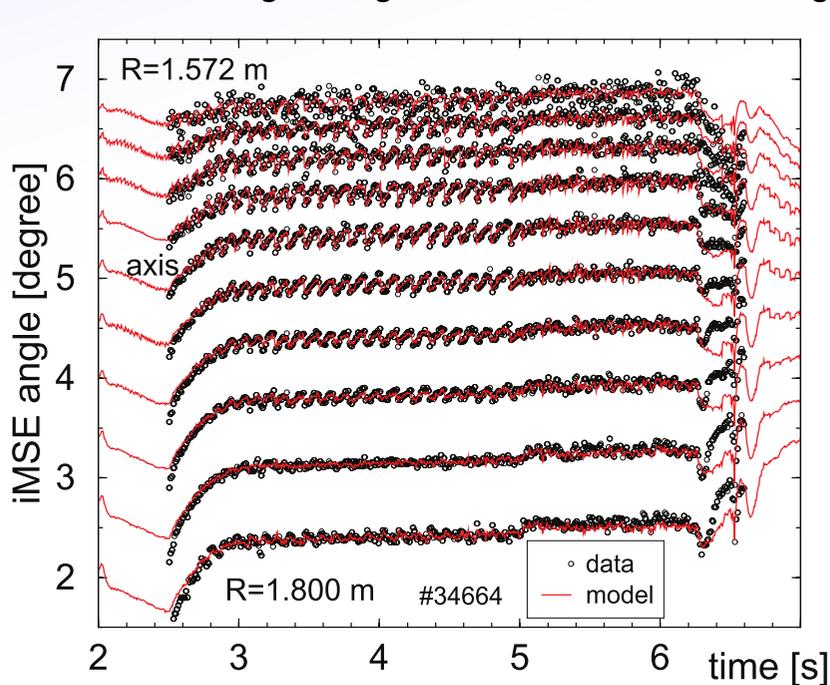
IDE:
- Modeling of effects.
- Tolerance to calibration systematics.

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Converge

IDE:
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Summary

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- Integrated Data Analysis
 - Current diffusion
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- Rigorous determination of uncertainty
 - Too computationally intensive for H-mode
 - Need internal measurements!
- Excellent internal measurements.
 - Good dynamics from approximate derivation of Δj_φ
 - Calibration very difficult to required accuracy.
 - Need to include in equilibrium
- Excellent tool for practical analysis with available data.
 - Current diffusion provides realistic model of missing information when data incomplete.
 - Sawtooth models in good agreement with IMSE evolution.
 - Still need to converge IDE+IMSE to arrive at an absolute q .

Bayesian Inference

A simple example with electron density:

Physics model:

$$n_e(r) = n_0(1 - r^2)^q$$

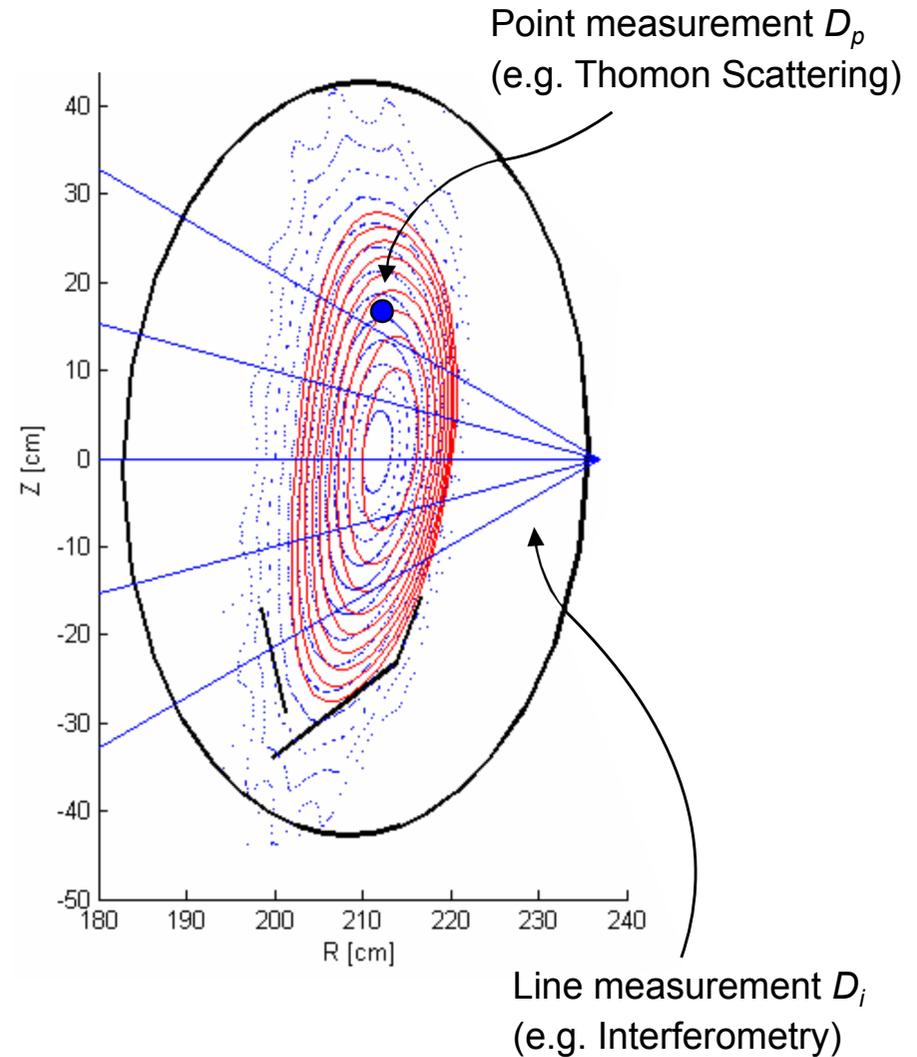
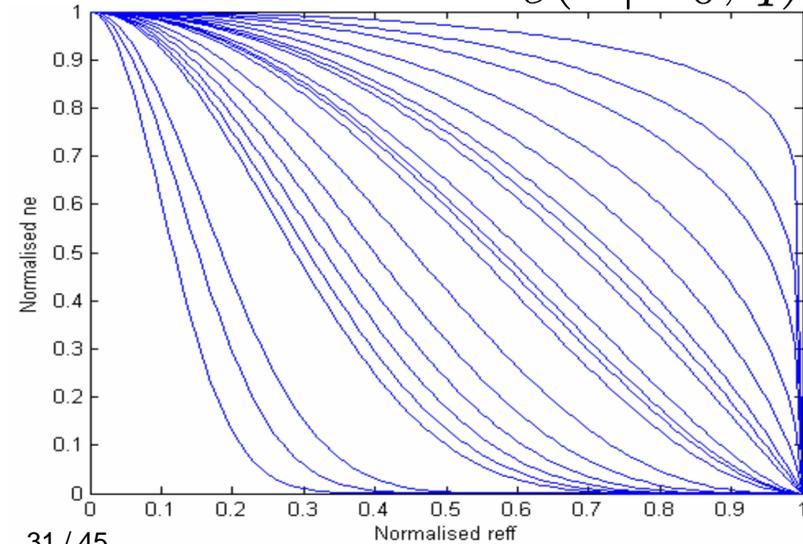
Parameters:

$$\mu = (n_0, q)$$

Forward model:

$$f(\mu) = \int n_e dl$$

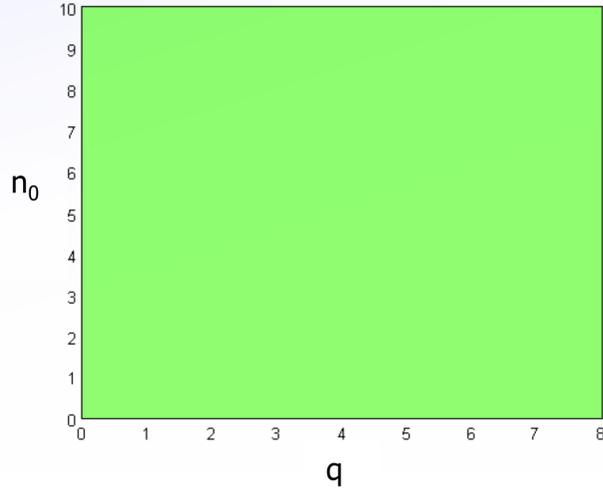
$$n_e(r | n_0, q)$$



Bayesian Inference

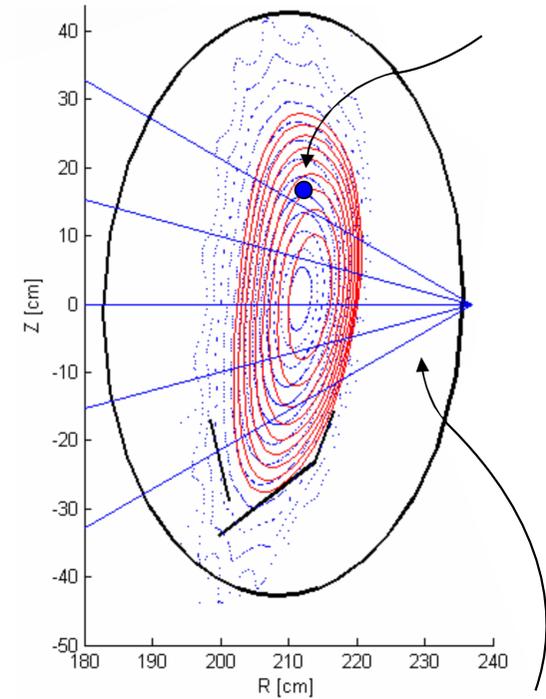
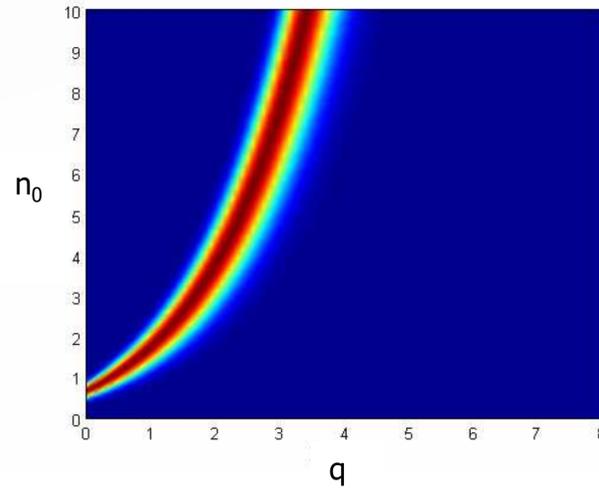
What should we take for a prior:
- Any n_0 and q is equally likely:

Prior

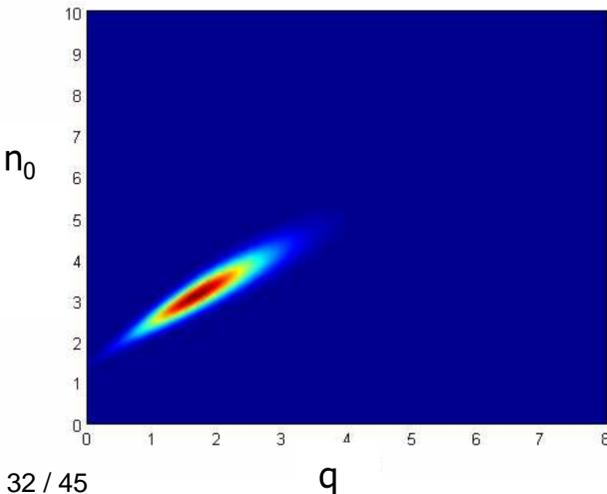


$P(D | \mu)$ for a single lie integral:

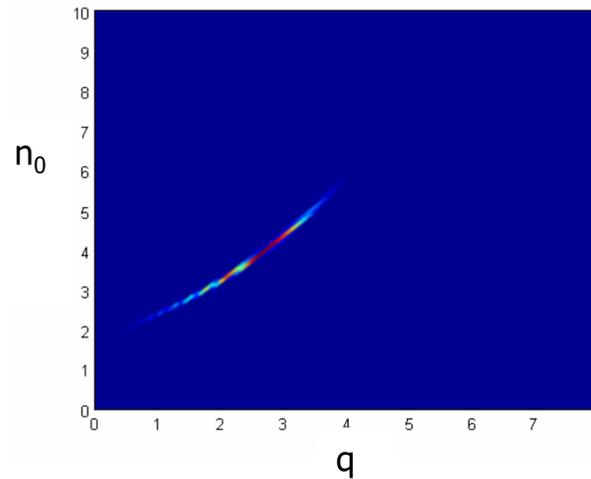
1 channel



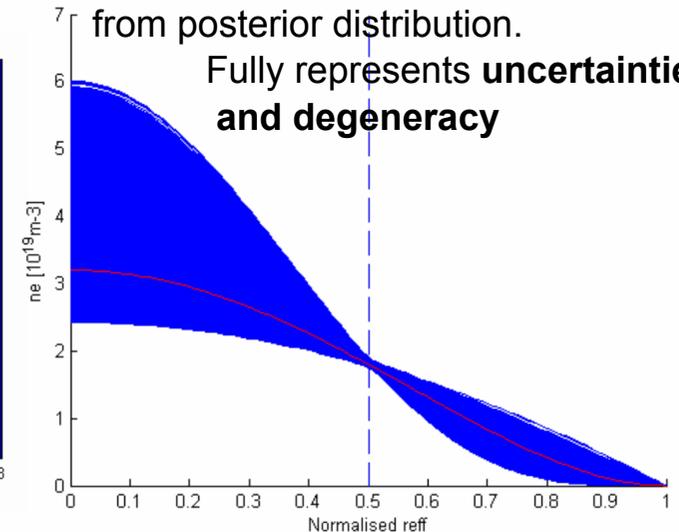
$P(D | \mu)$ for 5 line integrals:



$P(D | \mu)$ for 5 line integrals
and 1 point measurement:

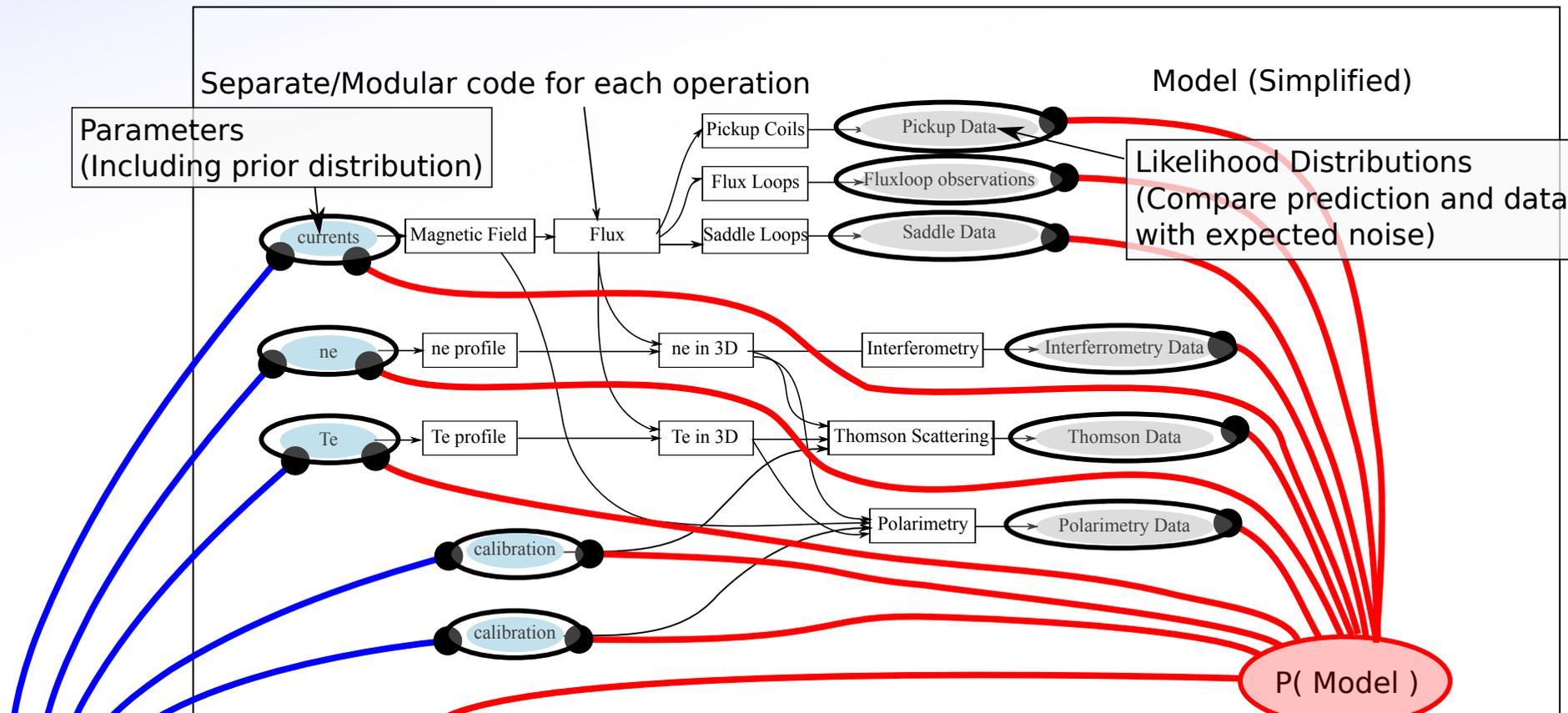


'Samples' = Profiles at points drawn
from posterior distribution.
Fully represents **uncertainties**
and **degeneracy**



Forward modelling and Bayesian Inference

Minerva framework for Bayesian combined modelling:



Bayes Theorem:
$$P(Te, Ne, J | Data) \sim P(D | Ne, Te, J) P(Te, Ne, J)$$

Practically: Solve and explore using external algorithms:

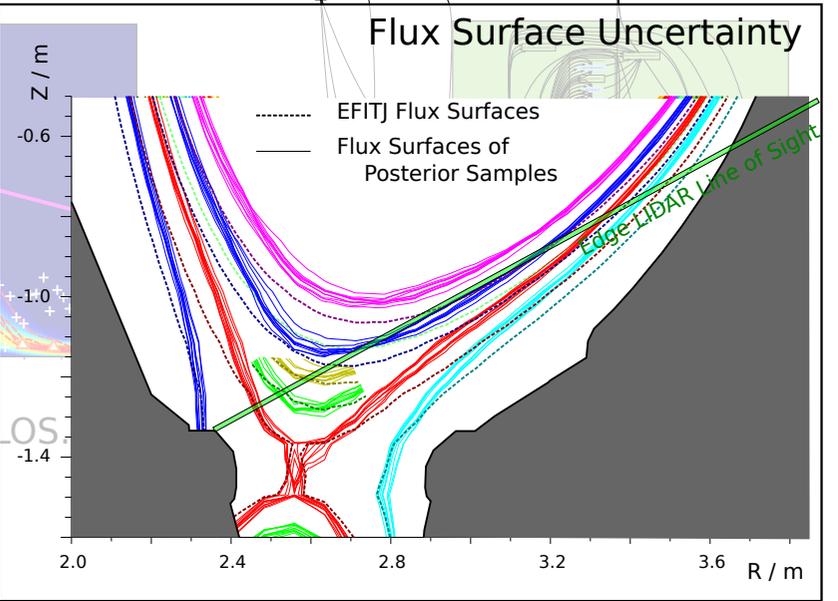
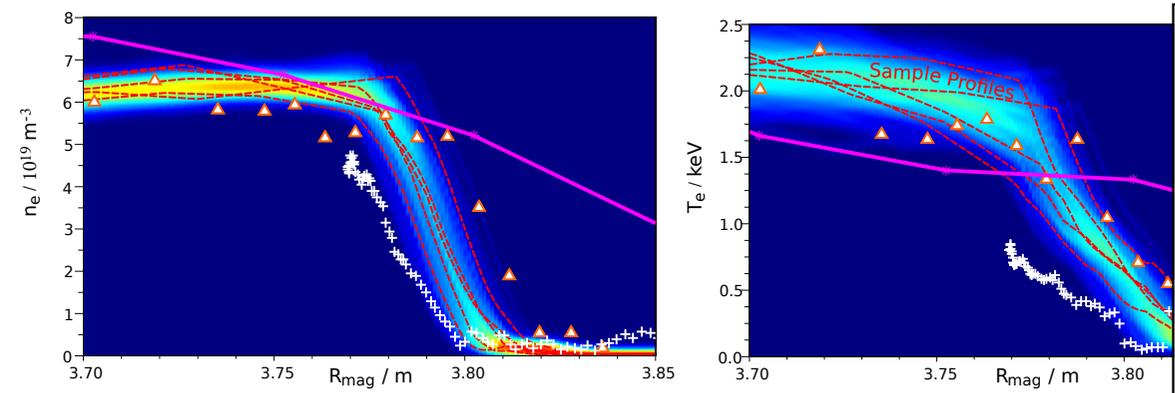
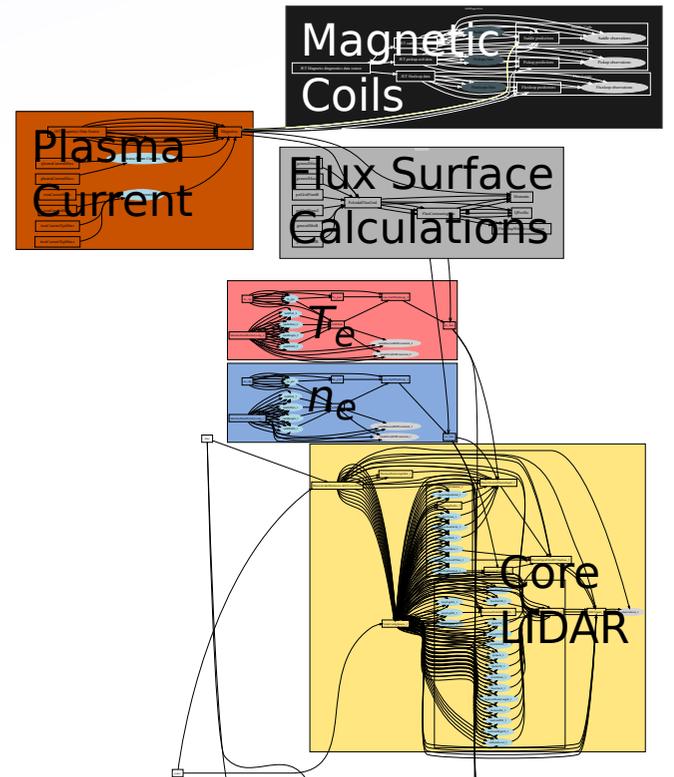
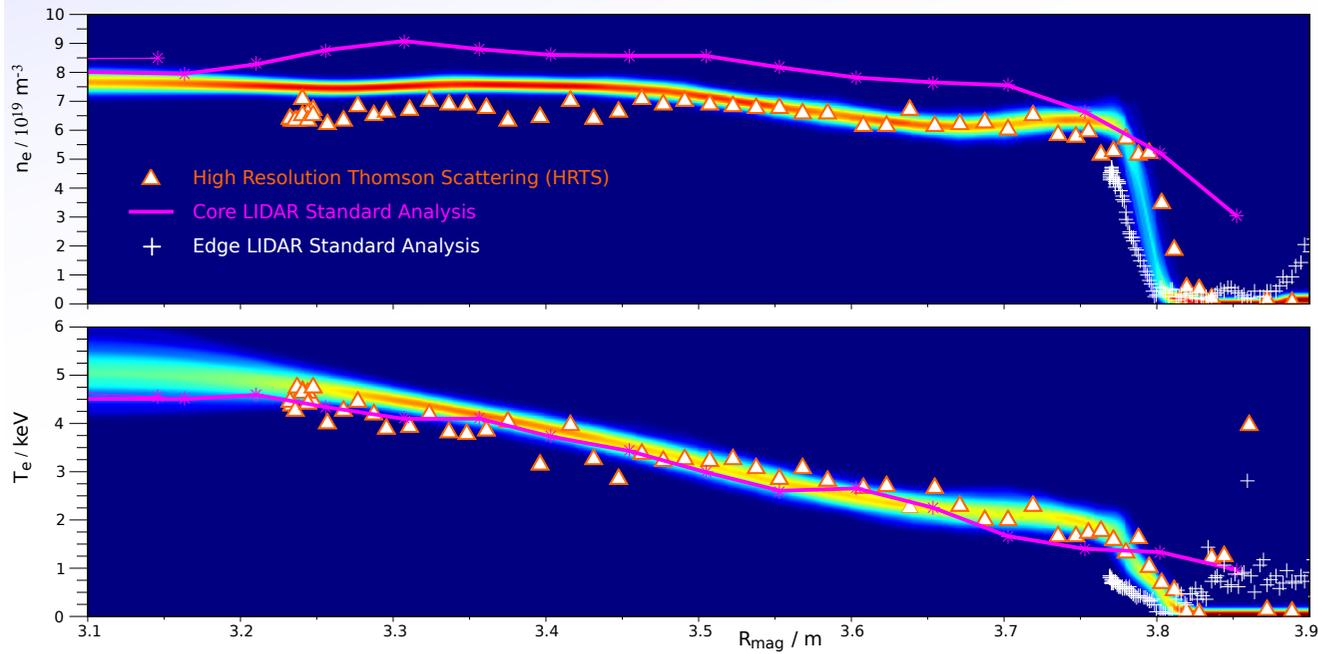
Linear Gaussian Solver
(Best fit and PDF covariance)

Genetic Algorithms
(Non-linear best fit)

Metropolis Hastings
MCMC Non-linear Exploration:
--> Uncertainty



Connect magnetics model and run inversion.

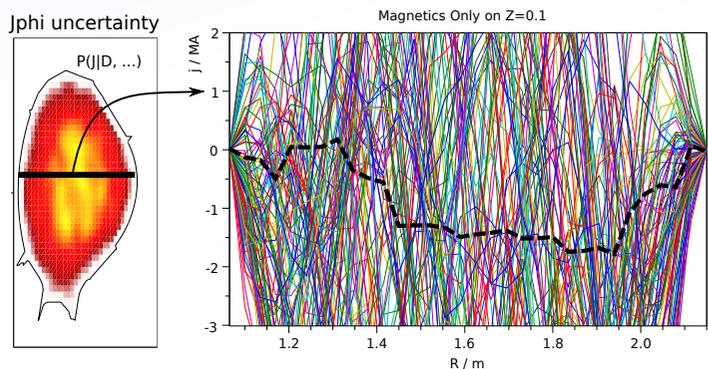


The TS diagnostics provide information on plasma current near LOS
Plasma current one of the most important and least diagnosed
parameters in Tokamaks.

IMSE + Current Tomography

Put description of AUG coils and some pickups into Minerva so we can now do Current Tomography and Bayesian Equilibrium for AUG.

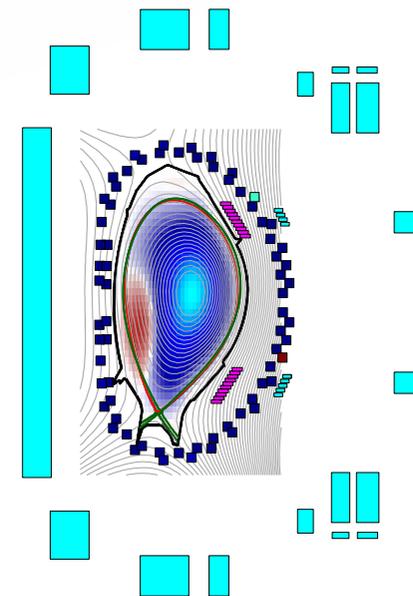
For magnetics only, we have the usual tomography situation:



(Almost) no prior/regularisation



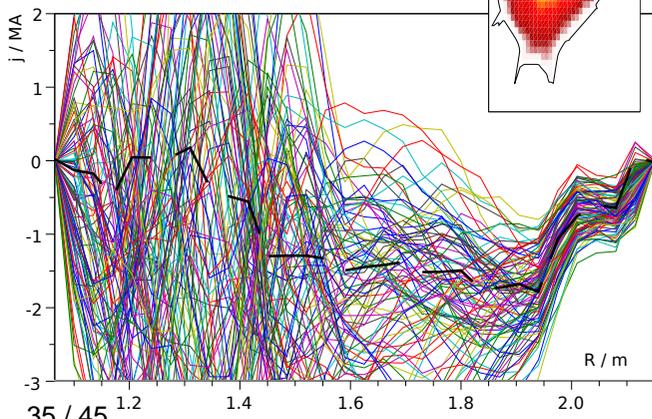
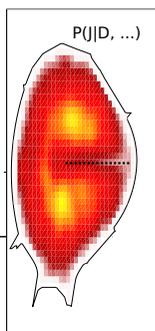
(Almost) infinite uncertainty
(but B/psi still good)



Each case has 900 measurements at $\sigma = 10\text{mT}$.
So difference is only in the **type** of information.

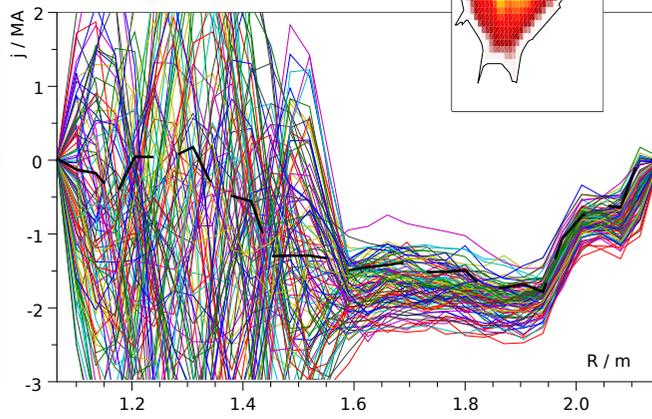
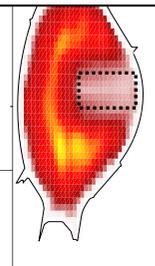
Normal MSE system:

30 x Bz at 30
positions along
NBI centre.



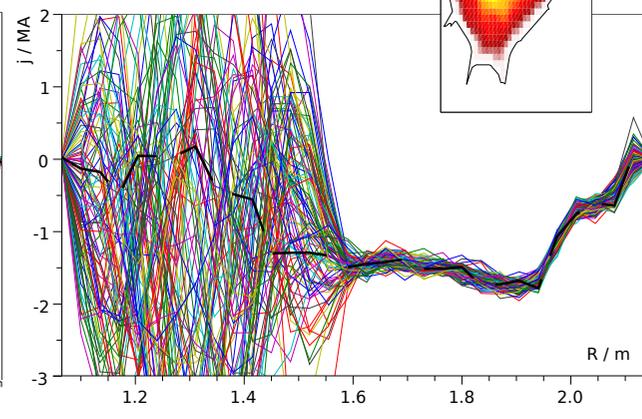
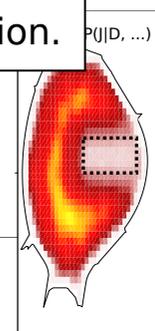
IMSE System:

30x30 grid of Bz
measurements.



Just for interest:

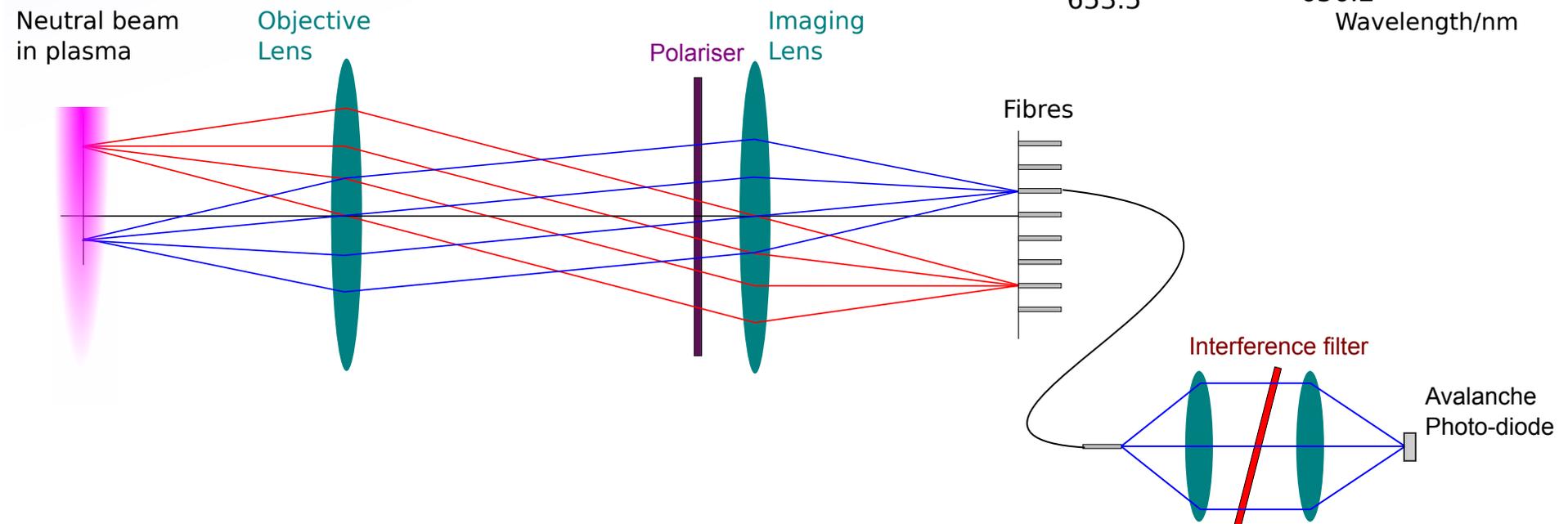
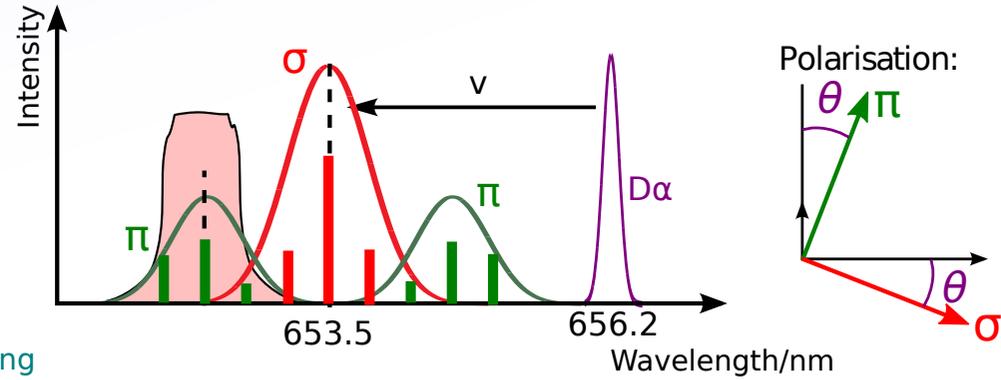
30x15 grid of Bz
30x16 grid of Br.



Motional Stark Effect Polarimetry

Typical hardware:

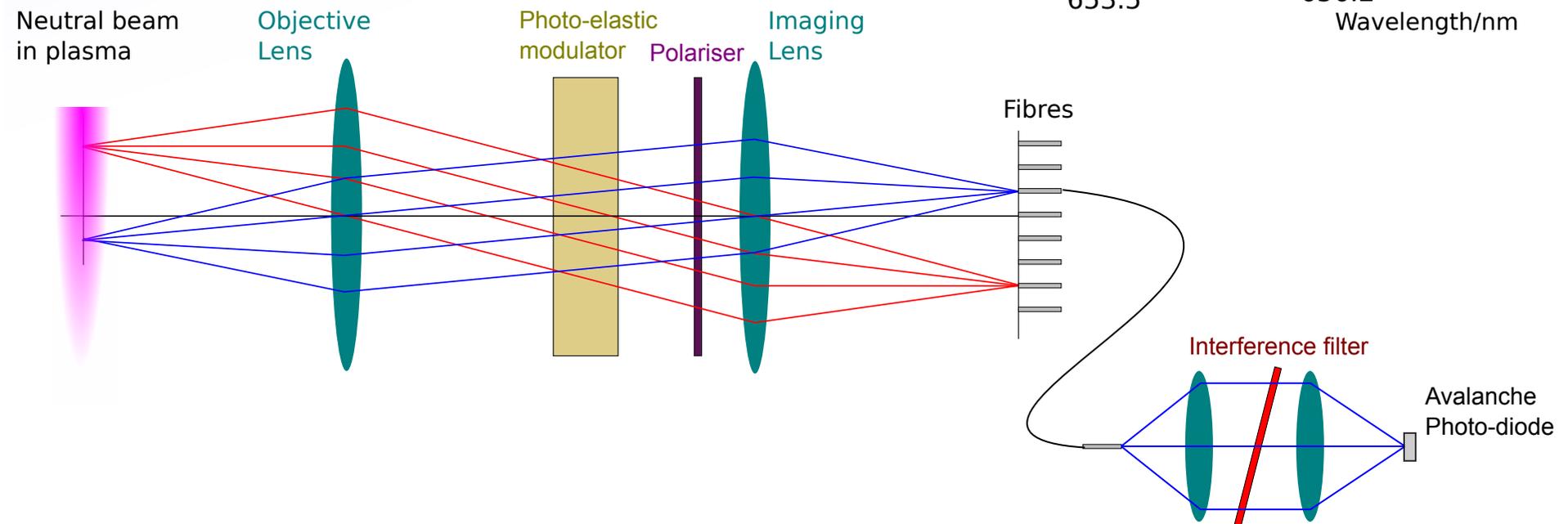
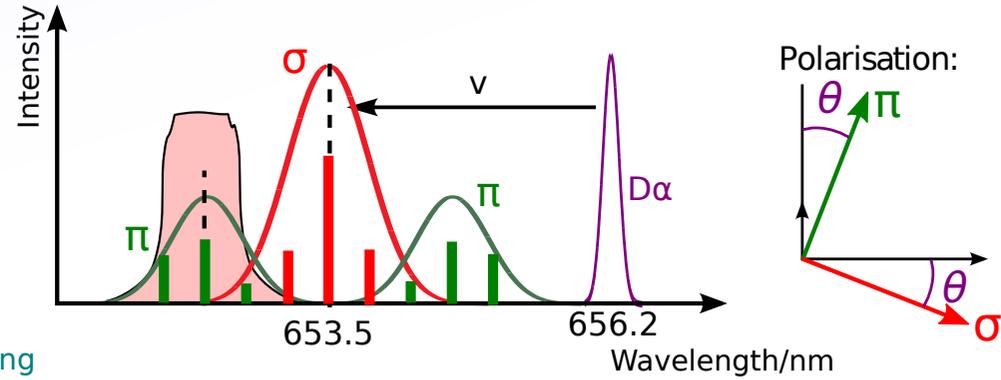
- Temperature/tile tuned interference filter



Motional Stark Effect Polarimetry

Typical hardware:

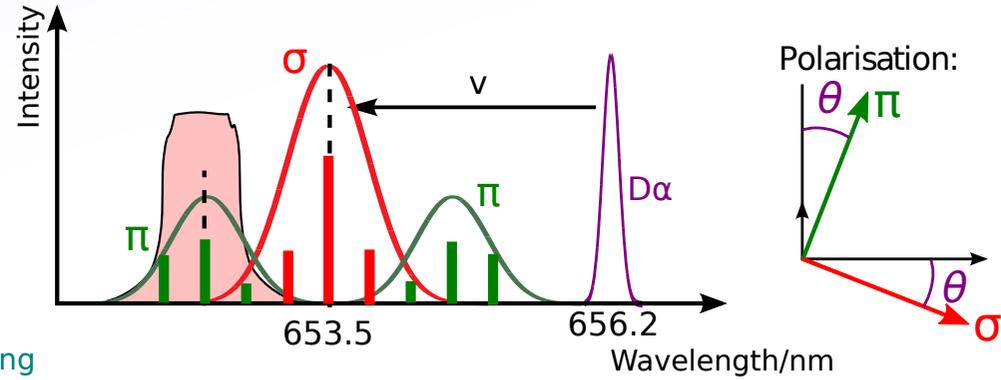
- Temperature/tile tuned interference filter
- Photo-elastic modulator (PEM)



Motional Stark Effect Polarimetry

Typical hardware:

- Temperature/tile tuned interference filter
- Photo-elastic modulator (PEM)



Neutral beam
in plasma

Objective
Lens

Photo-elastic
modulator

Polariser

Imaging
Lens

Fibres

Interference filter

Avalanche
Photo-diode

Initial
polarisation

Ordinary

Phase
delay

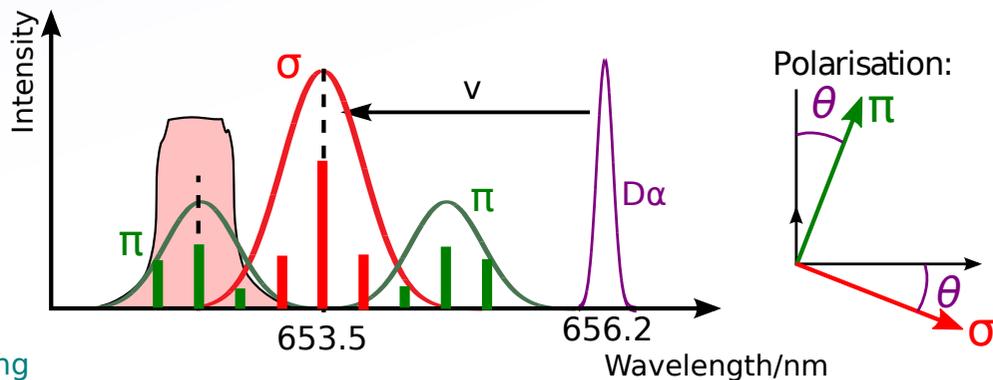
Extraordinary

- PEMS at $\sim 45^\circ$ to initial polarisation
- Splits incoming polarisation into E and O waves.
- Introduces phase shift ϕ between E and O

Motional Stark Effect Polarimetry

Typical hardware:

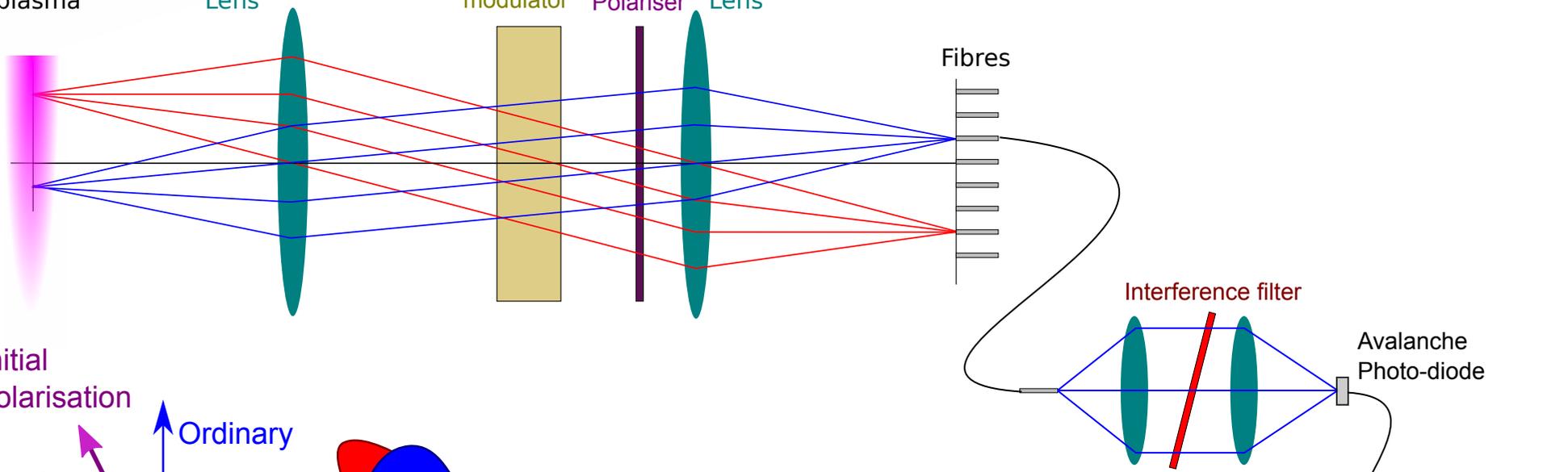
- Temperature/tile tuned interference filter
- Photo-elastic modulator (PEM)



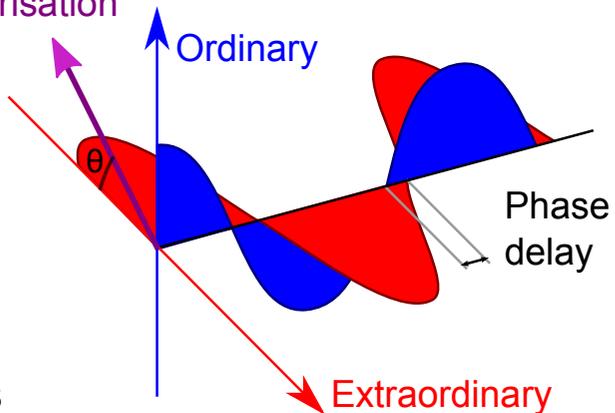
Neutral beam
in plasma

Objective
Lens

Photo-elastic
modulator
Polariser
Imaging
Lens



Initial
polarisation

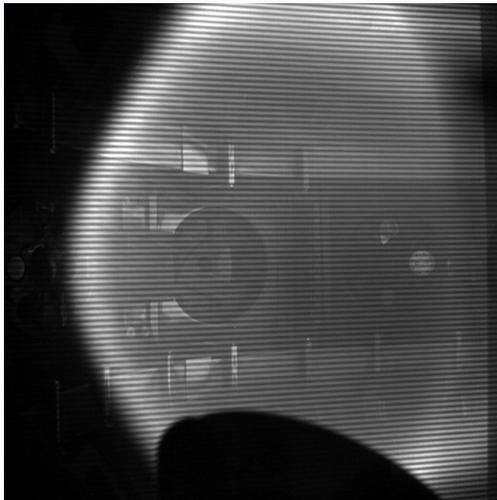


- PEMS at $\sim 45^\circ$ to initial polarisation
- Splits incoming polarisation into E and O waves.
- Introduces phase shift ϕ between E and O
- Modulates $\Delta\phi$ in time
- E and O waves interfere due to final polariser.
- Modulation of $\Delta\phi$ in time from which θ can be recovered.

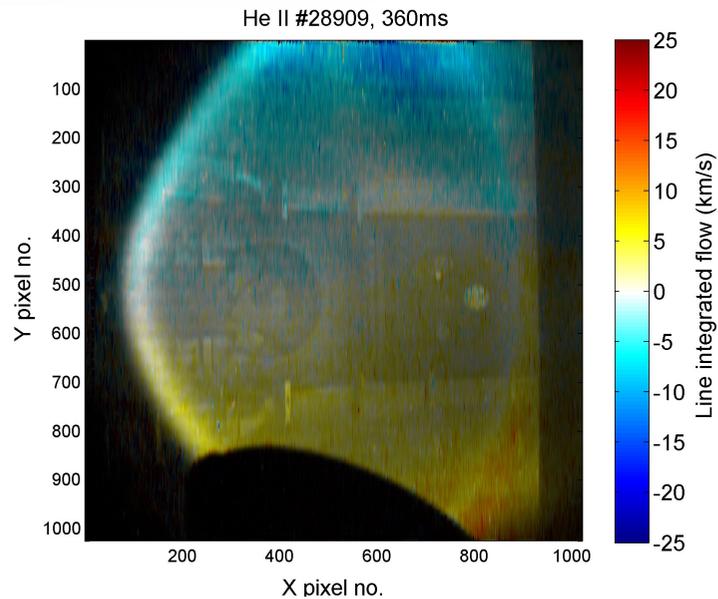
Coherence Imaging

Some results of neutral Helium flow in the (relatively) cold edge of MAST:

Raw Image:

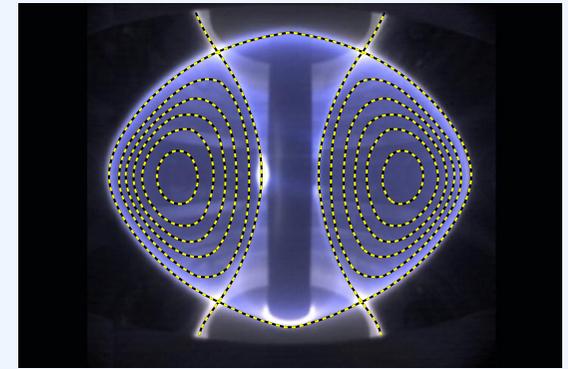


Helium Flow Velocity:



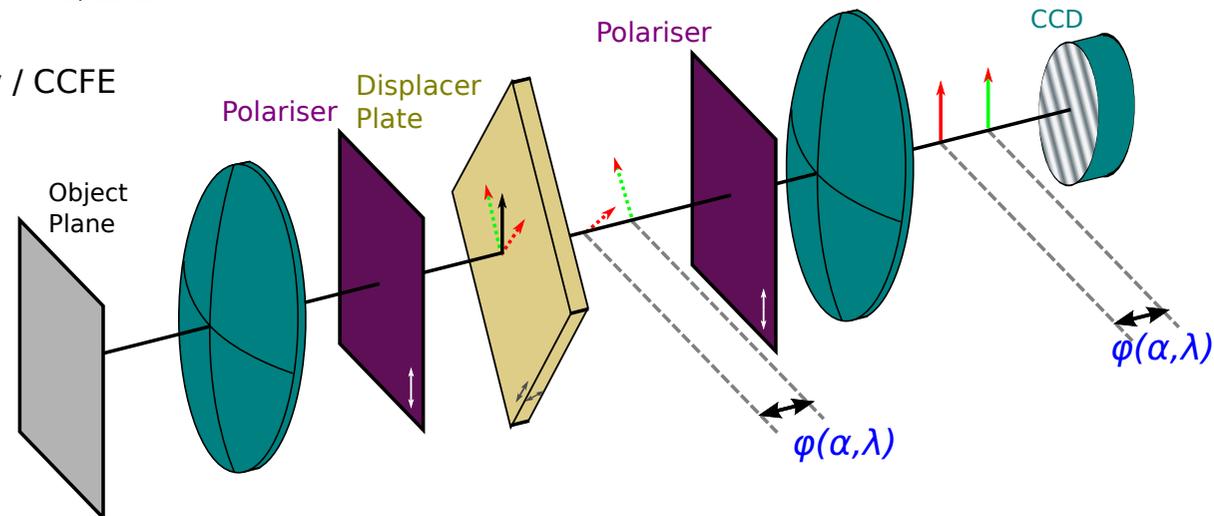
MAST

Mega Amp Spherical Tokamak,
CCFE, Culham, UK



MAST is a 'spherical' Tokamak. The torus has a very small major radius compared to its minor radius, but is still a Tokamak.

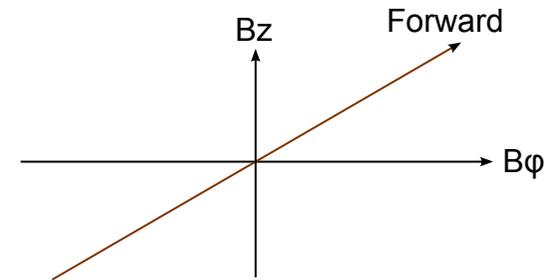
*With thanks to Scott Silburn, Durham University / CCFE
[S. Silburn et. al. 40th EPS Conf. on plasma phys. 2013]



IMSE - Calibration

Absolute q_0 requires absolute $d\theta/dR$
How can we calibrate θ , (or $d\theta/dR$)?

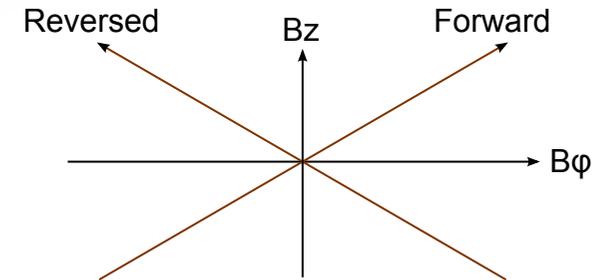
- Run the same plasma with reversed field --> Reversed pitch angle



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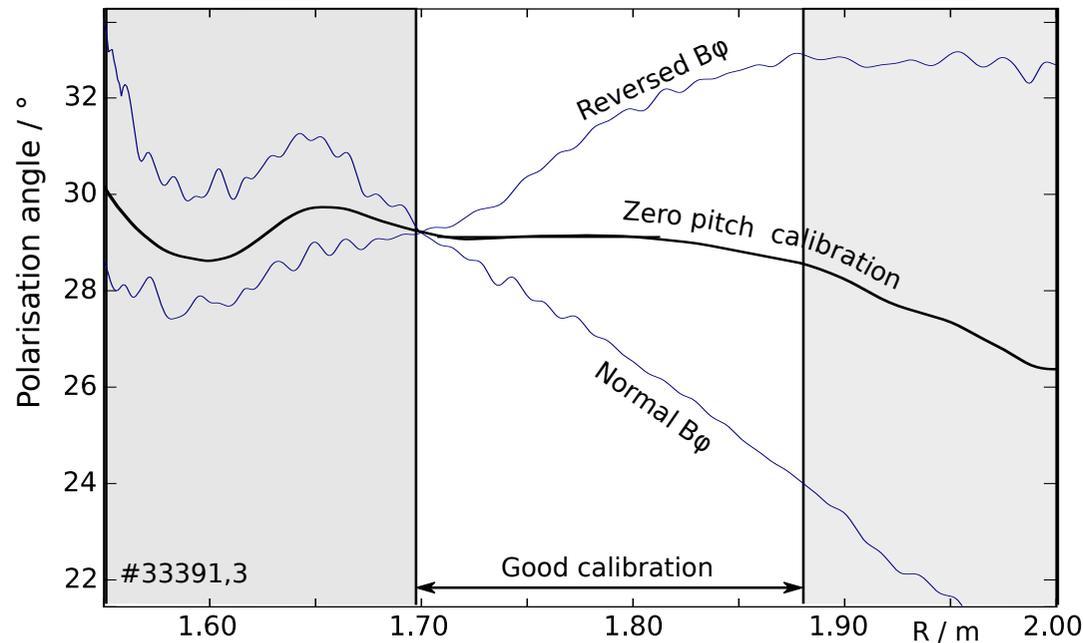
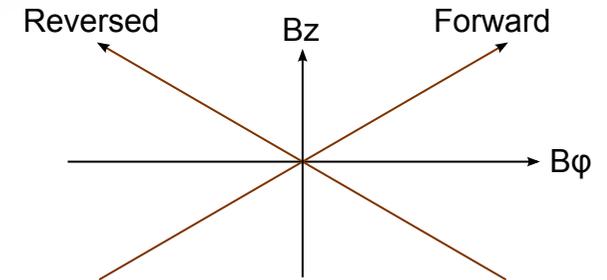
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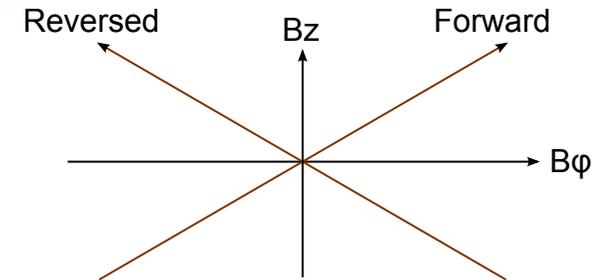
- Run the same plasma with reversed field --> Reversed pitch angle



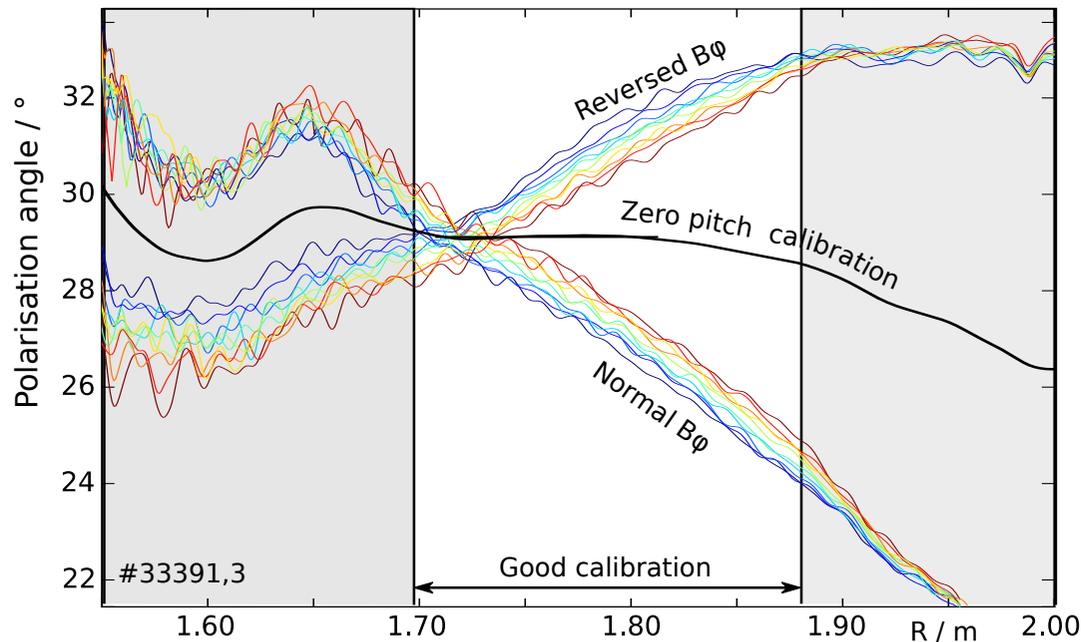
IMSE - Calibration

Absolute q_0 requires absolute $d\theta/dR$
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- Run the same plasma with reversed field --> Reversed pitch angle
- Also scan axis position to confirm meeting point (magnetic axis) agrees with 0 pitch angle.



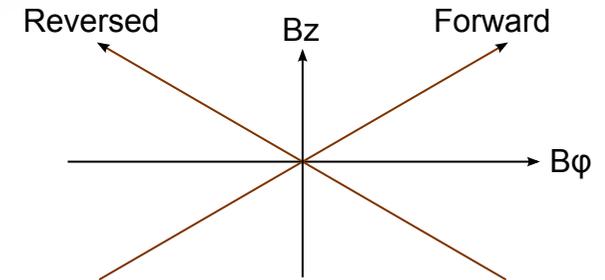
Plasma moved by 5cm
to scan axis position.



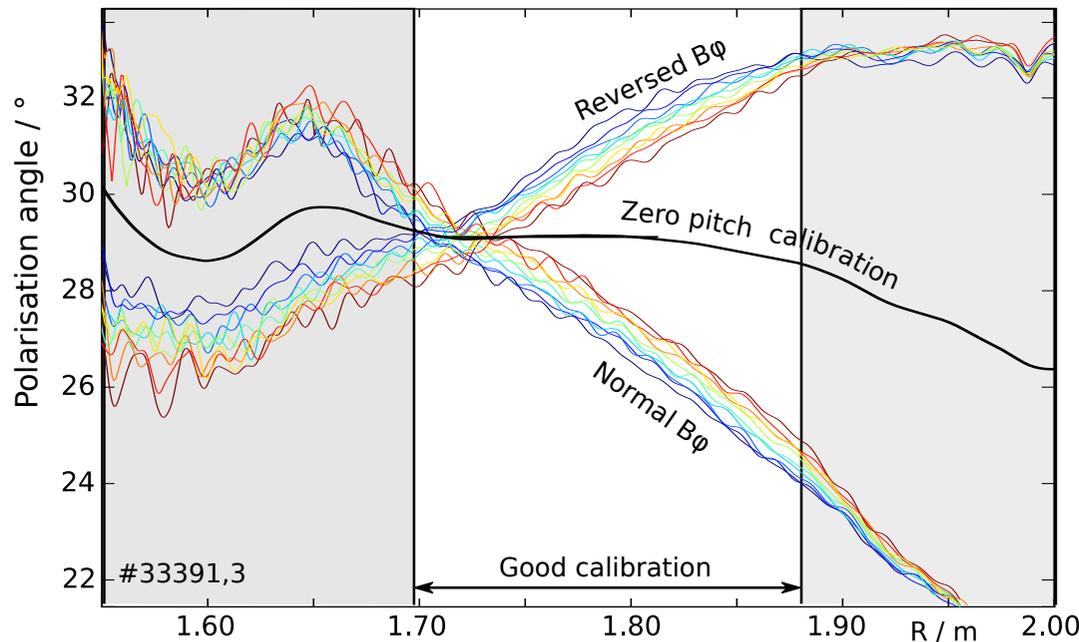
IMSE - Calibration

Absolute q_0 requires absolute $d\theta/dR$
How can we calibrate θ , (or $d\theta/dR$)?

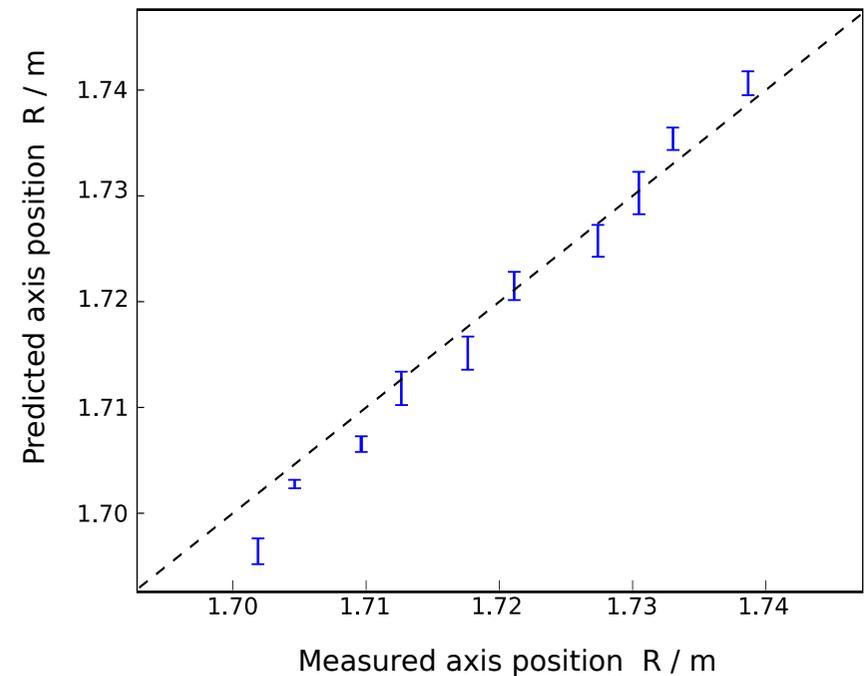
- Run the same plasma with reversed field --> Reversed pitch angle
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Plasma moved by 5cm
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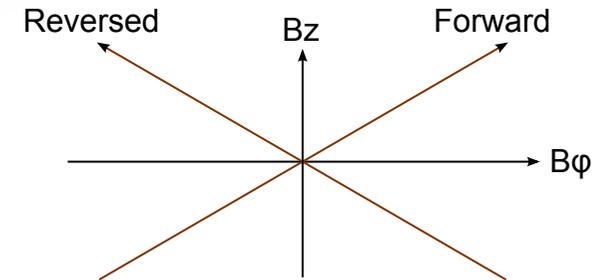
Meeting point well predicted by equilibrium:



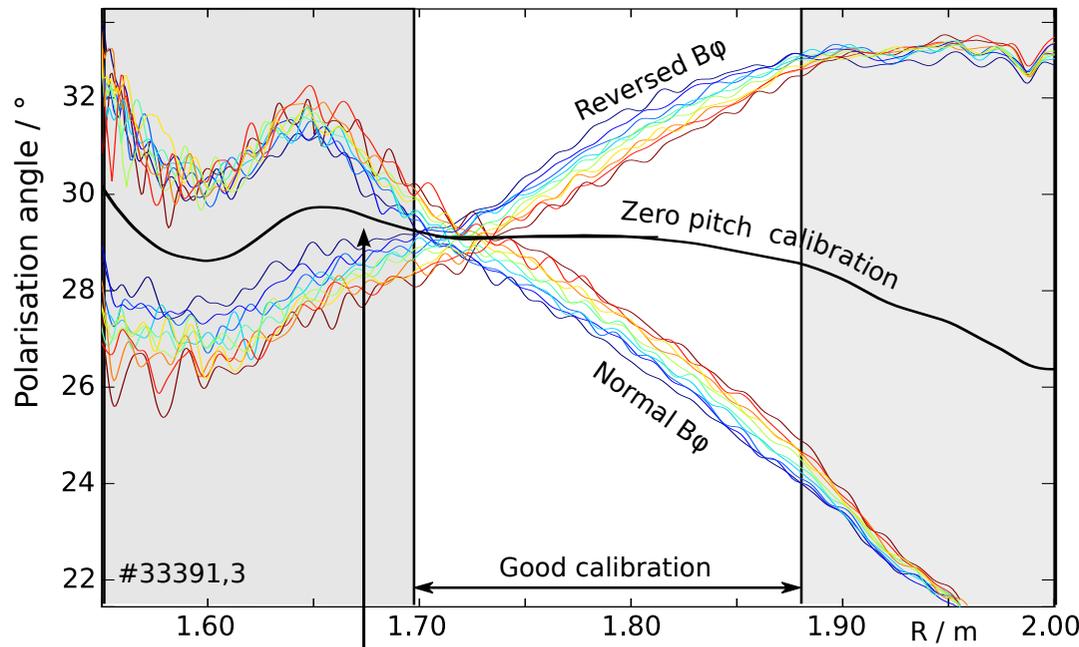
IMSE - Calibration

Absolute q_0 requires absolute $d\theta/dR$
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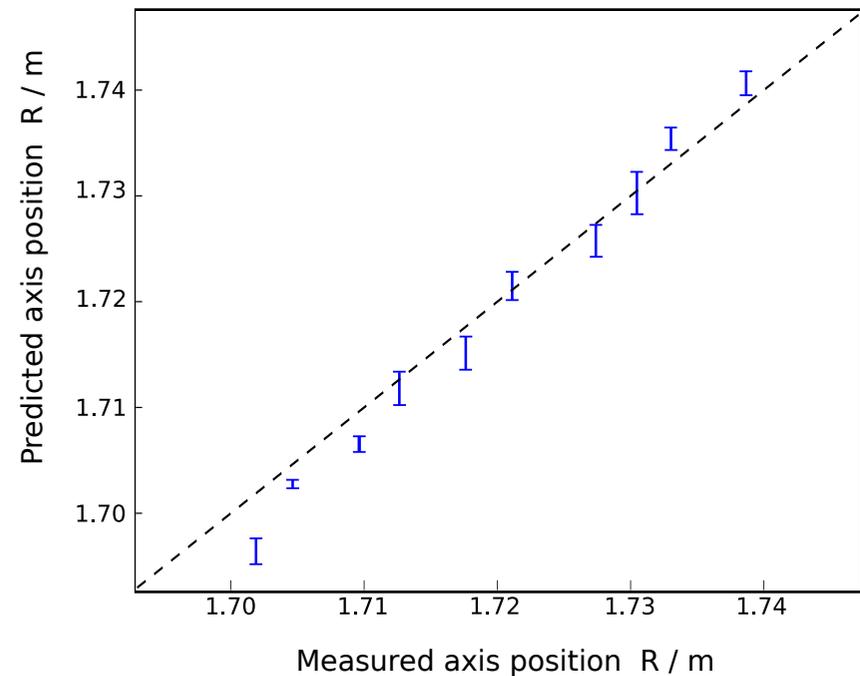


Plasma moved by 5cm
to scan axis position.

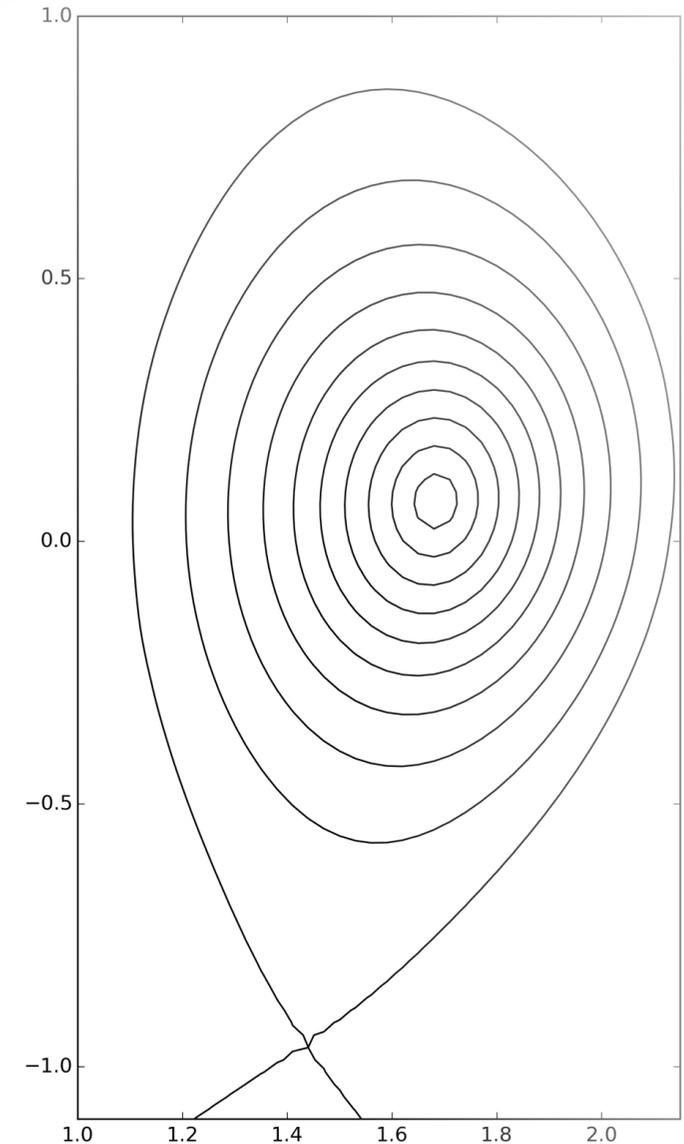
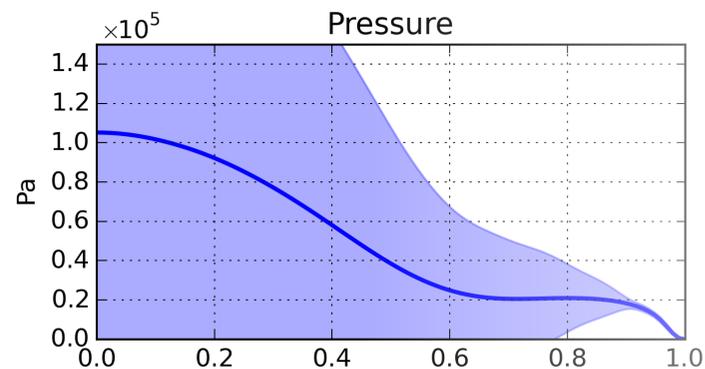
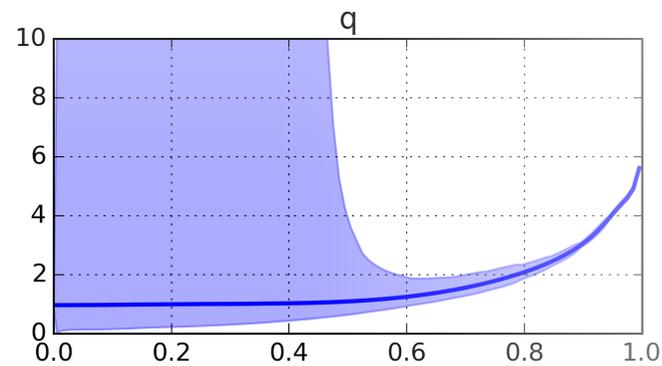
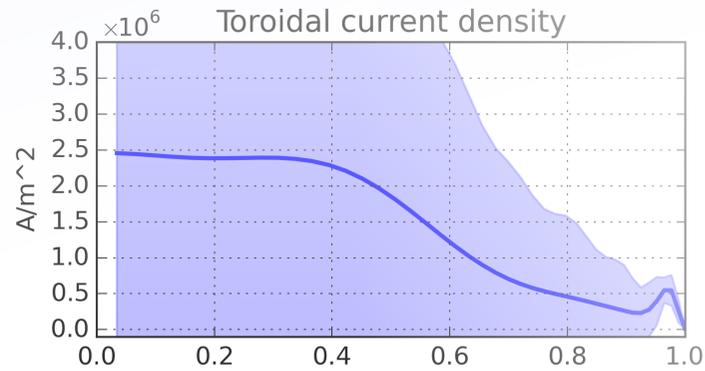


Systematic features $\sim 0.5^\circ / (3 \text{ cm}^{-1})$!!

Meeting point well predicted by equilibrium:

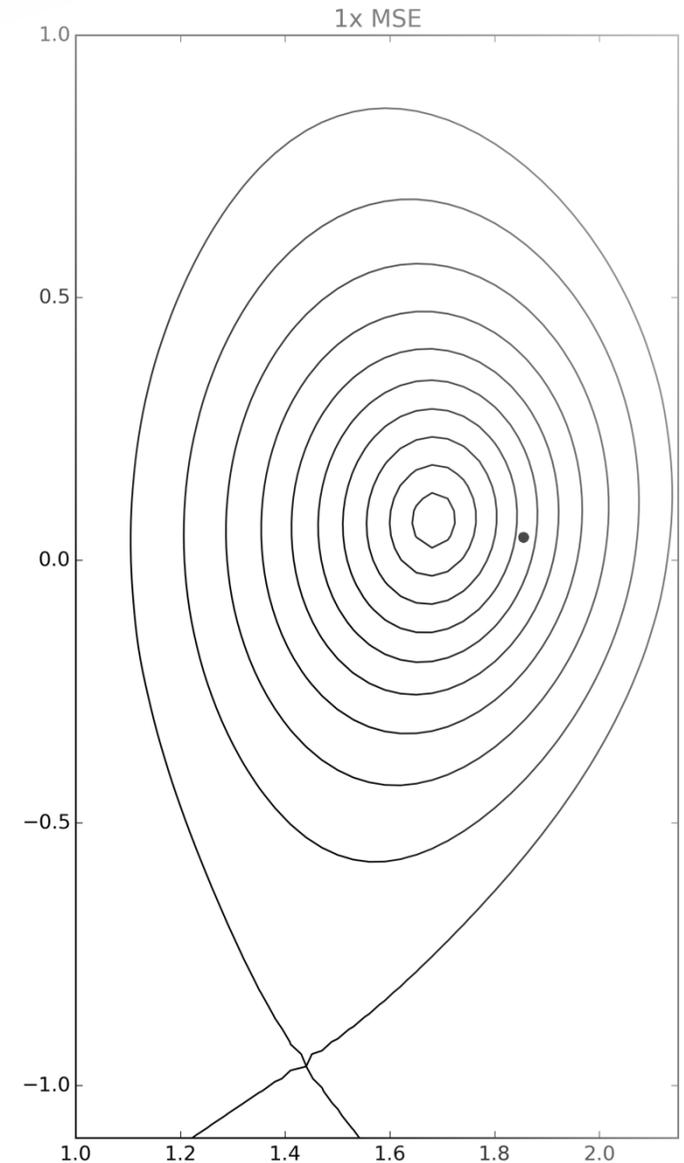
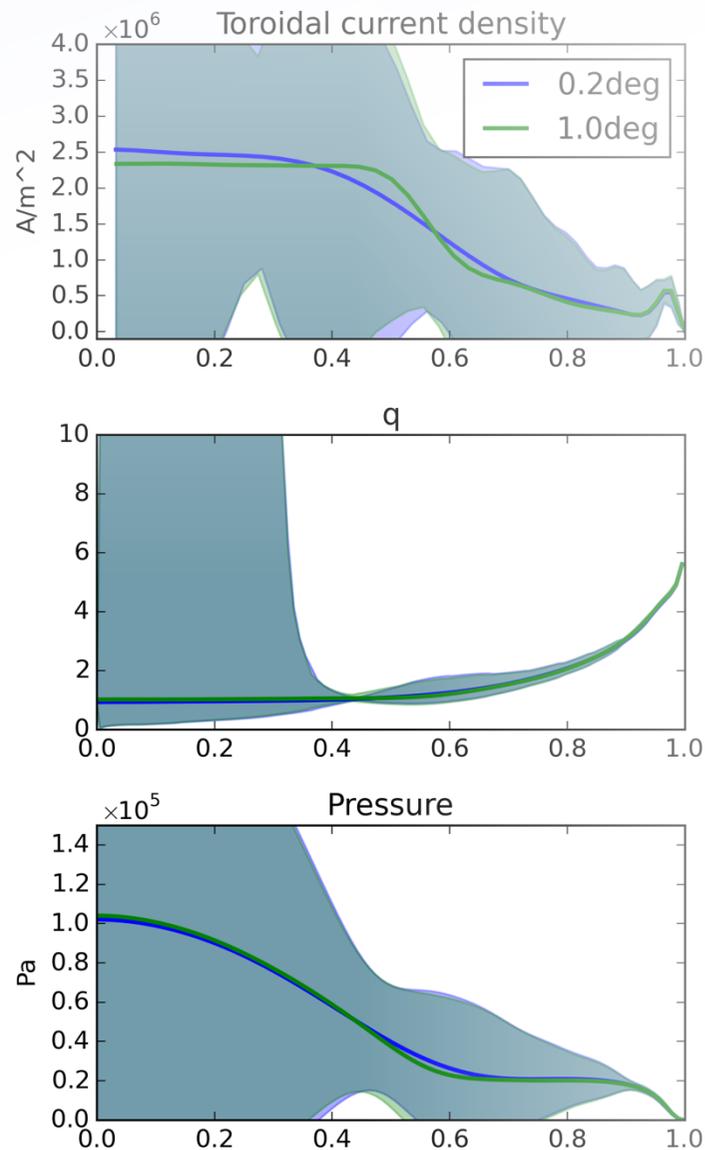


Uncertainties in j/q without internal measurements



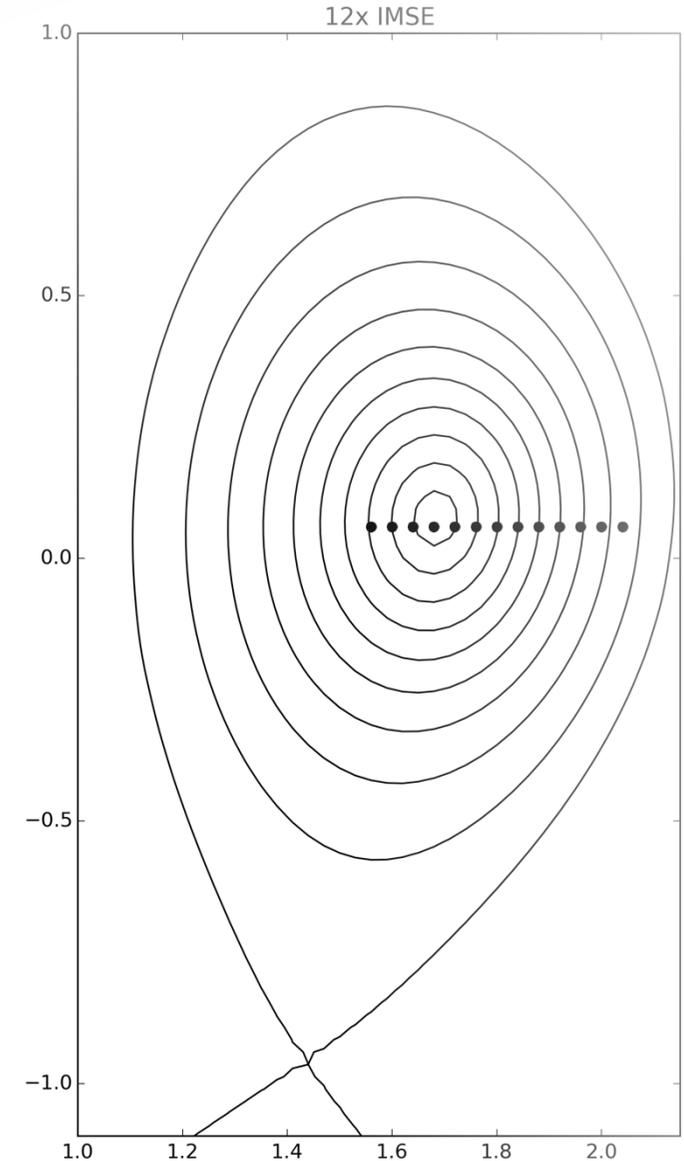
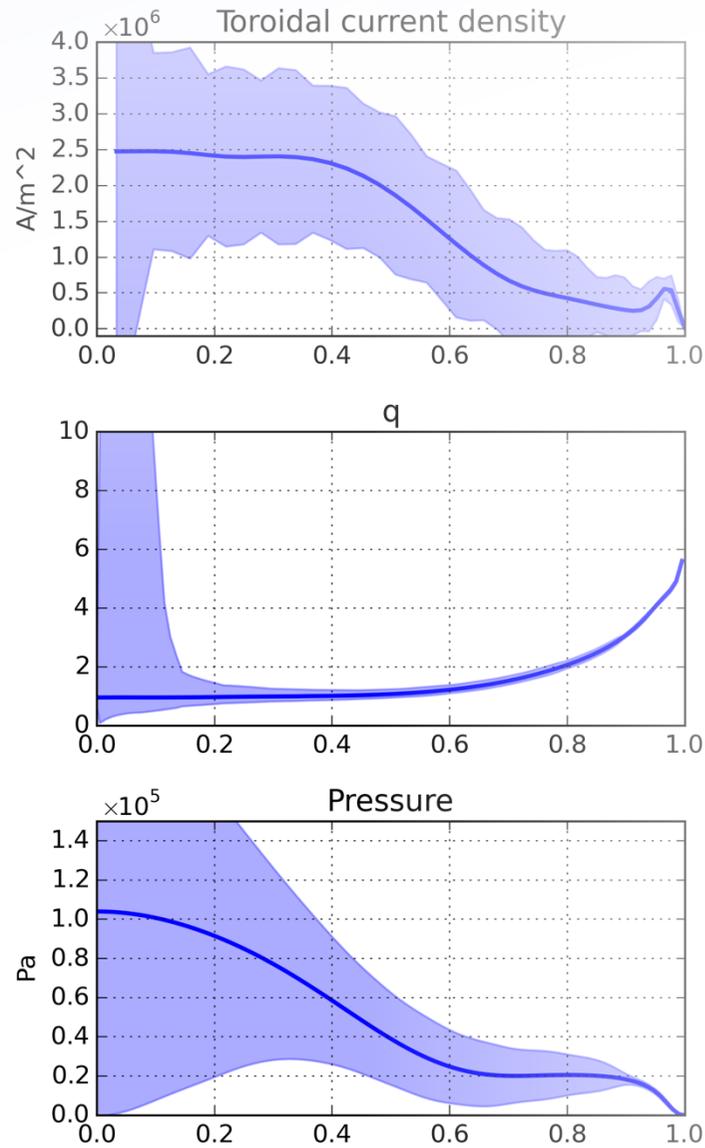
#32232 3.0s

Uncertainties in j/q 1 MSE LOS



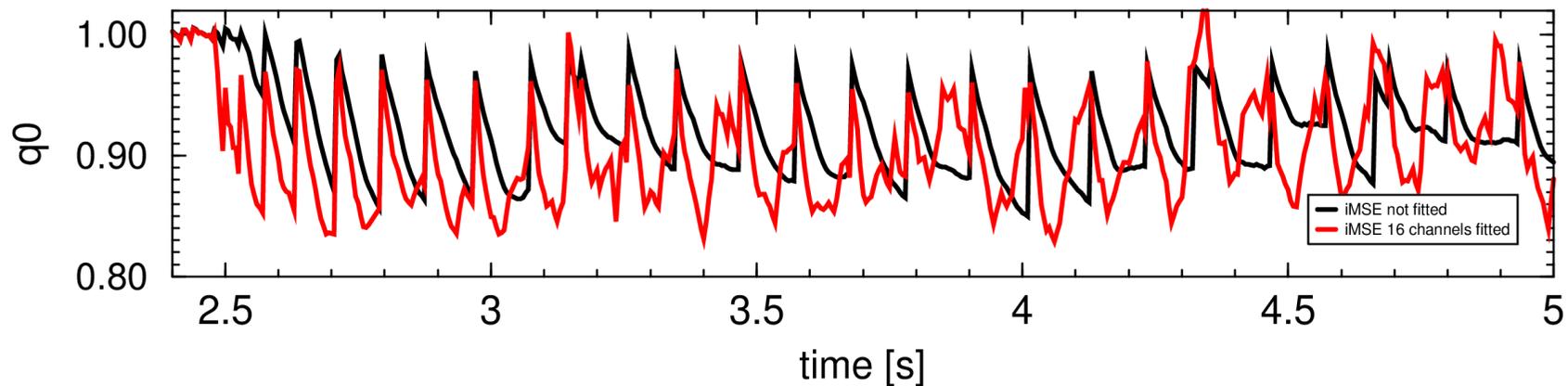
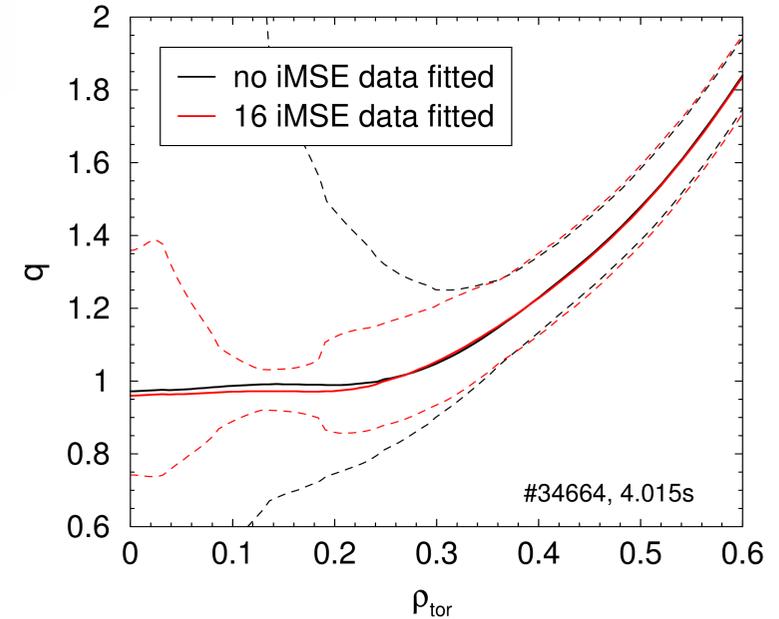
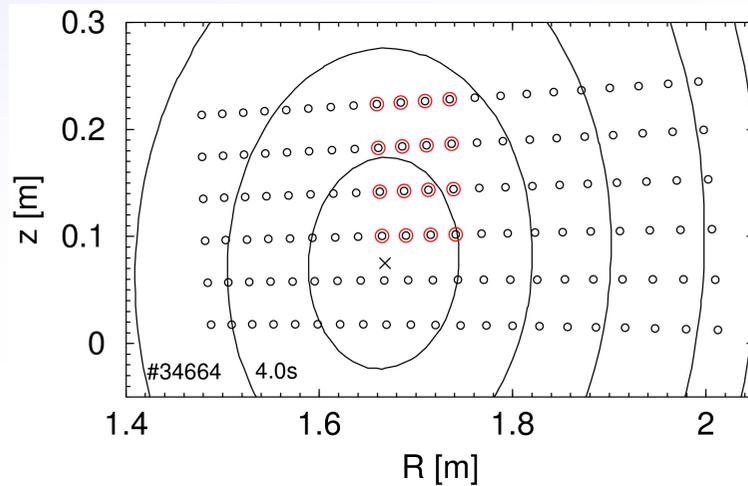
#32232 3.0s

Uncertainties in j/q 12 IMSE LOS



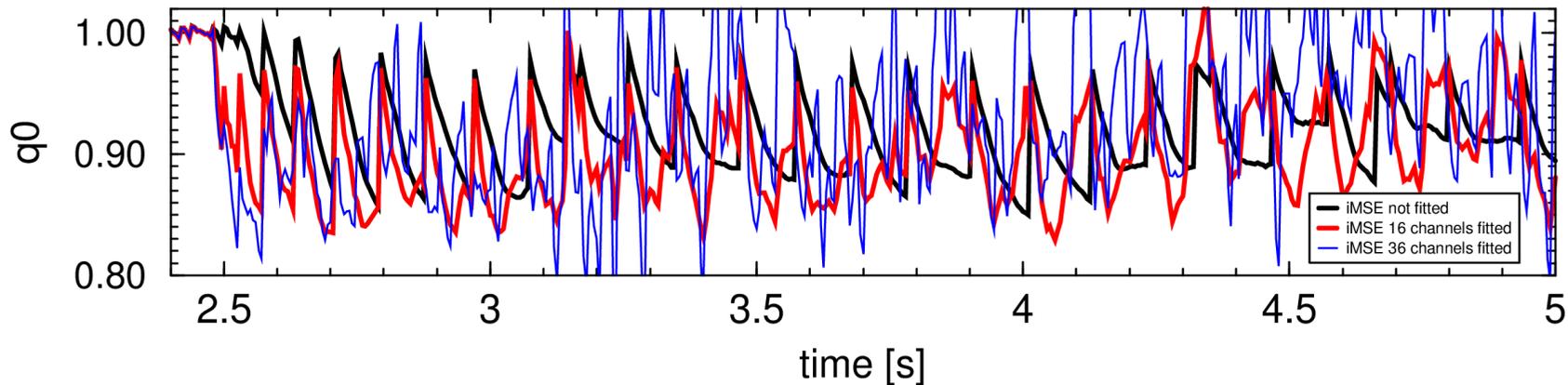
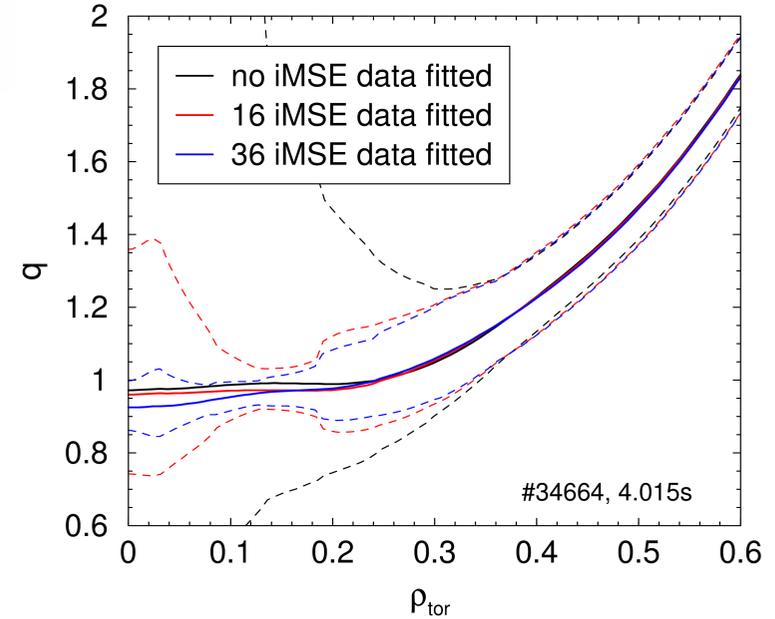
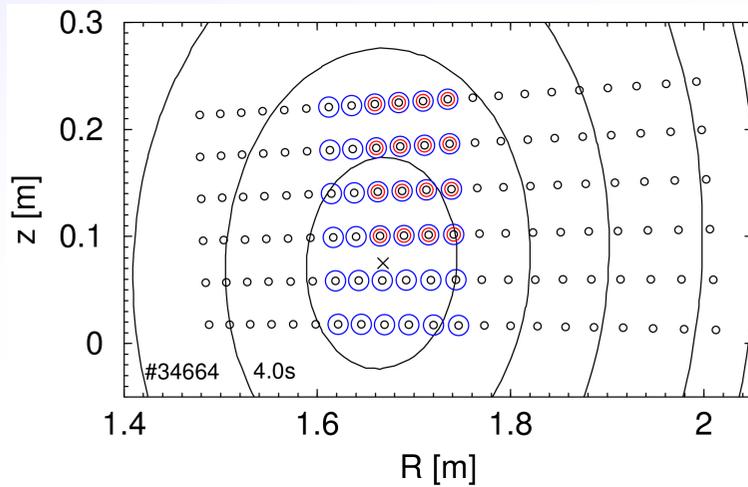
#32232 3.0s

Status: iMSE fit (2)



- 16 iMSE channels appear reasonable, but large q_0 uncertainty

Status: iMSE fit (3)



- 16 iMSE channels appear reasonable, but large q_0 uncertainty
- 36 iMSE channels not conclusive

Note: q_0 estimation is the most challenging problem in current profile reconstruction!