



# Power and particle transport in NBI vs ECRH plasmas

## Profiles Topical Group, May 2021

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(... Interferometry, Diamagnetic Loop, ECRH ...)

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...

## Background:

- Gas fuelled ECRH discharges:
  - Flat density profiles
  - $T_i$  clamped at 1.5keV because:
    - 1) Poor coupling at low collisionality,
    - 2) Turbulence and stiff profiles
    - 3) Te/Ti exacerbates turbulence
  - Low and flat impurity densities
- Pellets:
  - Core fuelling --> peaked density profiles
  - Turbulence suppression.  $Q_i$  reduced to  $O(\text{neoclassical})$ 
    - > Highest observed  $T_i$

Great but.... Only seen *after* rapid pellets. Can steady state pellets give peaked density?

If not, what can we do?

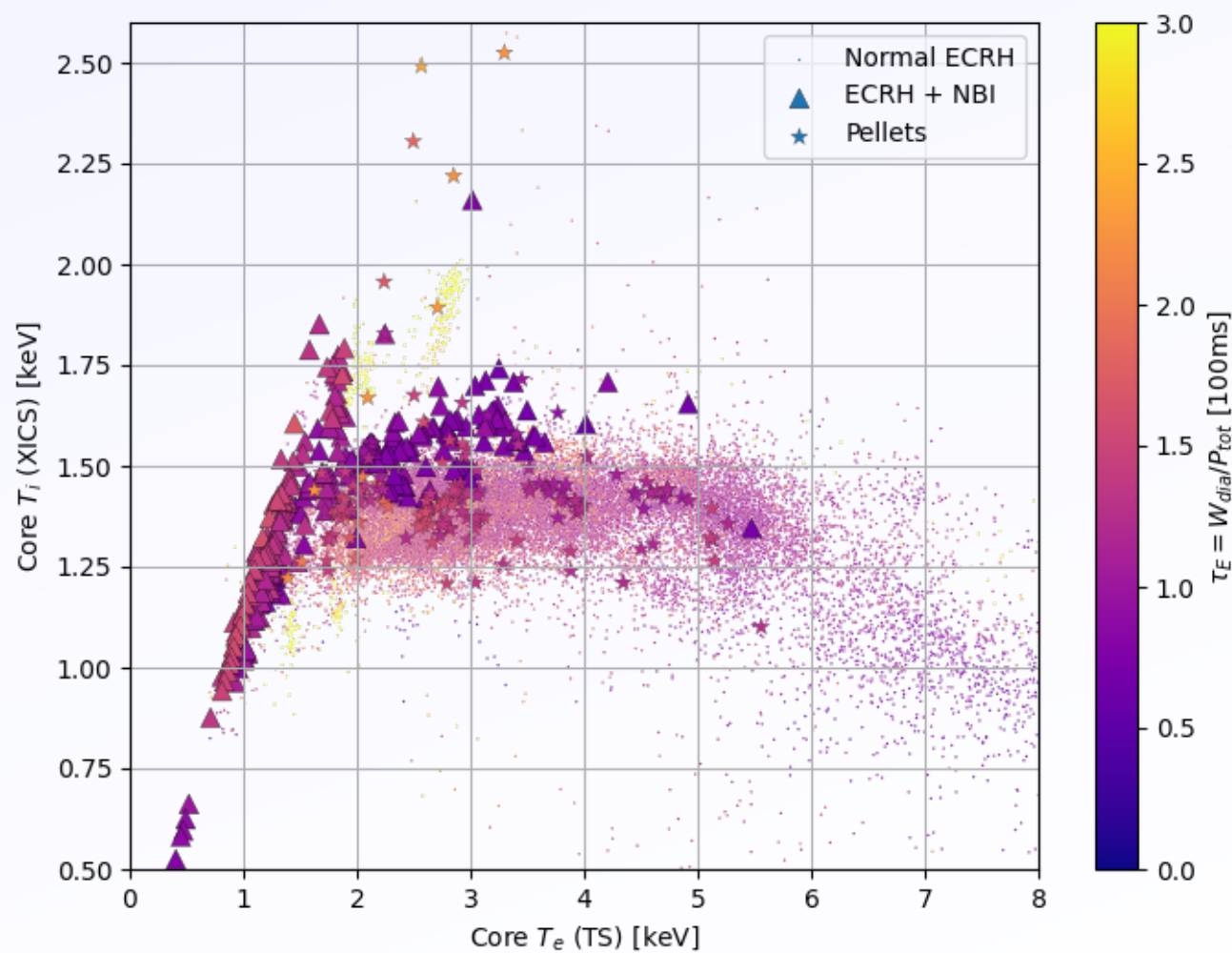
NBI gives continuous core fuelling and ion heating. Can NBI provide a route to improved performance (higher  $T_i$ )?

Note: I focus here on high  $T_i$ , not  $\beta$ , so I am ignoring  $T_e$  'performance'.

## Global view

How does global  $T_i$  look?

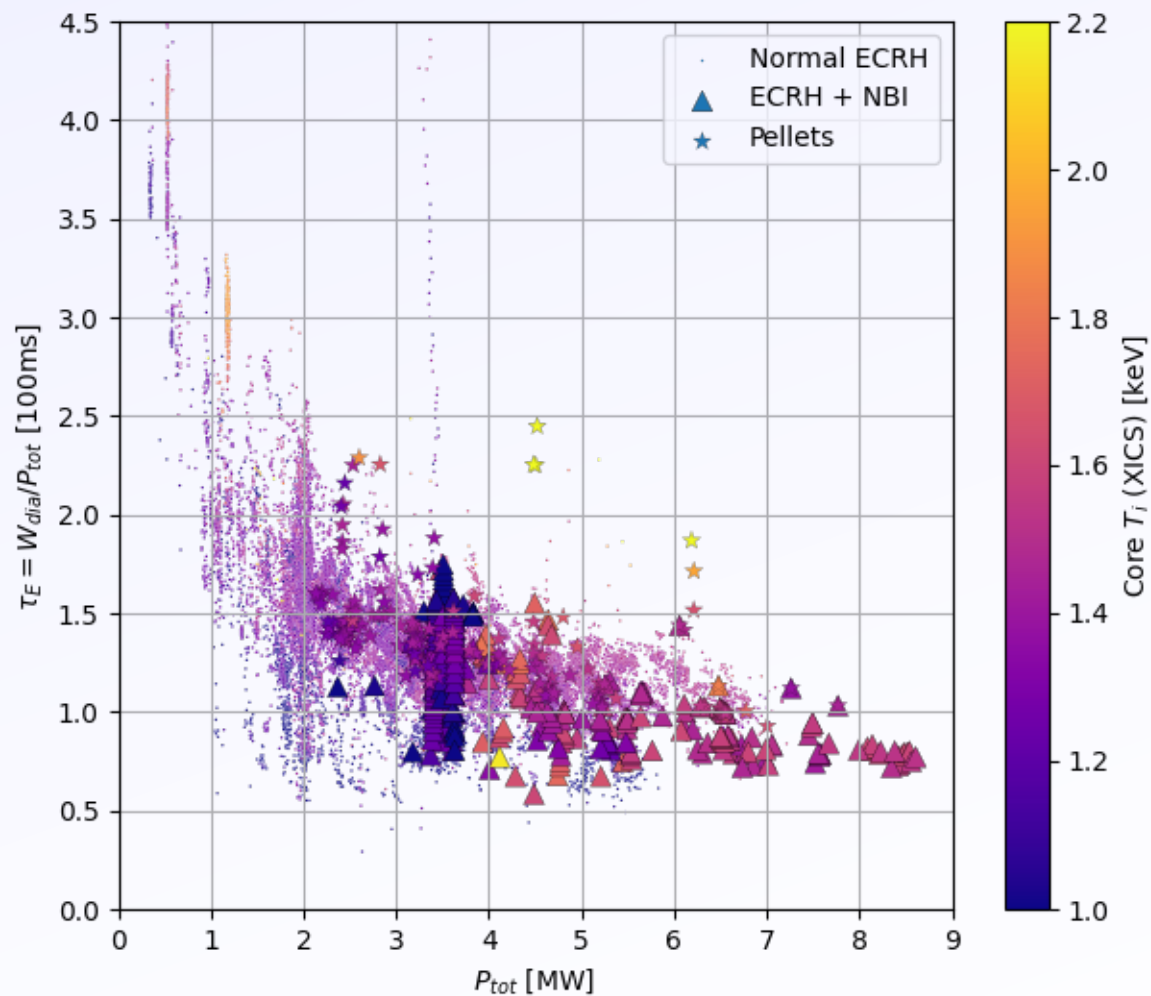
- Most shots are ECRH + NBI.
- $T_i$  still around clamping limit, maybe slightly higher but generally not as high as post-pellets plasmas.
- Some of the highest  $T_i$  are at lower  $T_e$ .



## Global view

Global confinement generally lower for NBI compared to ECRH due to lower efficiency of NBI heating physics:

- 1) Significant fast-ion loss fraction  $\gg$  ECRH stray radiation
- 2) Power deposition profile much broader

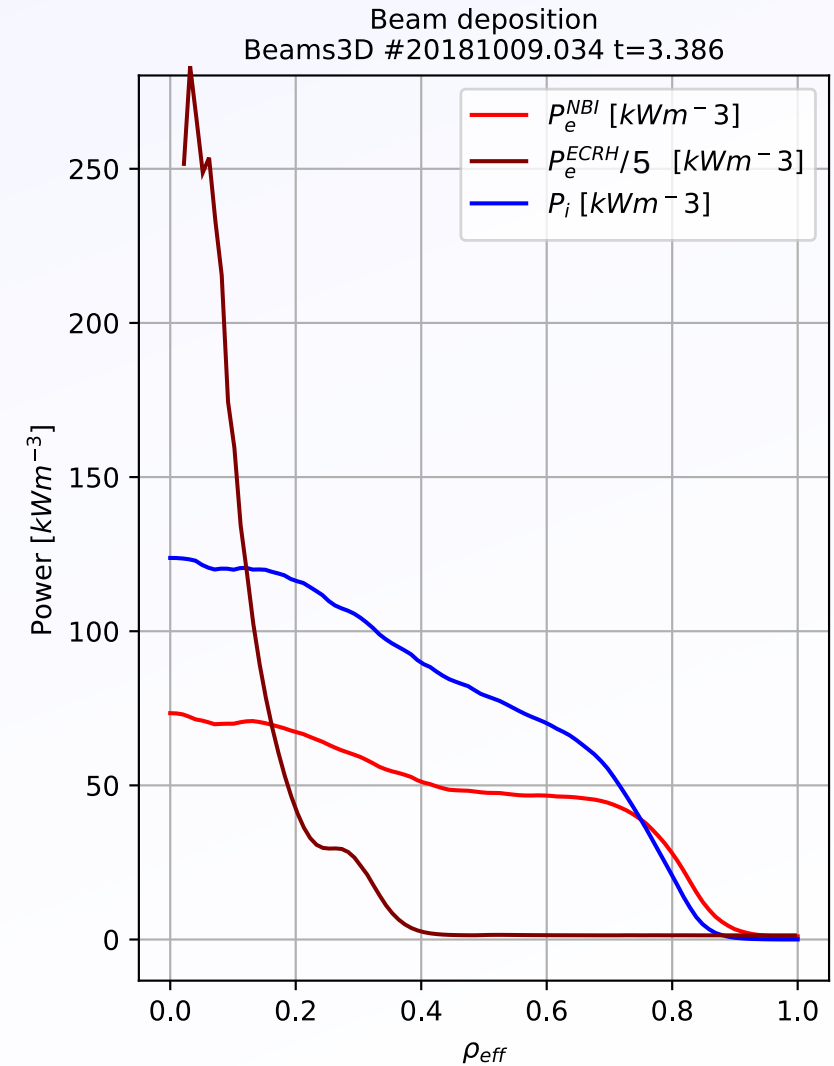
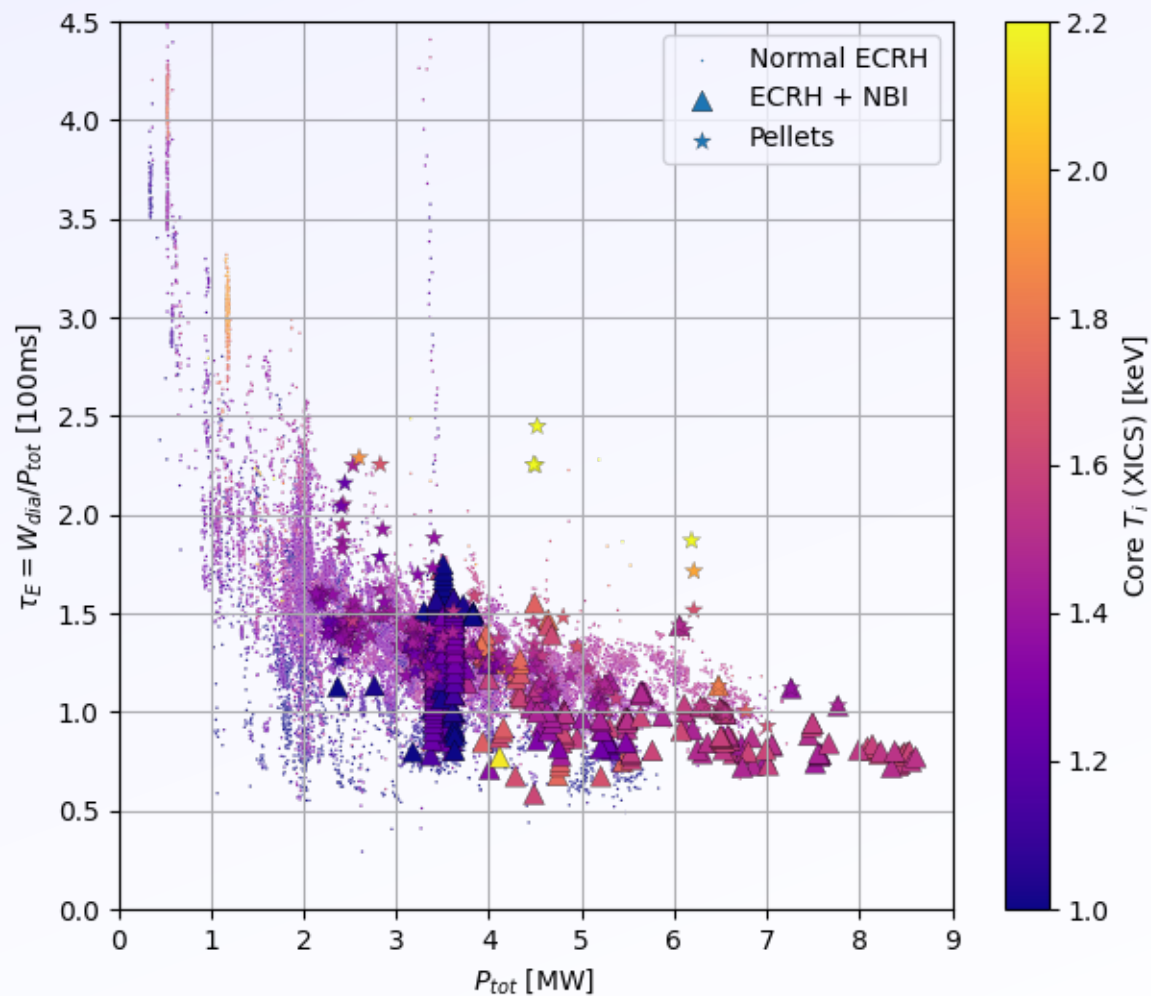




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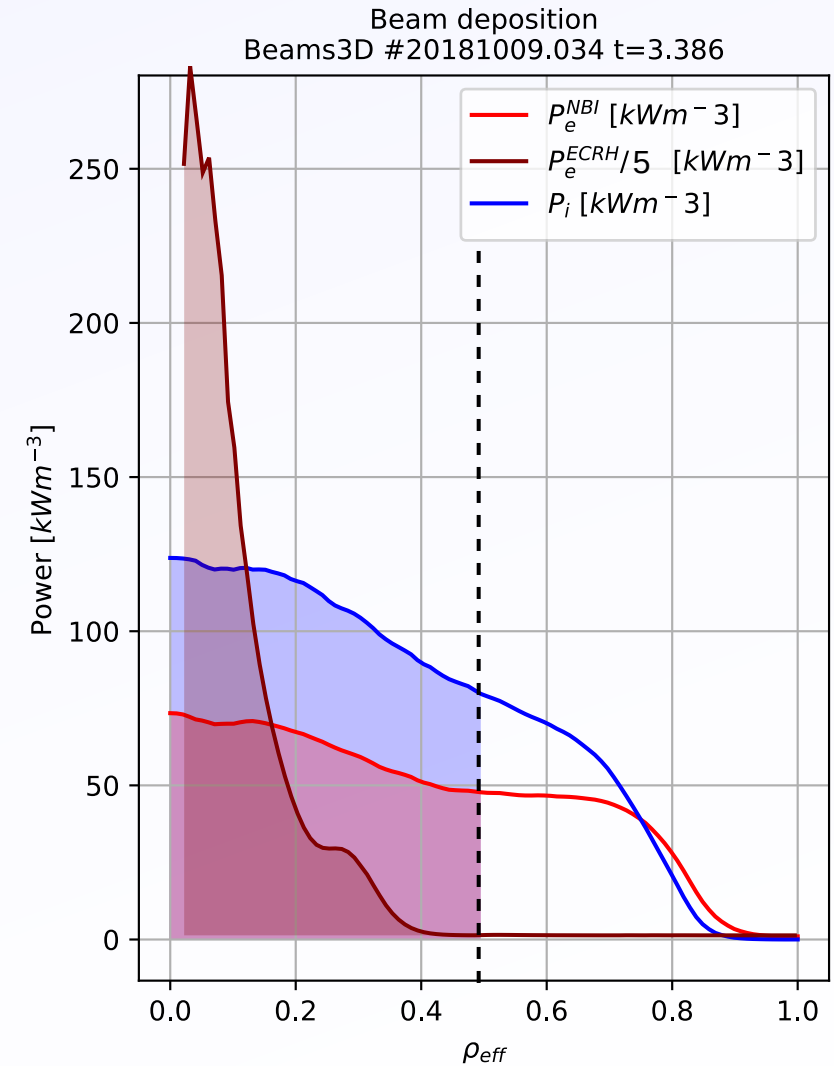
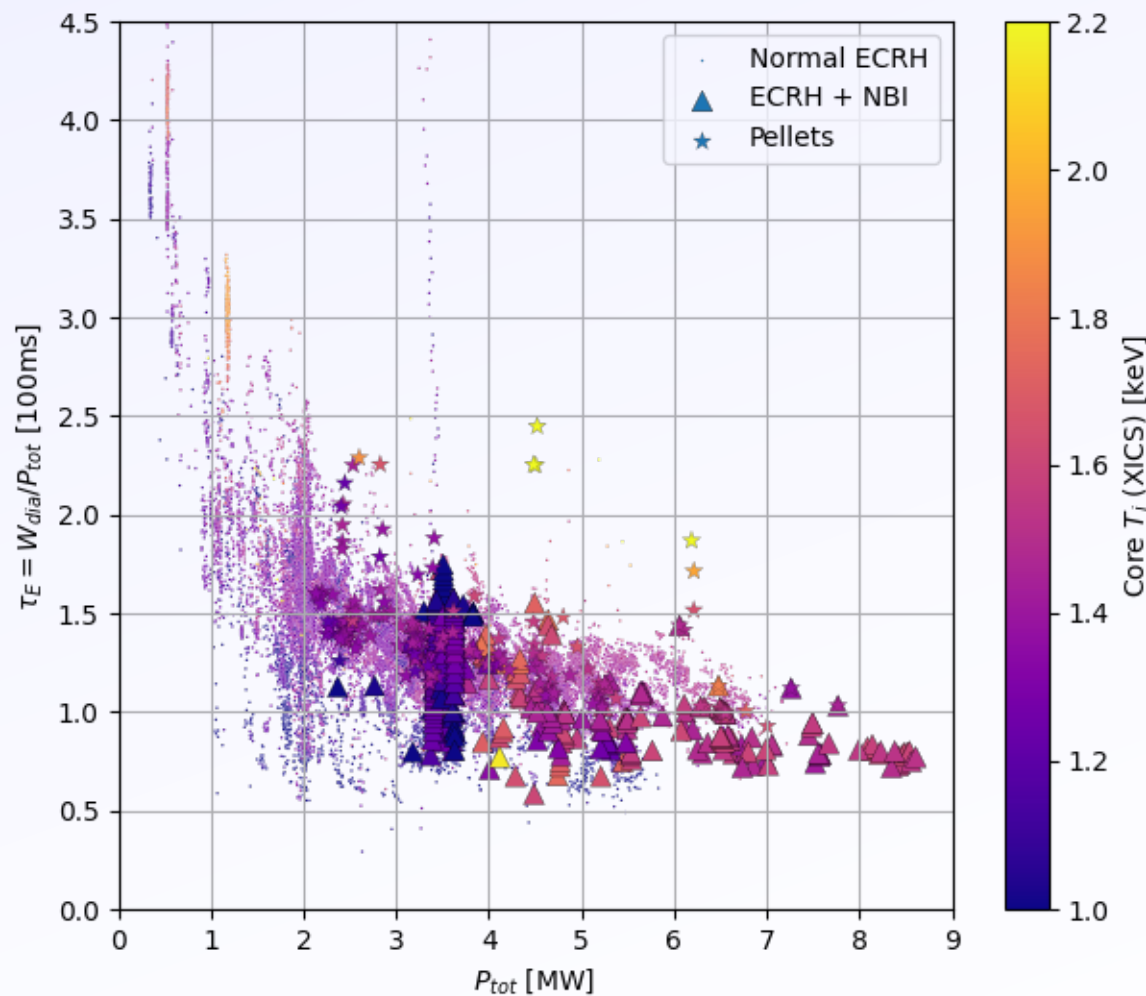
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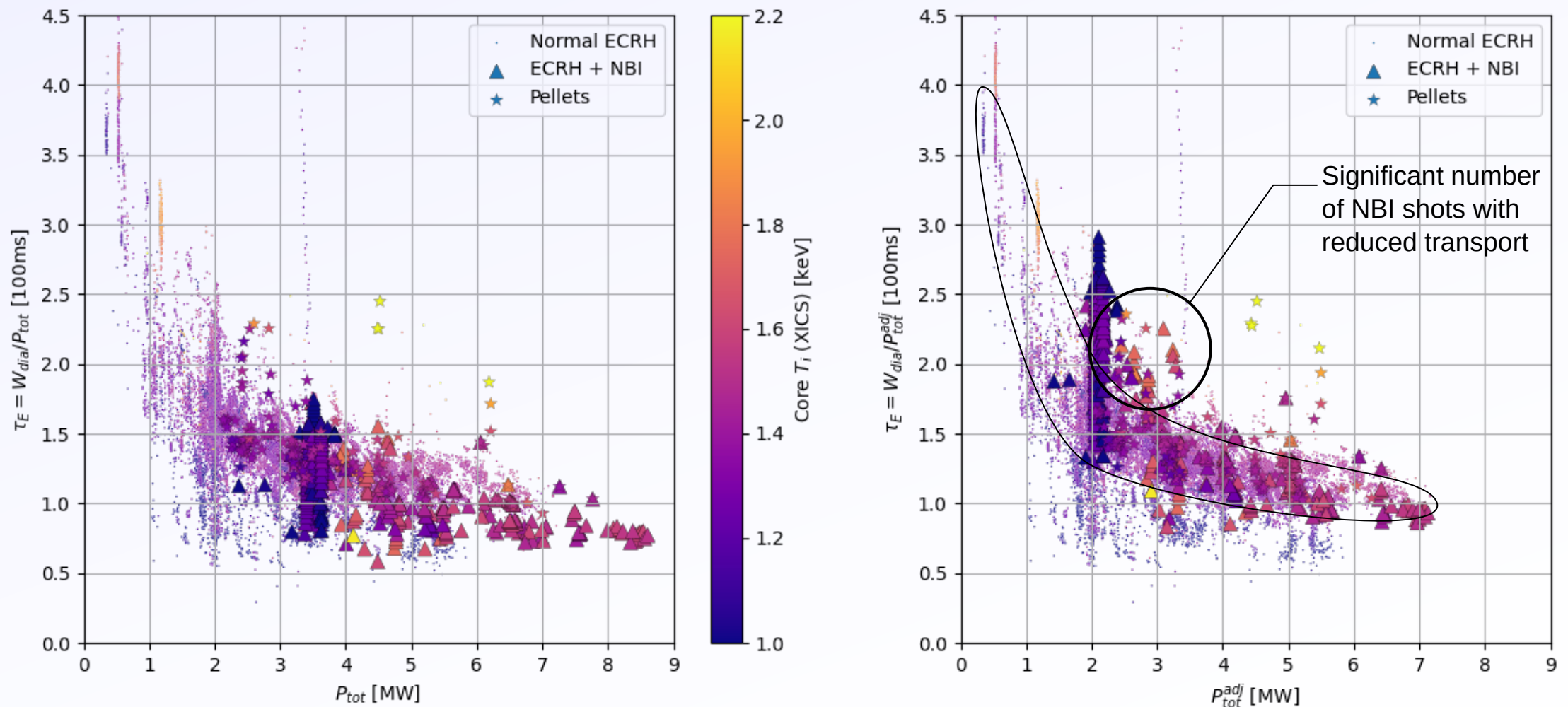


## Confinement vs Transport

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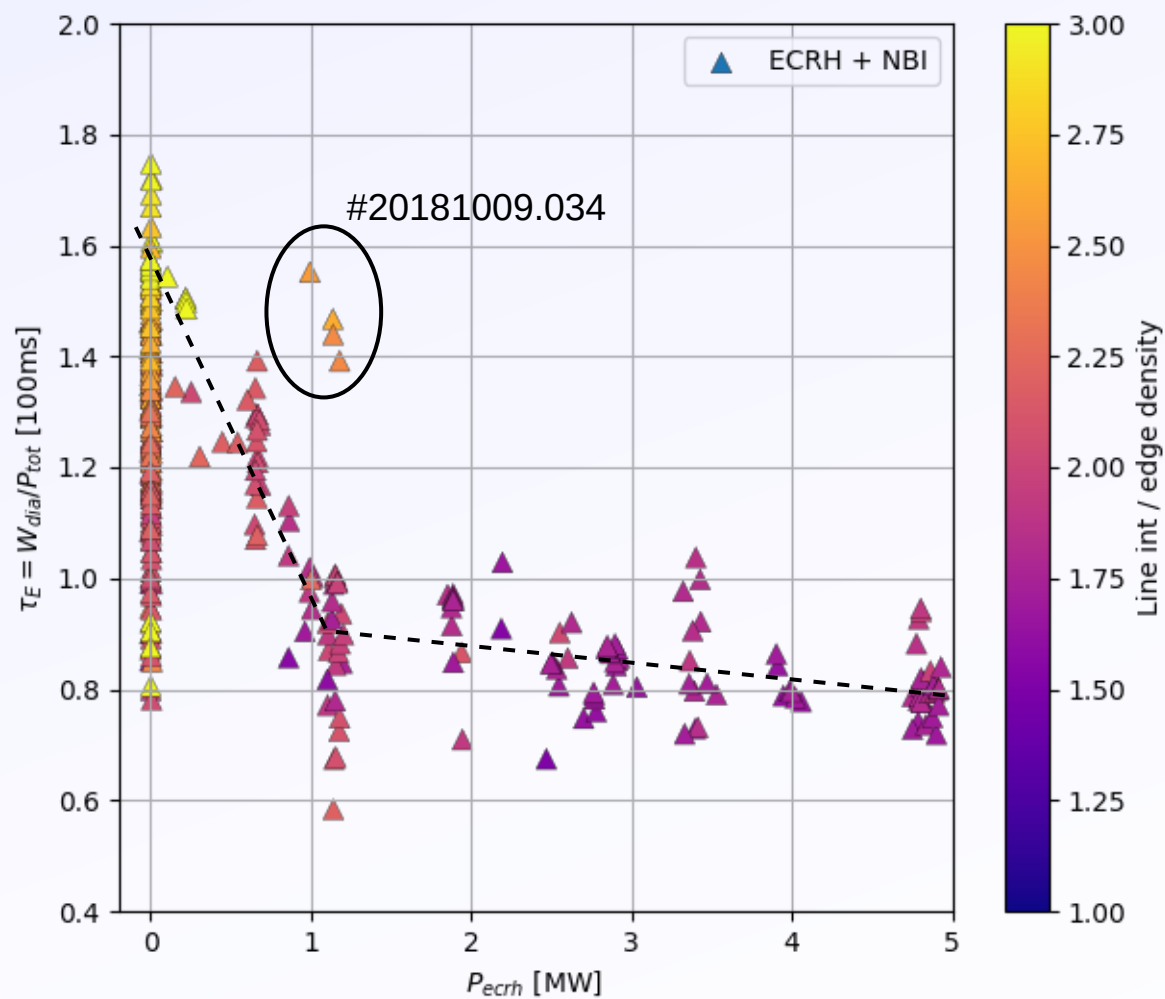
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To consider *transport* rather than *confinement*, examine adjust to e.g.  $P_{\text{total}} = P_{\text{ECRH}} + 60\% P_{\text{NBI}}$ :



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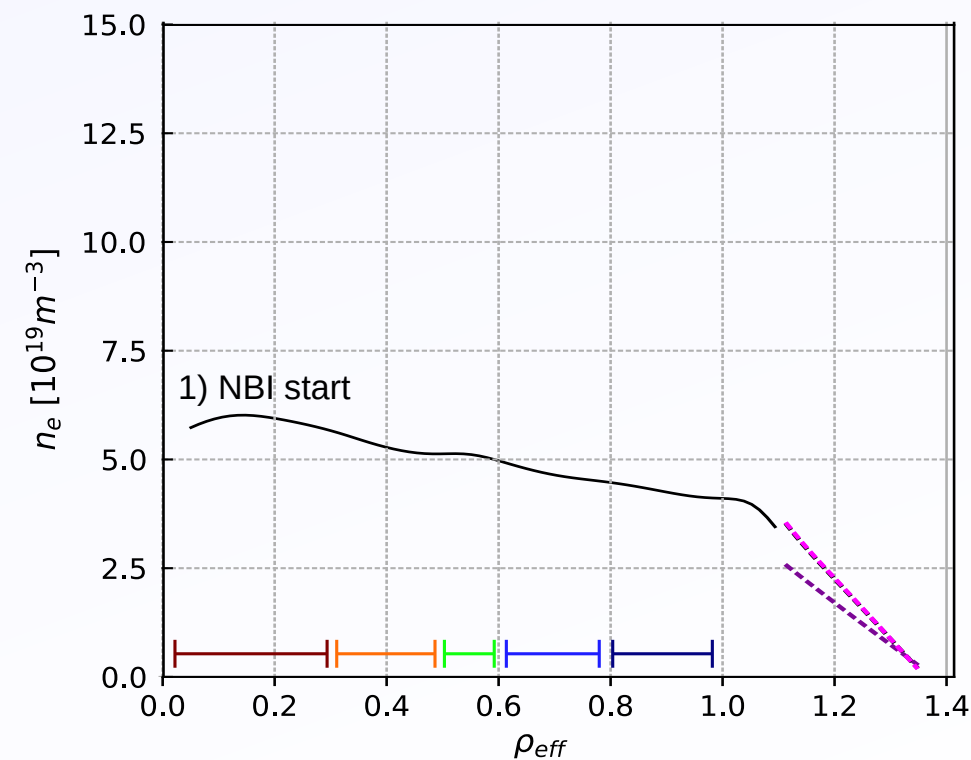
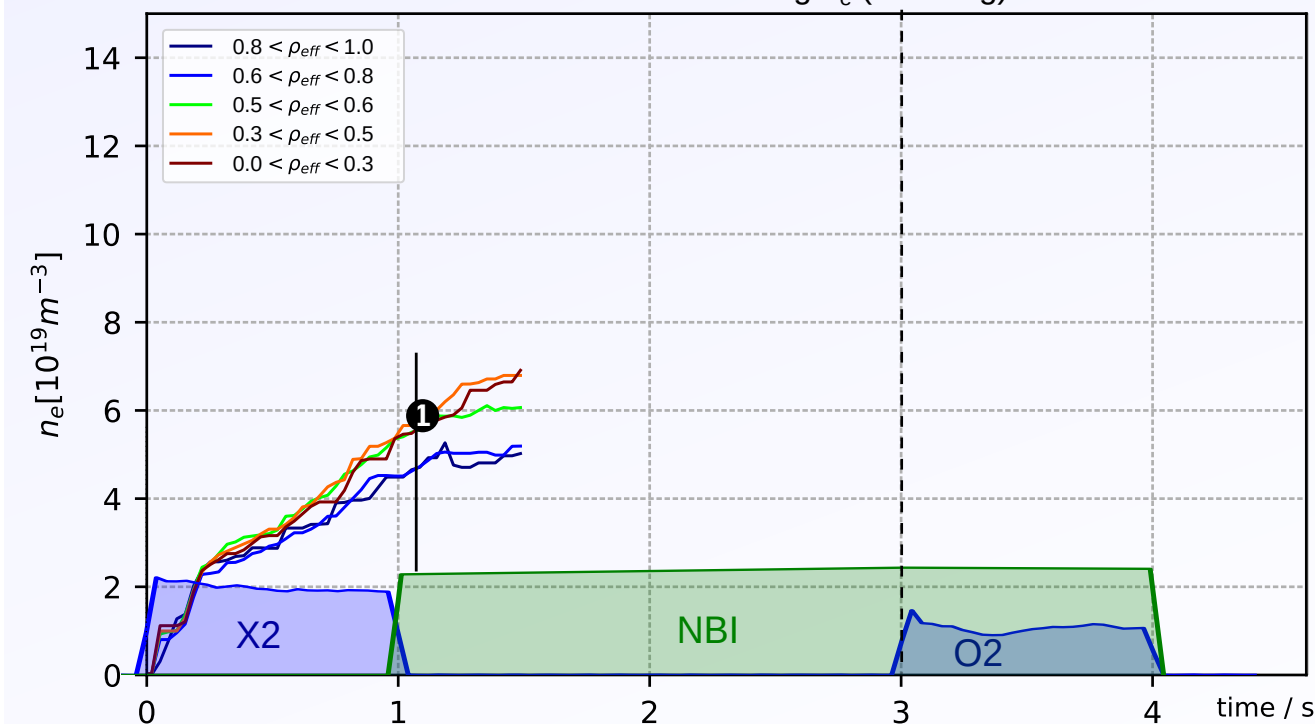
Within the NBI shots, ECRH  $\geq 1$  MW quickly degrades performance.  
ECRH  $< 1$  MW shots show some density peaking.



## Pure NBI

- Pure NBI discharges show core density and impurity peaking (*almost* all of the time!).
- Strong density rise occurs
  - $\rho_{eff} < 0.5$ .
  - $t > t_{onset}$ , which varies over 1 - 2s after NBI in different shots. No apparent correlation of  $t_{onset}$  with external events.

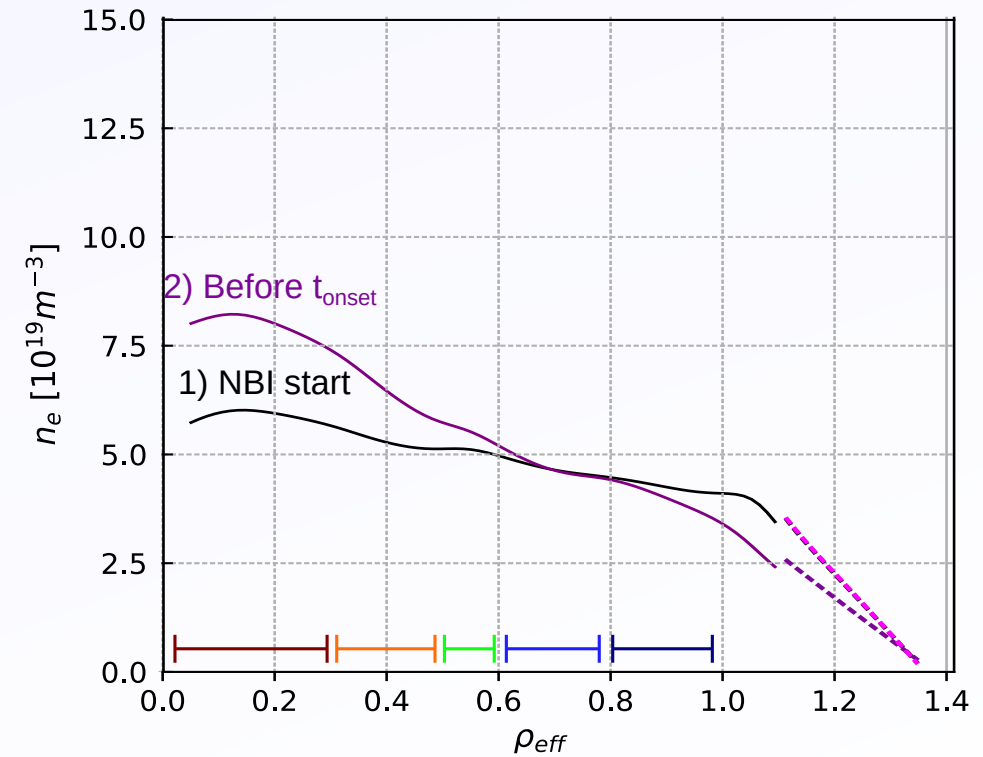
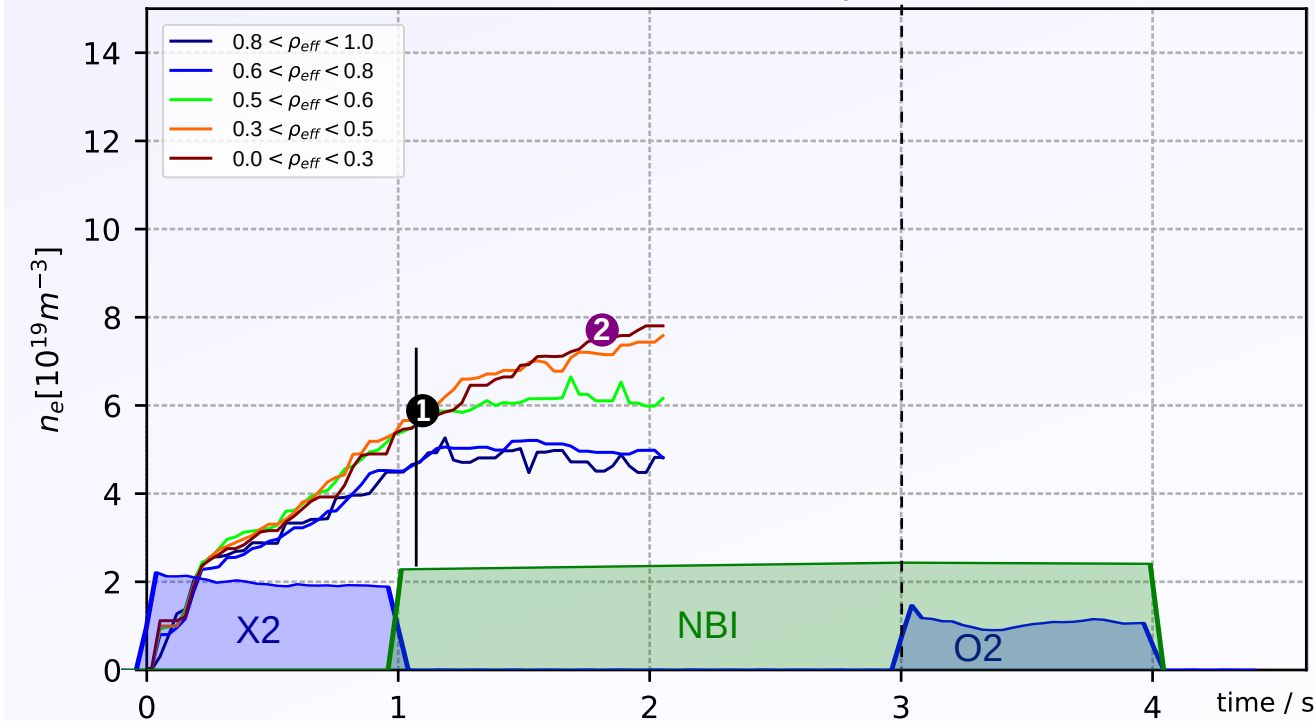
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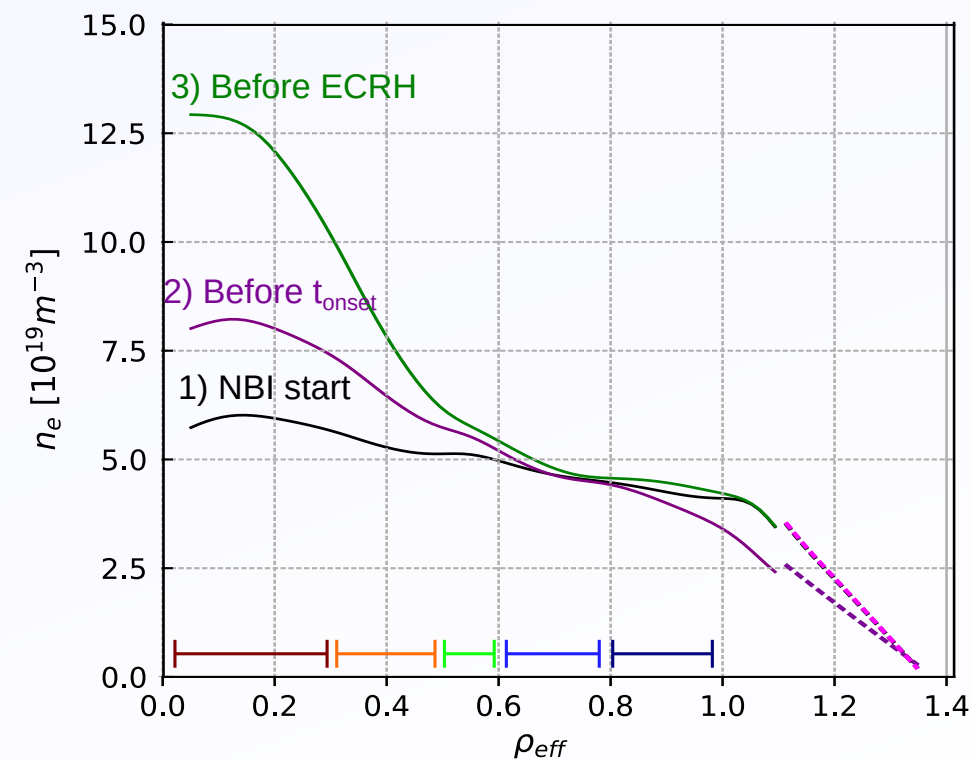
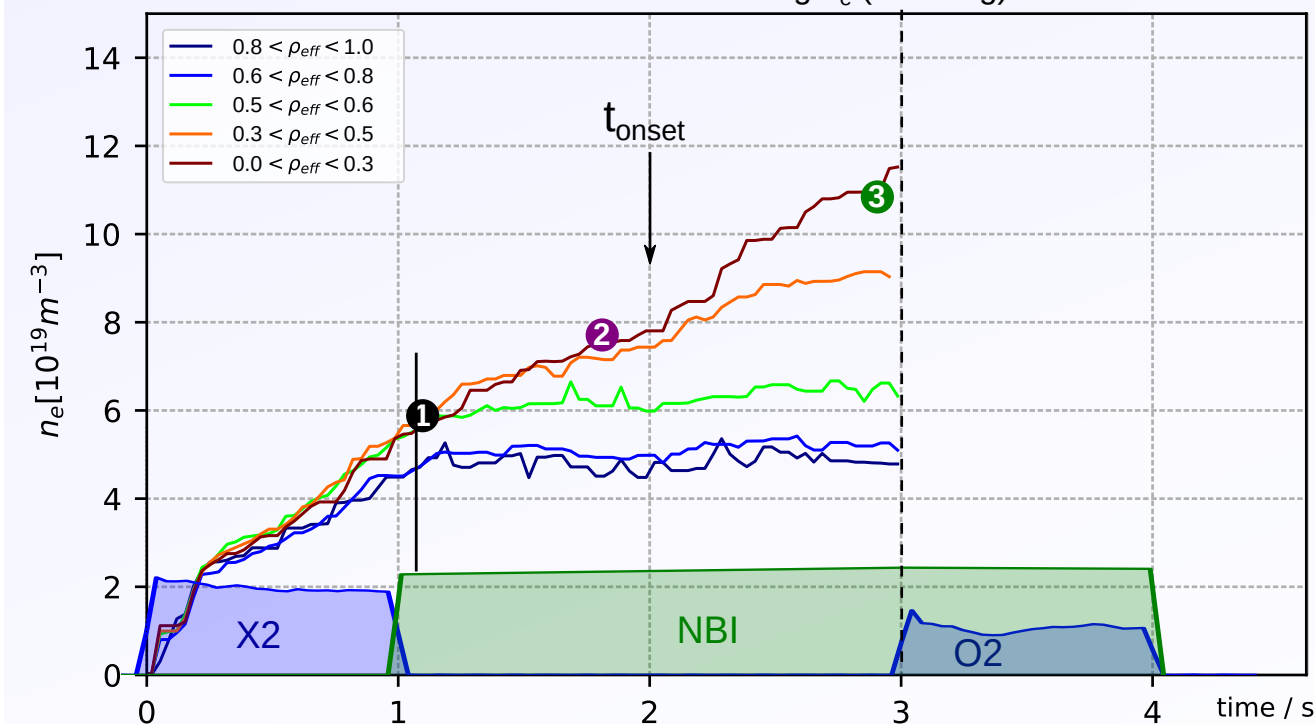
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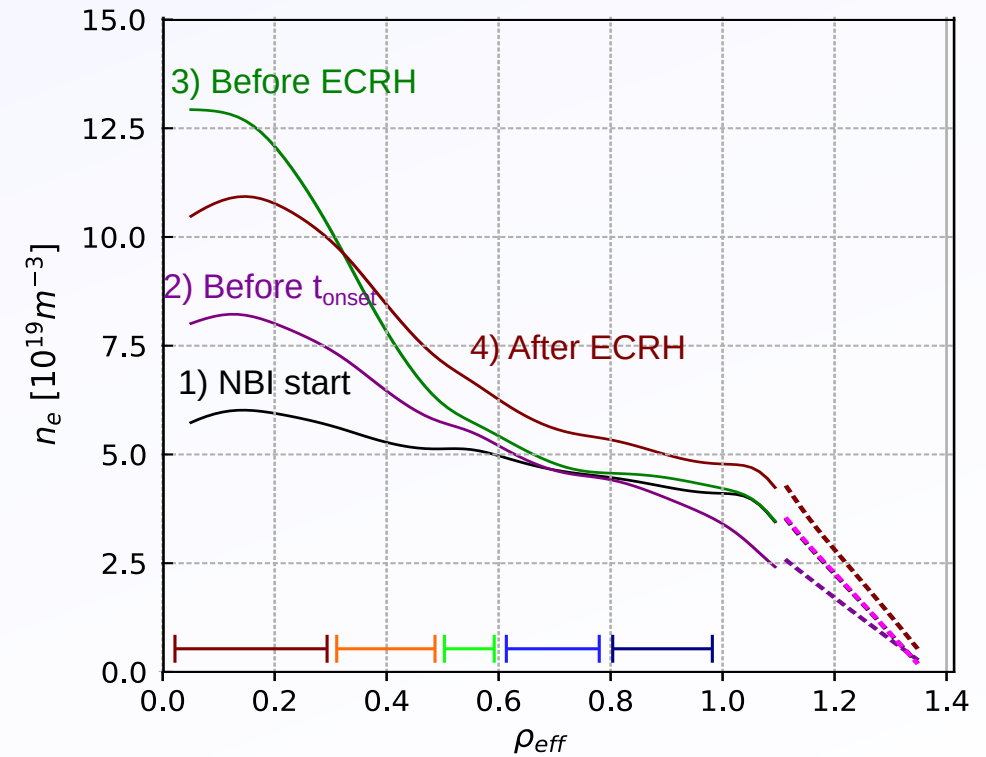
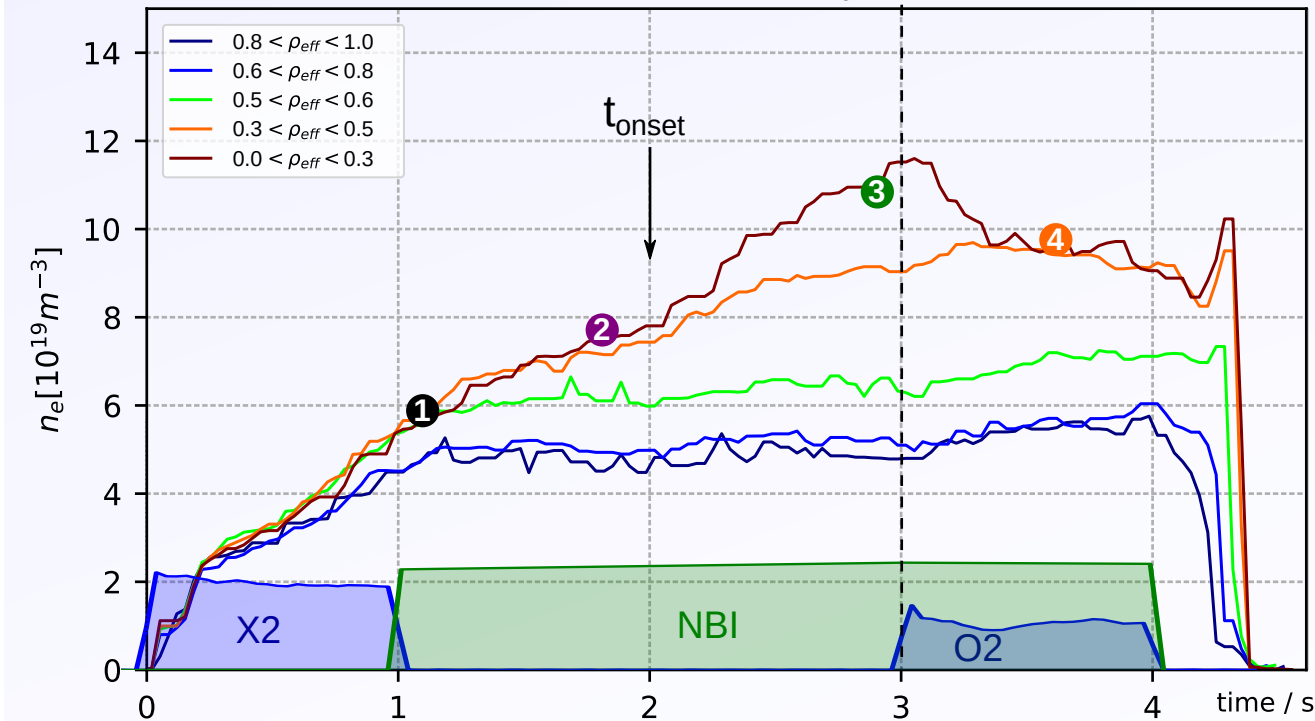




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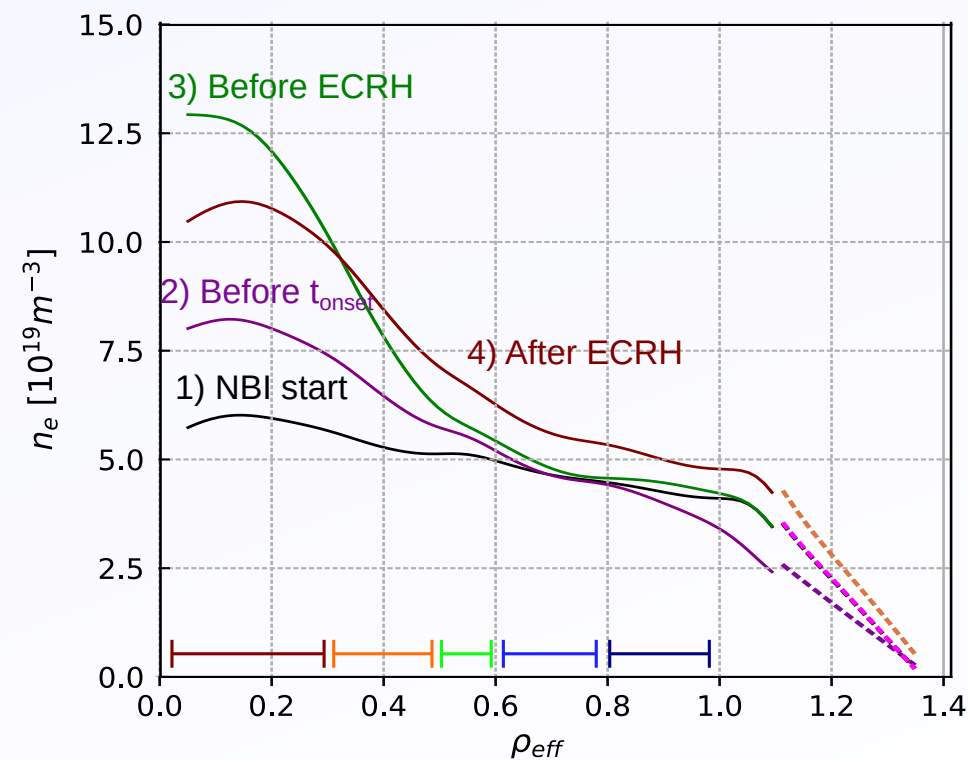
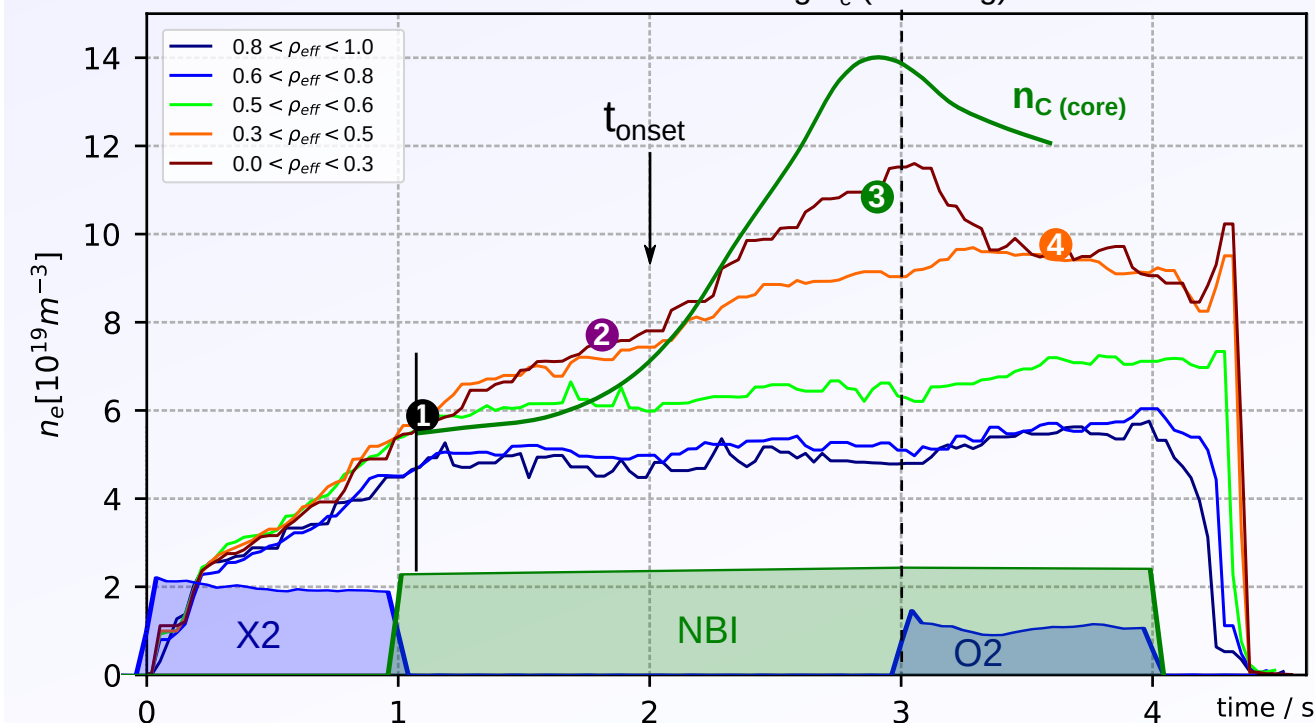




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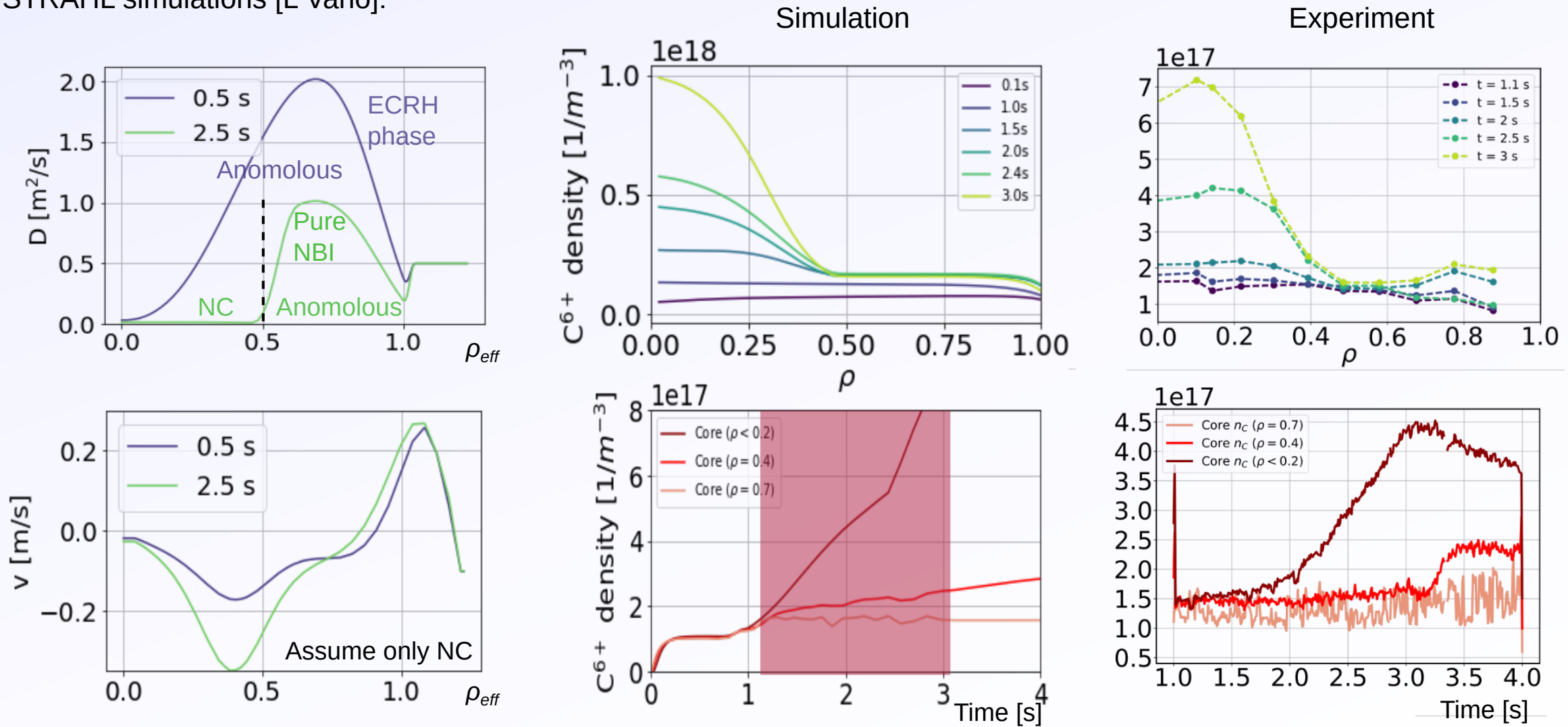
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## Pure NBI

Carbon peaking consistent with neoclassical transport for  $\rho < 0.5$  and some turbulent diffusion for  $\rho > 0.5$ .  
STRAHL simulations [L Vanó]:



Peaking (supressed turbulent diffusion) starts at same onset time of accelerated core electron density peaking ( $t=2.2$ s)

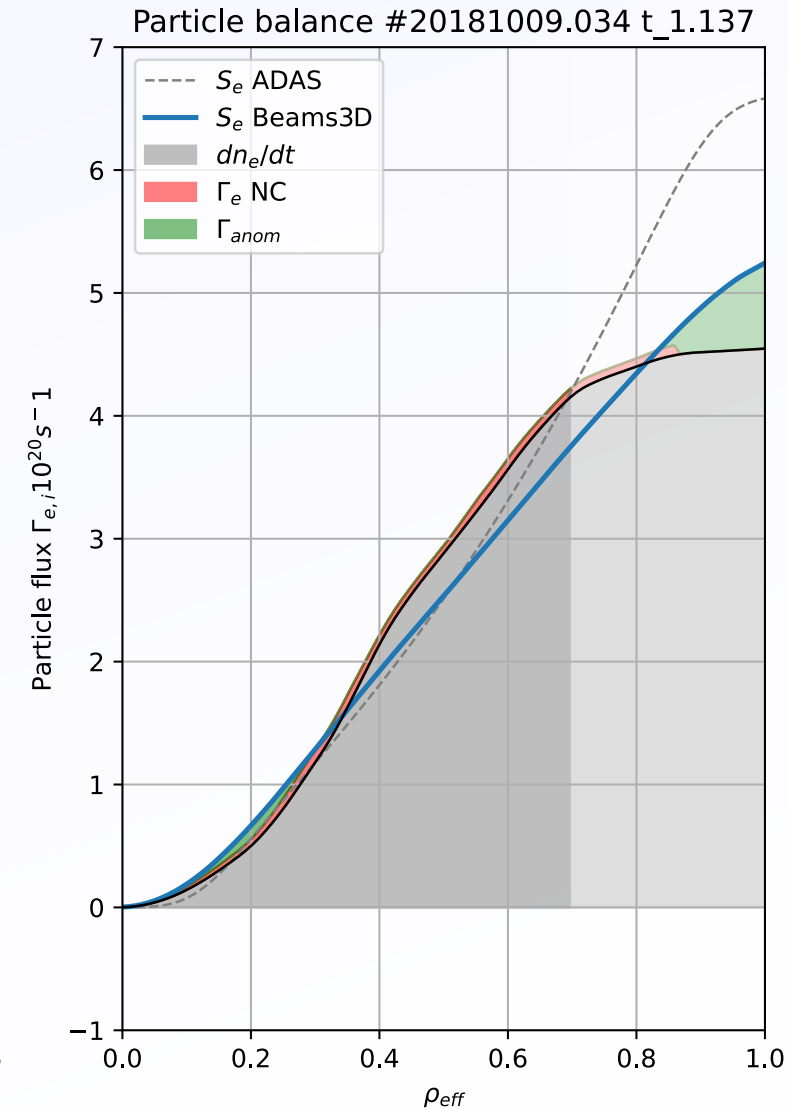
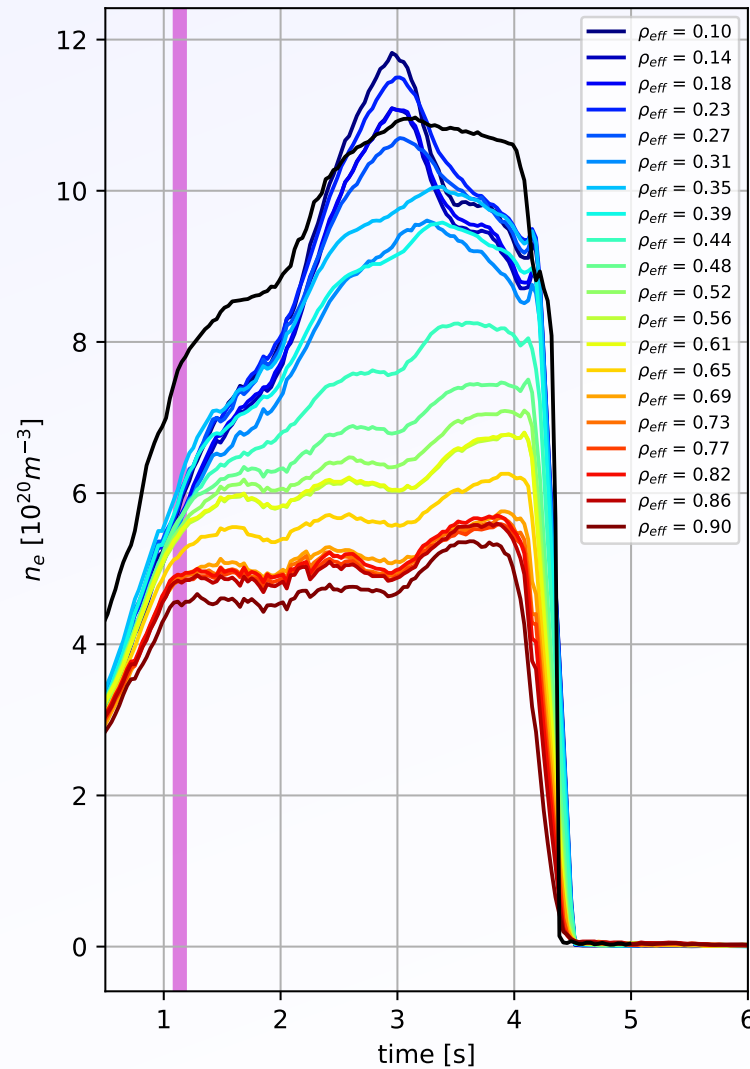
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### Particle balance:

- Source rate from Beams 3D. Roughly agrees with ADAS beam stopping. No Halo diffusion but not significant.
- NC particle fluxes calculated using NEOTRANSP. Robust to uncertainties: Profiles, Te-Ti, Zeff, Er --> no more than ±20%.
- Ignore gas fuelling and recycling --> Maybe invalid for rho > 0.7

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NC particle flux insignificant.  
Density rises at fuelling rate.  
--> **No anomolous flux**



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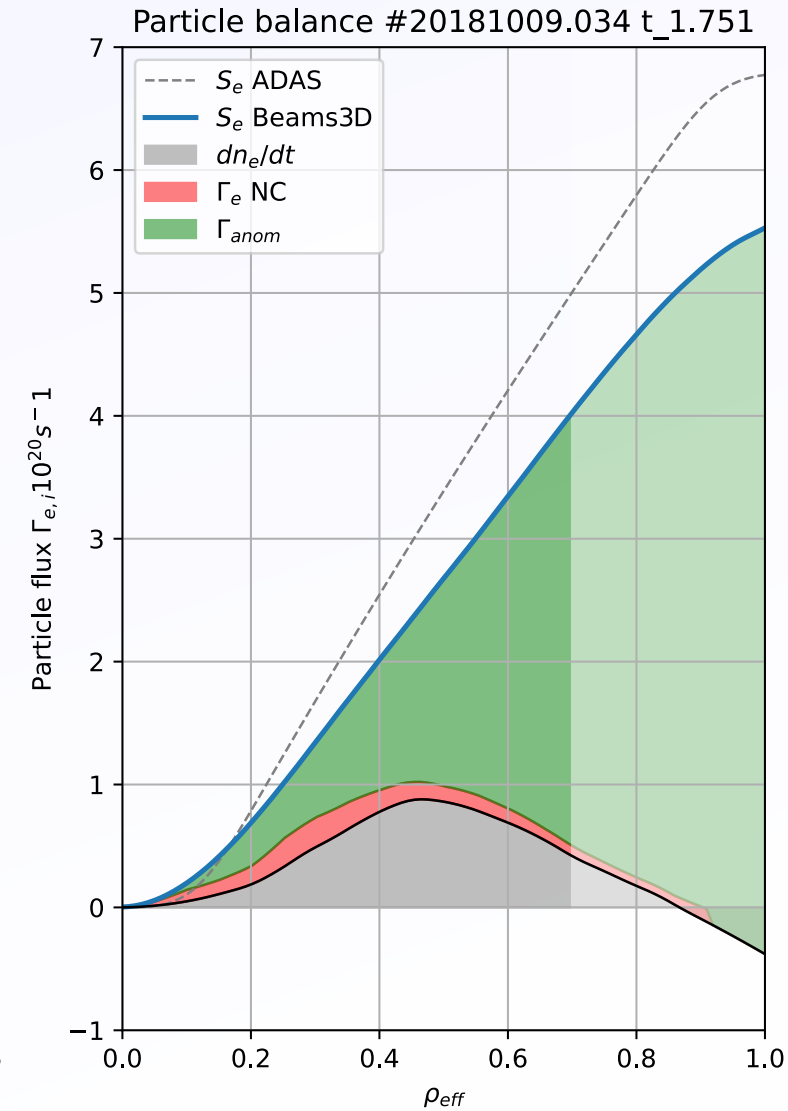
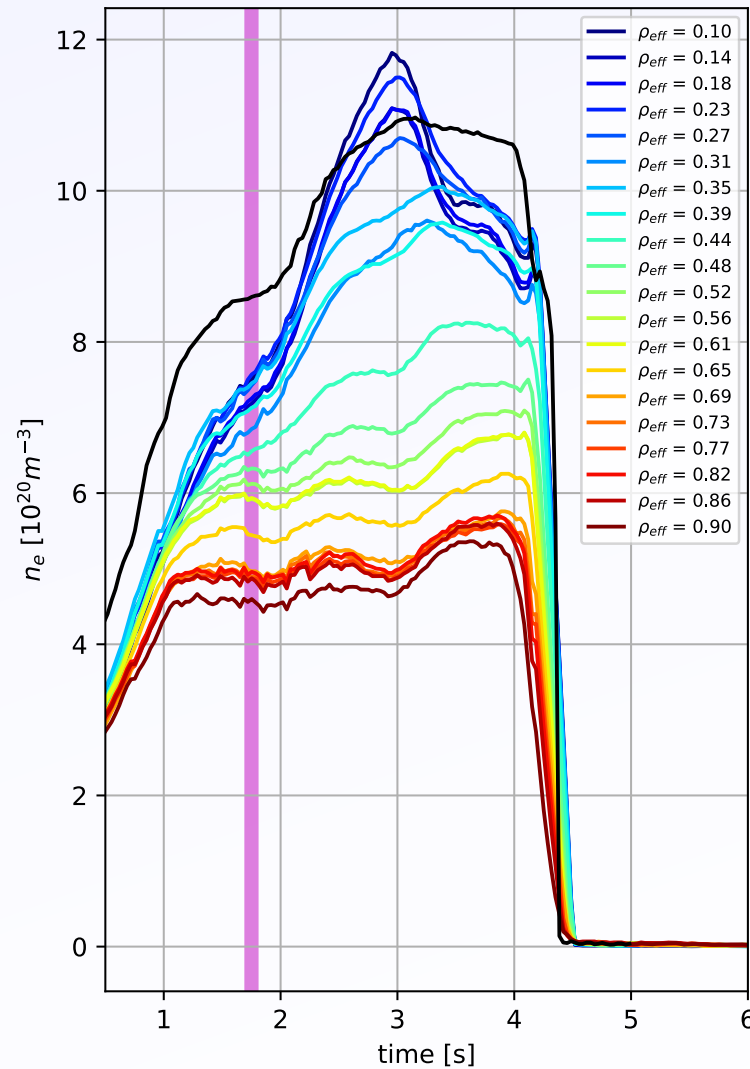
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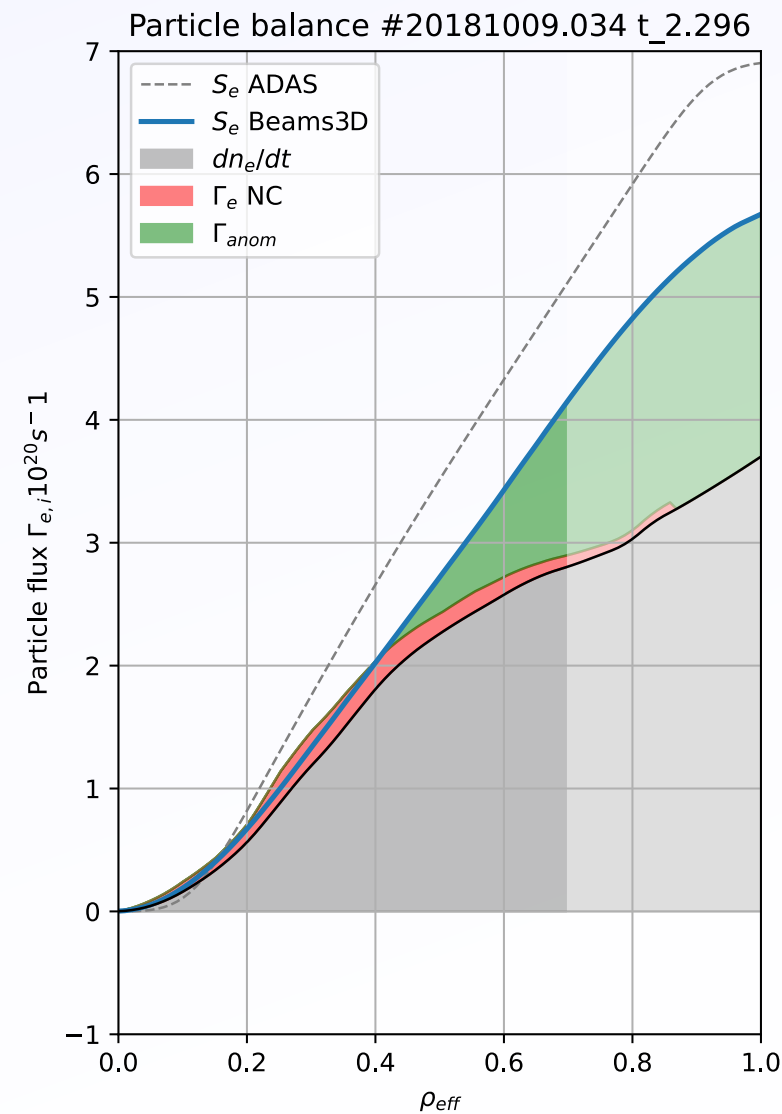
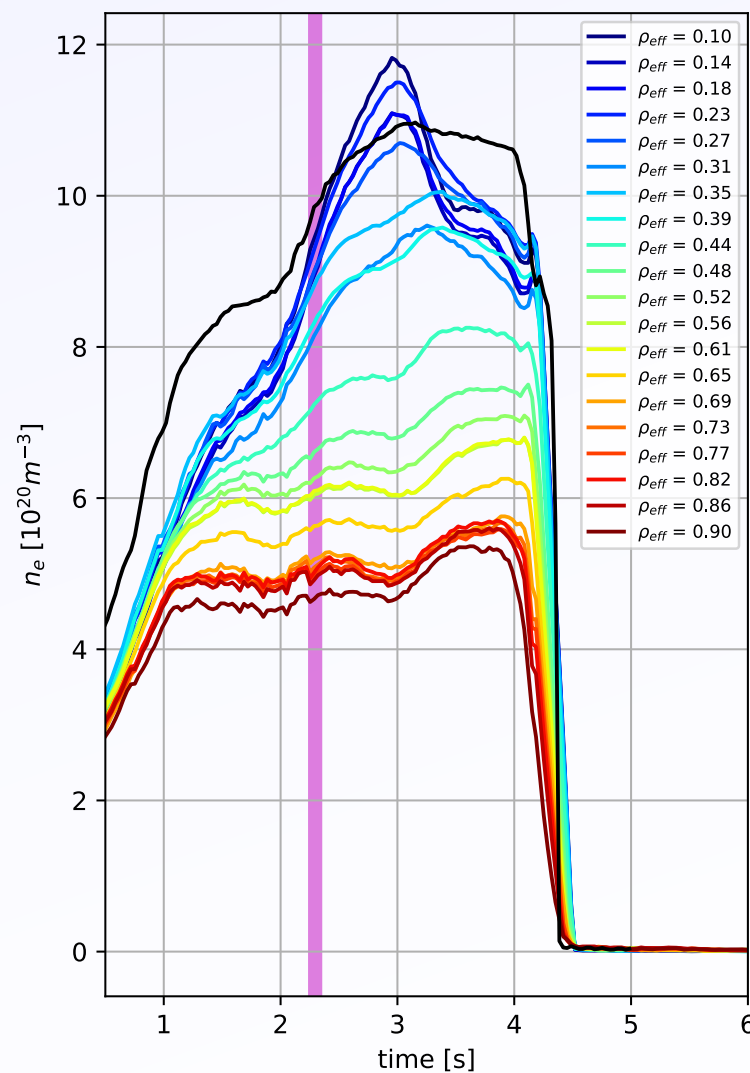
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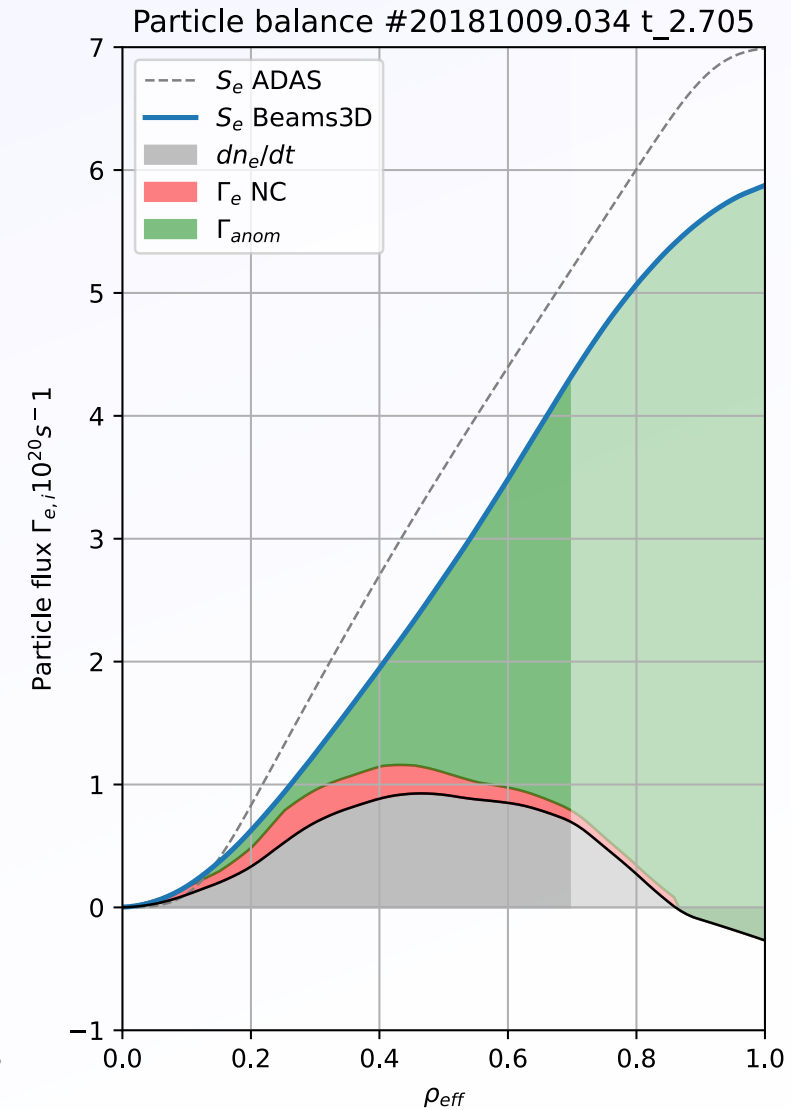
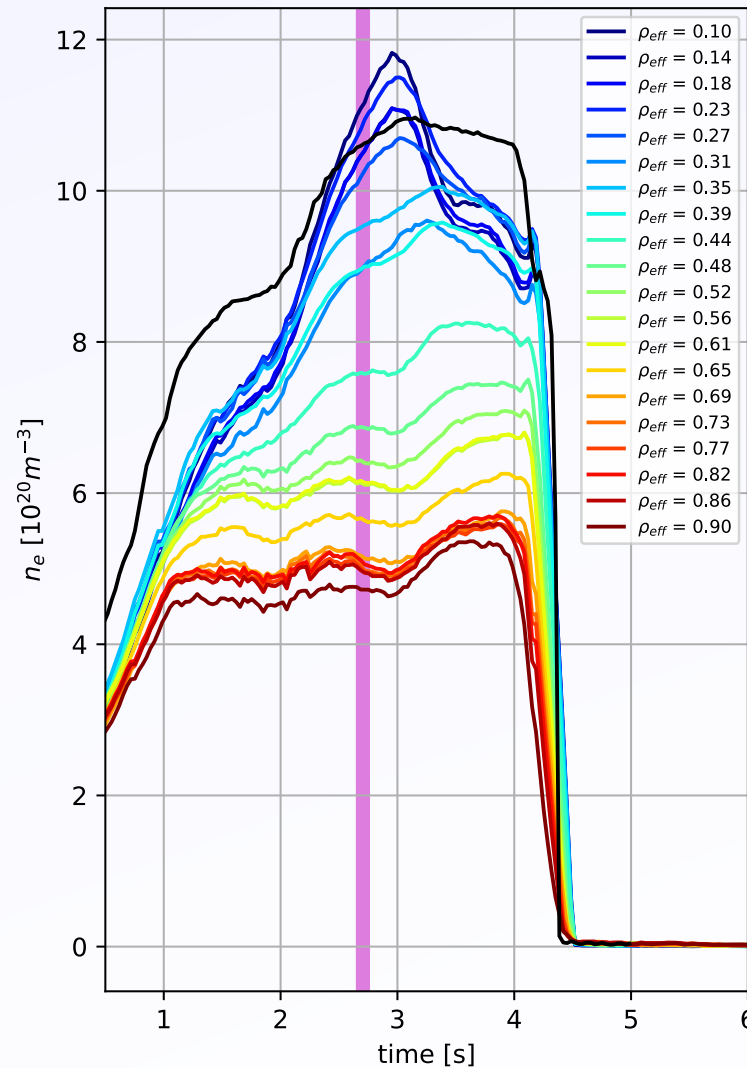
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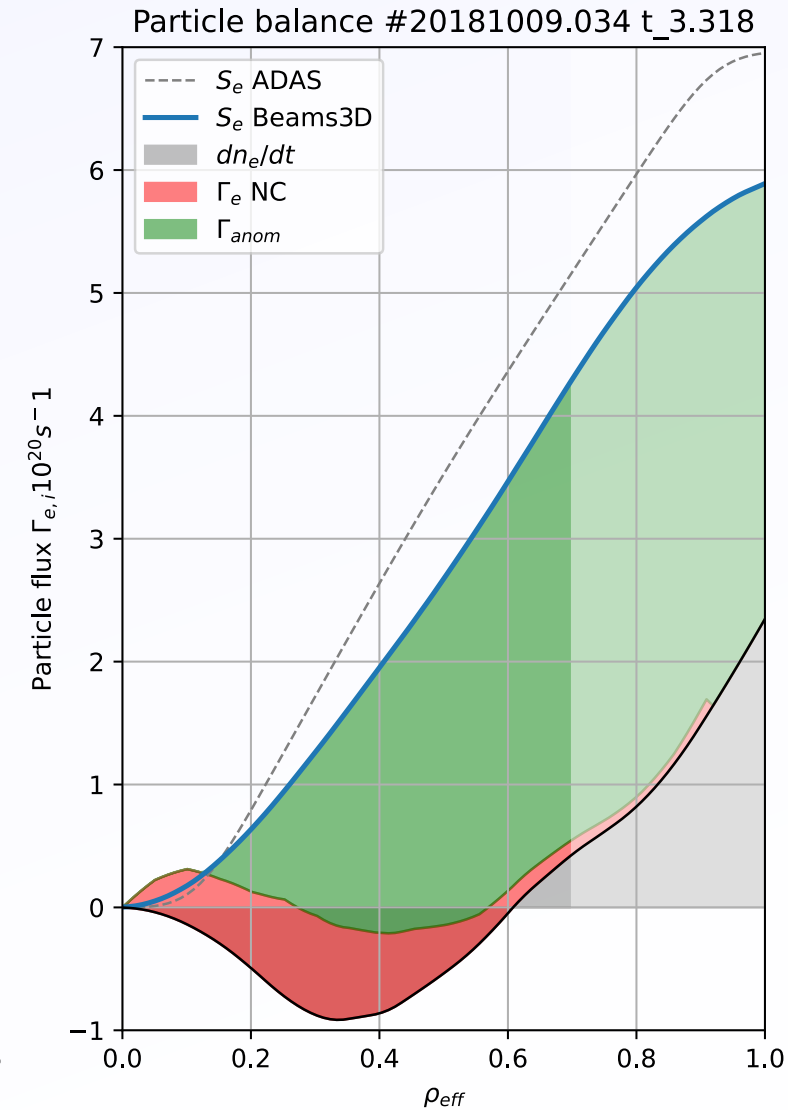
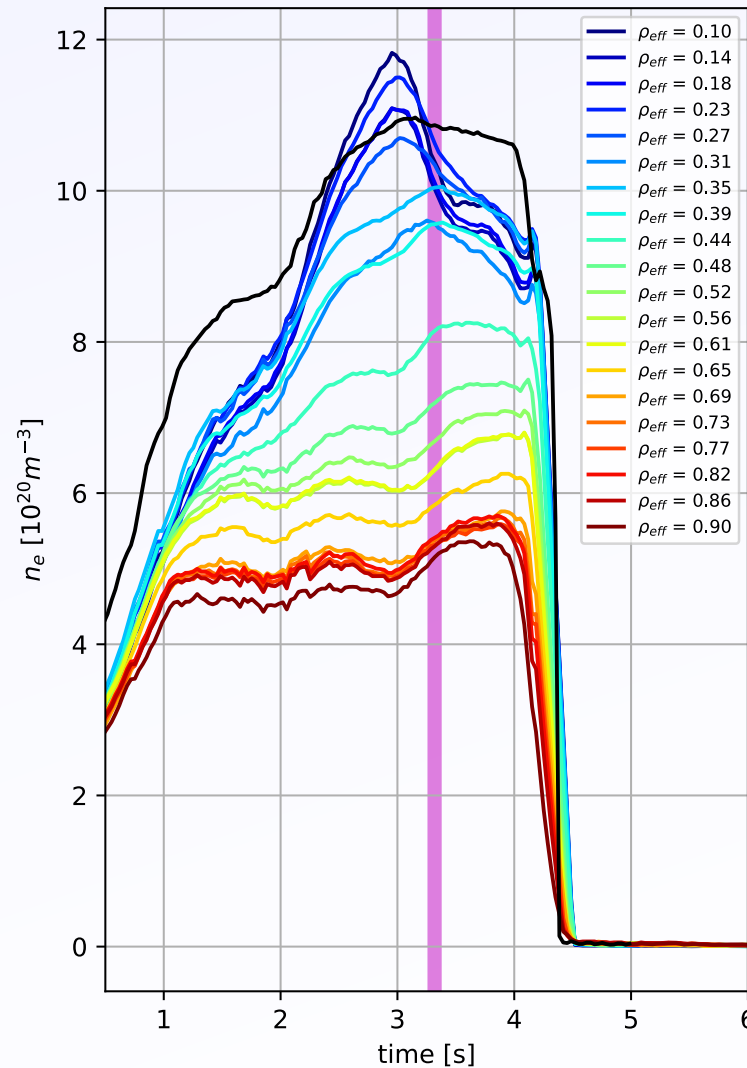
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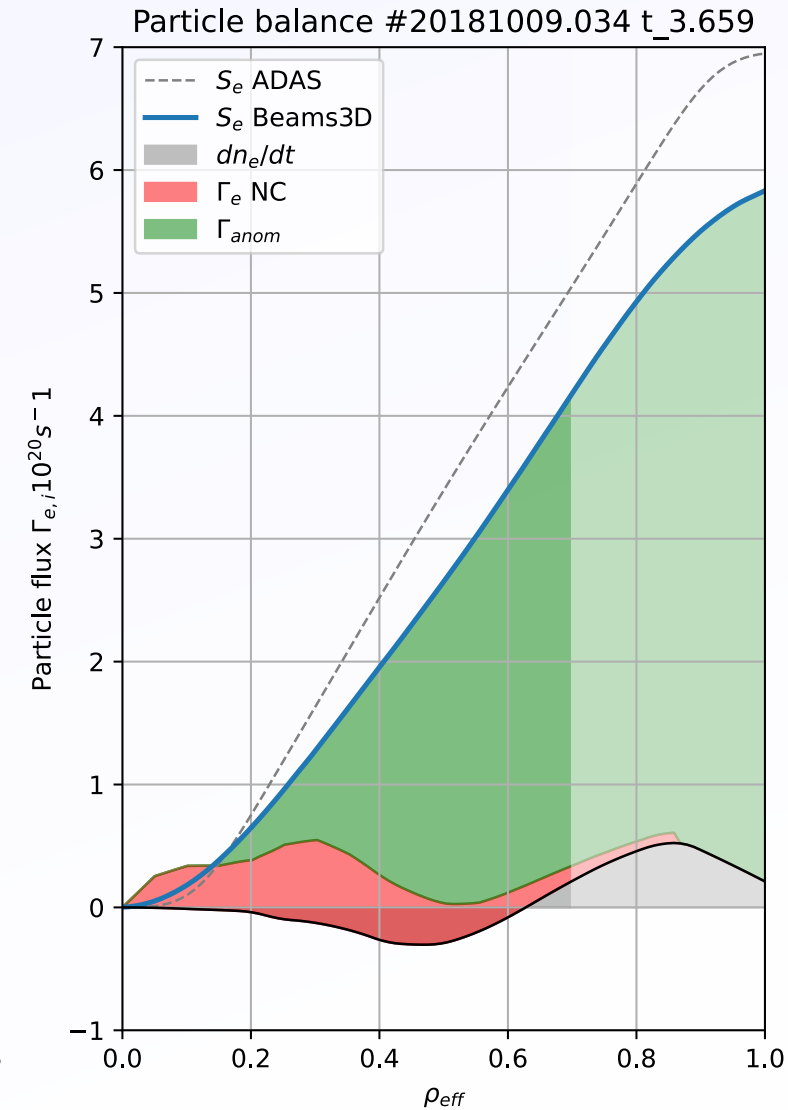
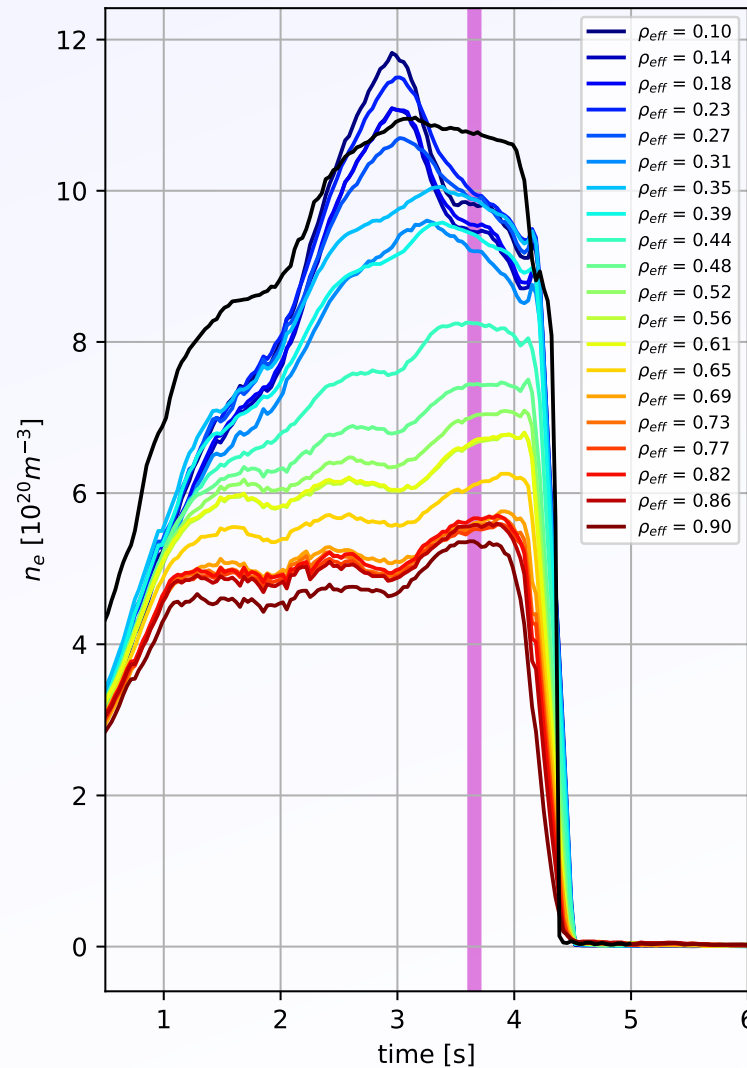
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#### 6) Density stabilises with balance of NC and anomolous in core. Strong anomolous at mid-radius to edge.

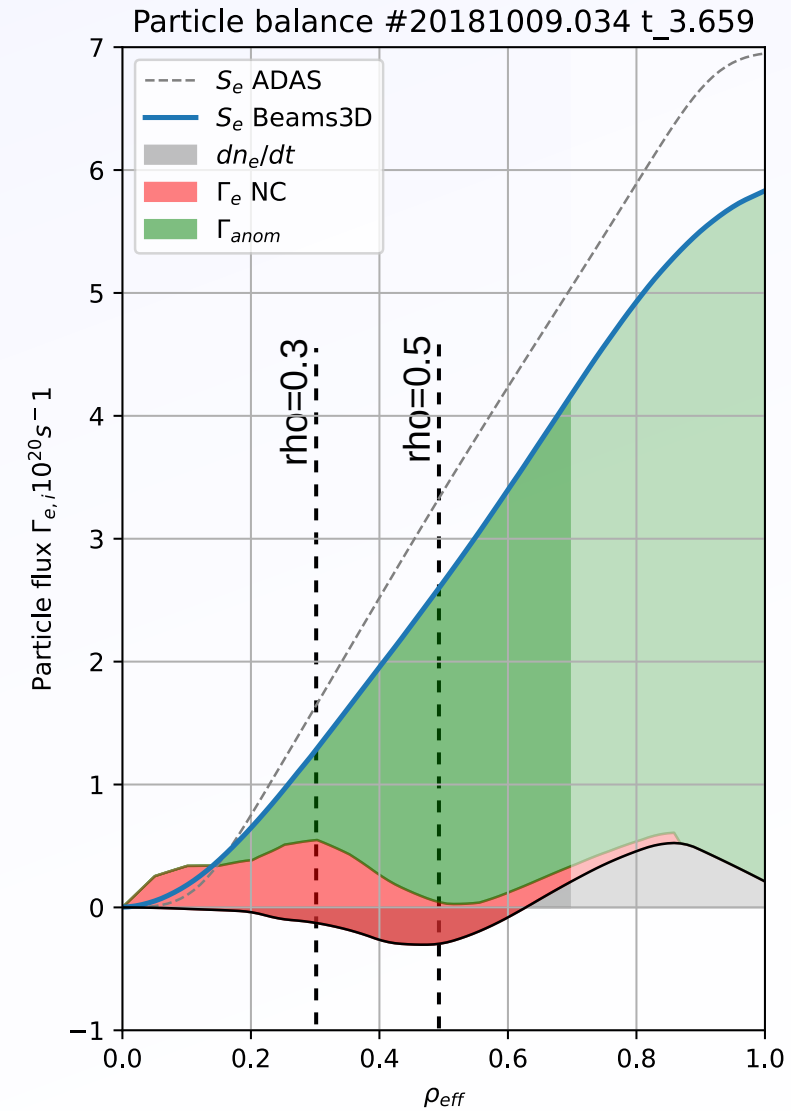
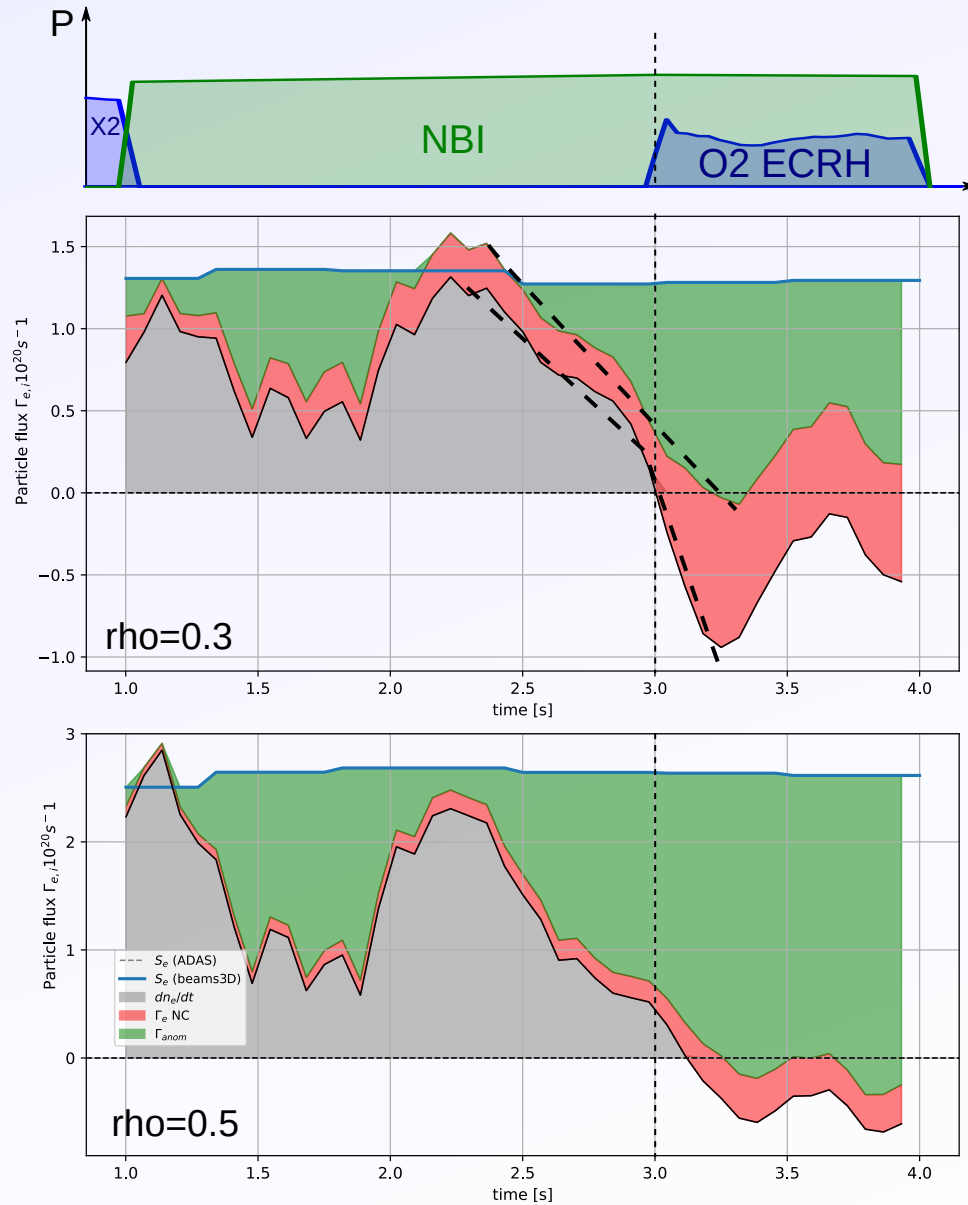




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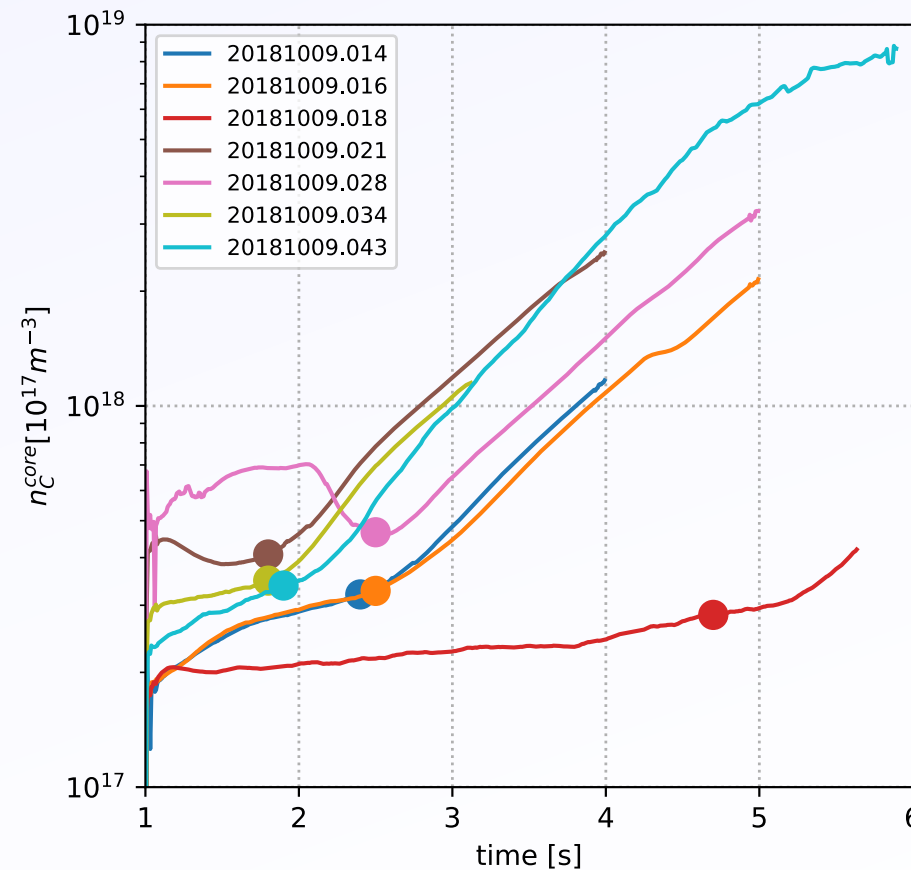
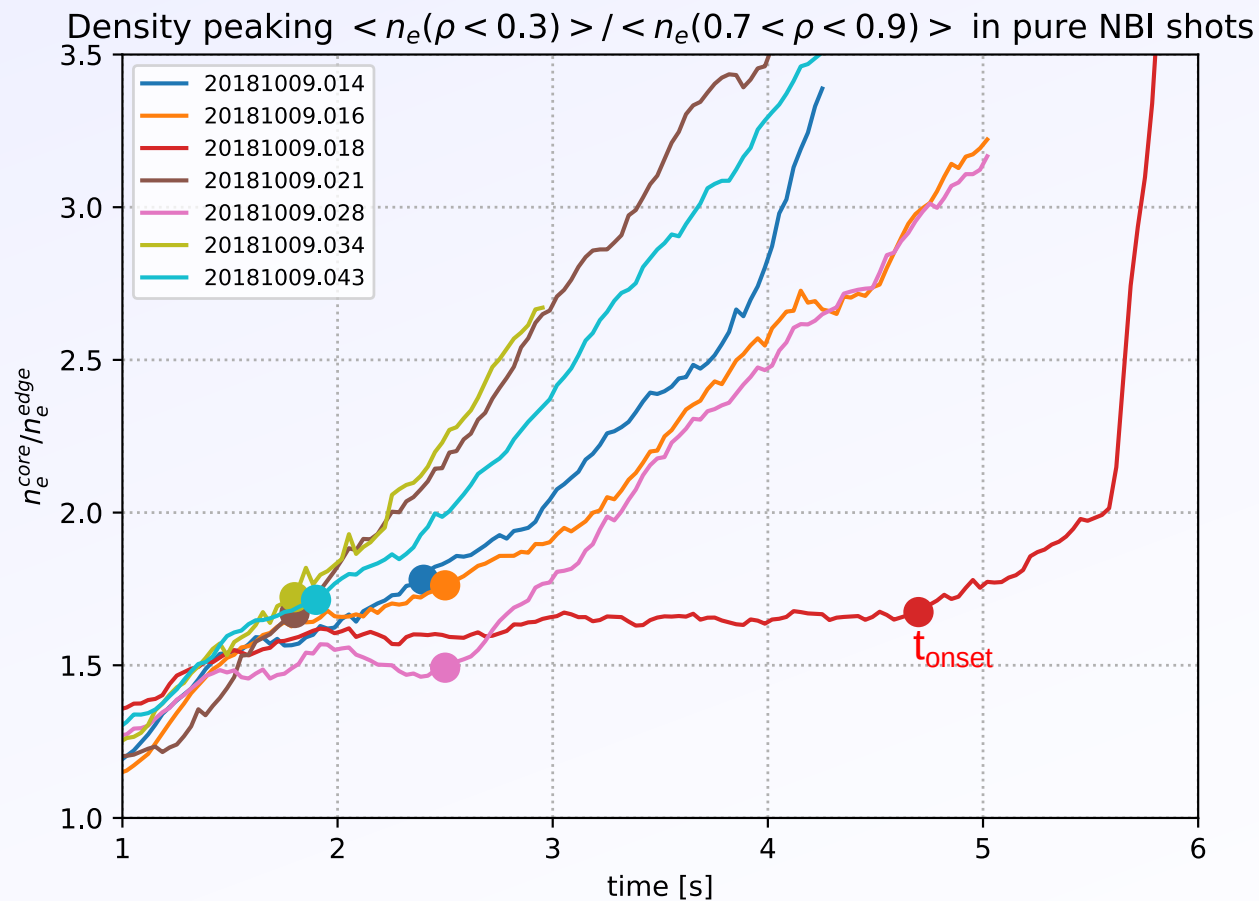
Particle balance temporal evolution:

At ECRH switch on, core  $d^2n_e/dt^2$  is consistent with increase of NC flux. i.e. new  $dn_e/dt$  matches  $\Gamma_{NC}$  with existing anomolous flux trajectory. ---> how quickly should turbulence react to profiles?



## Pure NBI - particle transport

The particle transport change appears in almost all NBI shots with  $P_{\text{ECRH}} < 1\text{MW}$ , at different on-set times. In some cases hard to see in  $n_e$ , but very obvious in  $\log(n_c)$  and almost coincident in time.



No change on any other signals at edge ( $T_e$ ,  $T_i$ ,  $H_\alpha$ ,  $P_{\text{rad}}$ )

In some cases  $n_e$  rises a little at all radii, in others the edge doesn't change.

Most consistent parameter at  $t_{\text{onset}}$  is  $a/L_n = 0.8 \pm 0.05$ , but this relies heavily the single red point (#018)

## Pure NBI - Species power balance

For power balance of individual species, we require the collisional power transfer  $P_{ei}$ :

$$P_{e-i} \approx 38 \cdot n_e^2 \cdot \frac{(T_e - T_i)}{T_e^{3/2}} \cdot \frac{Z}{A} \left[ \frac{\text{kW}}{\text{m}^3} \right]$$

At  $n_e \sim 10^{20} \text{ m}^{-3}$  and  $T \sim 1 \text{ keV}$  and integrating to mid radius:

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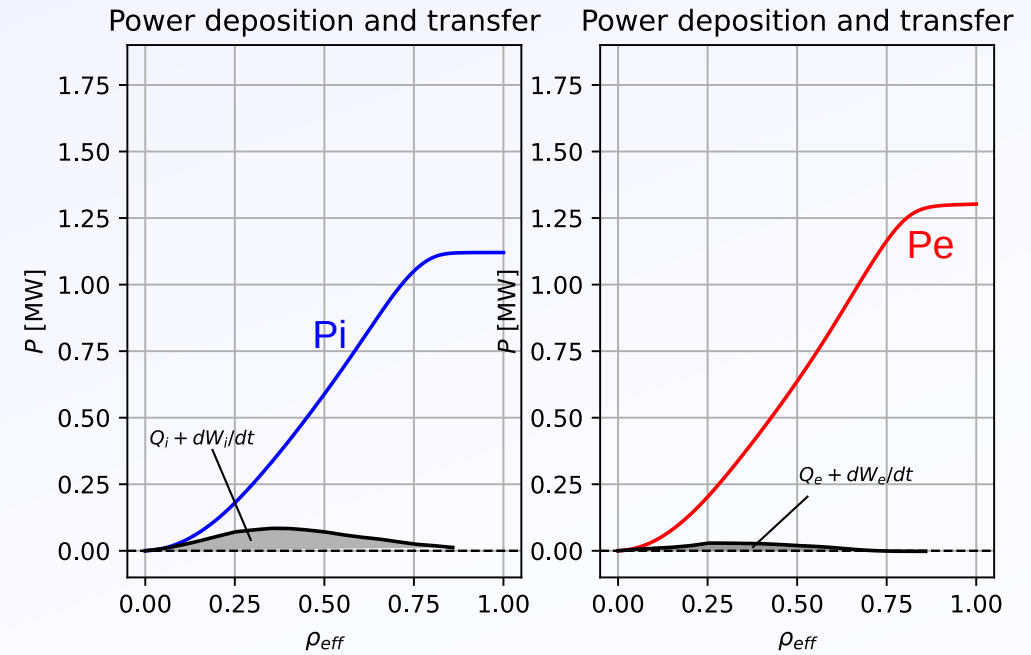
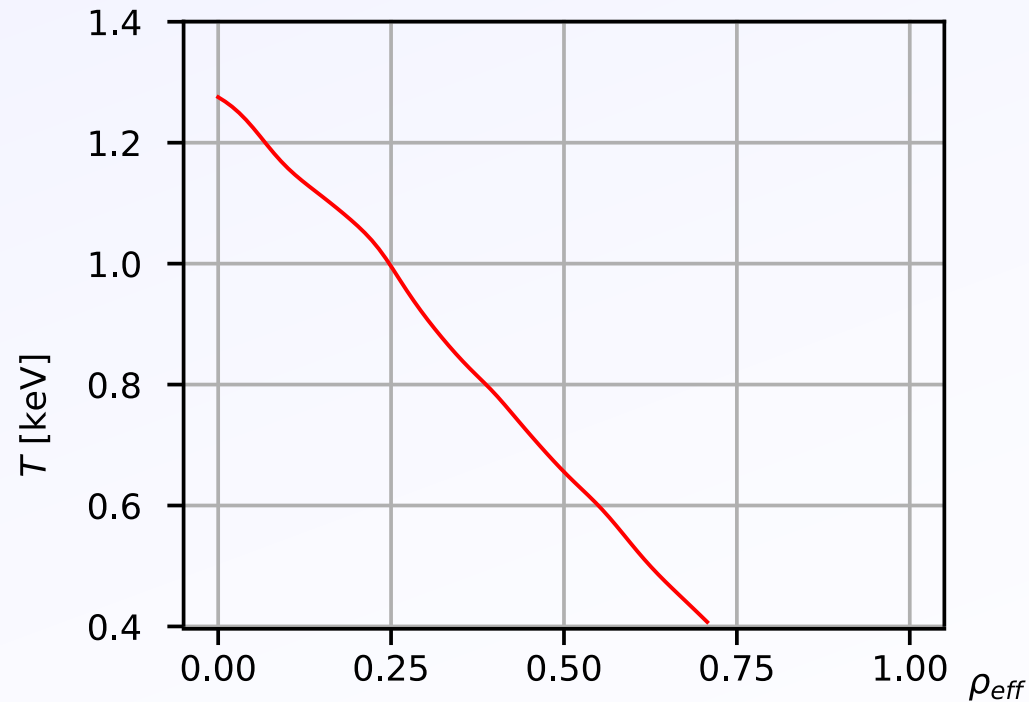
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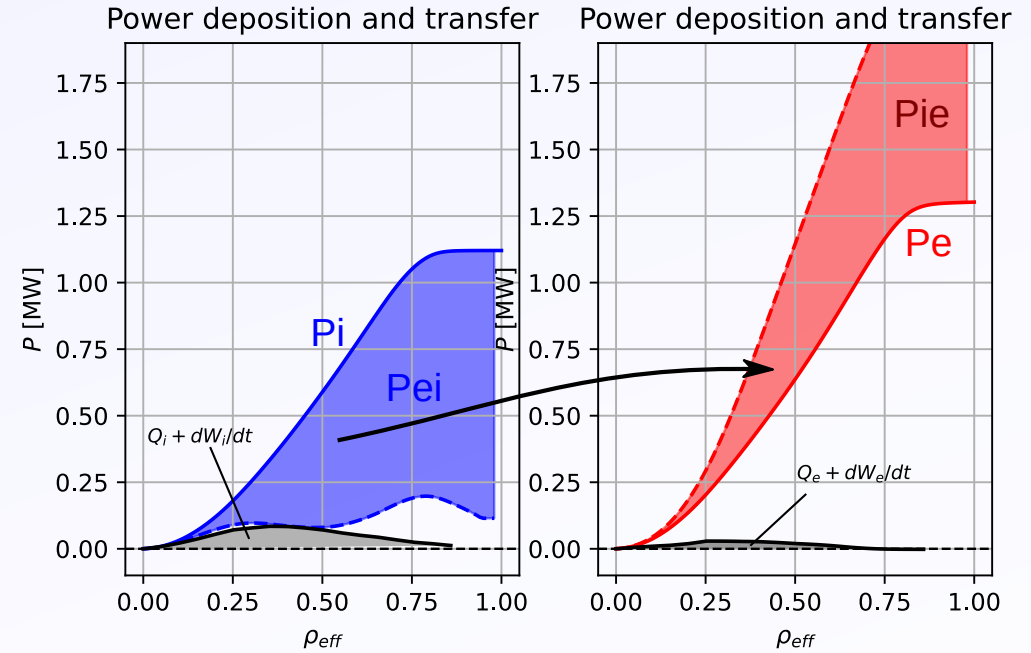
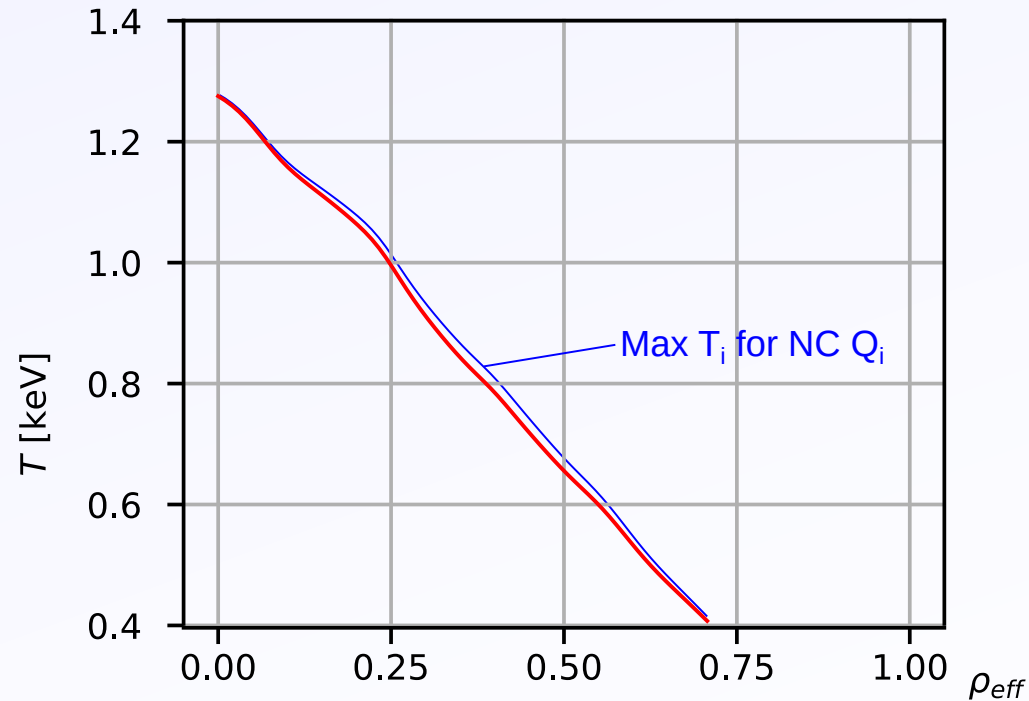
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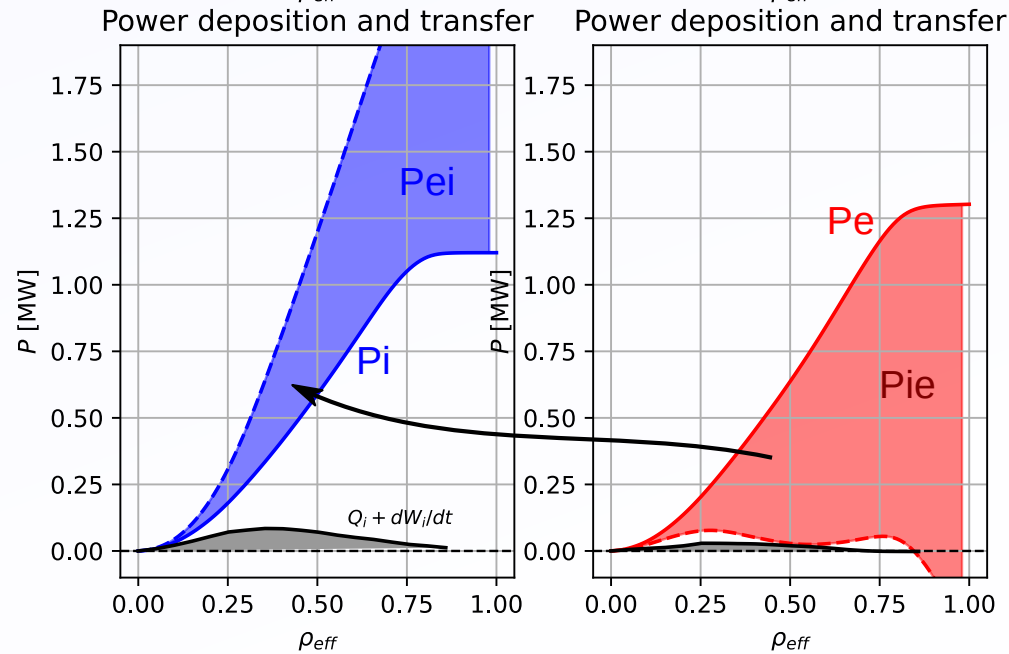
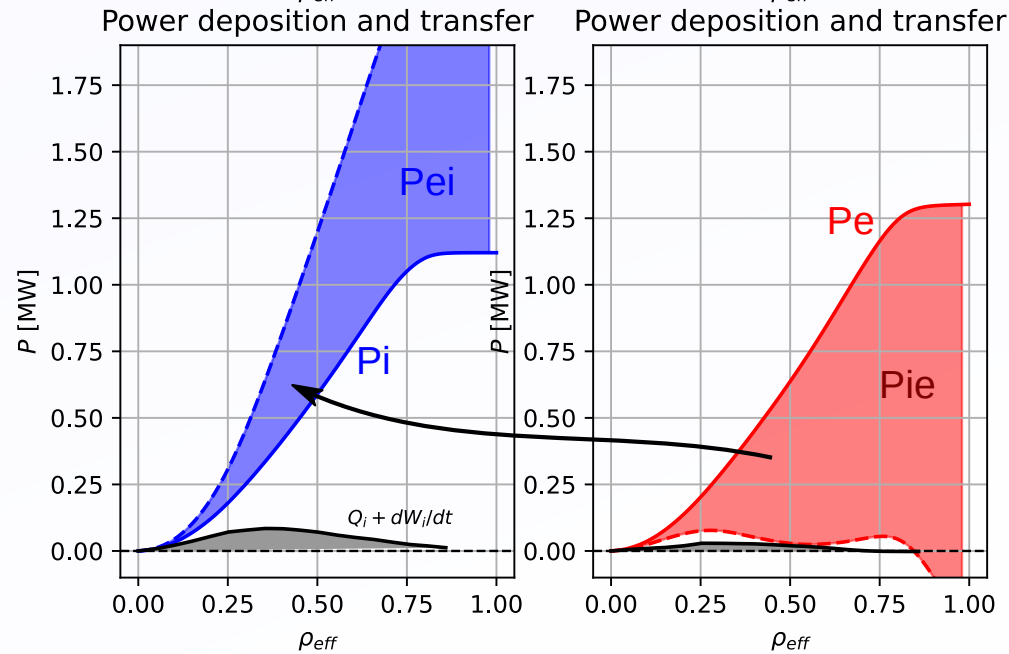
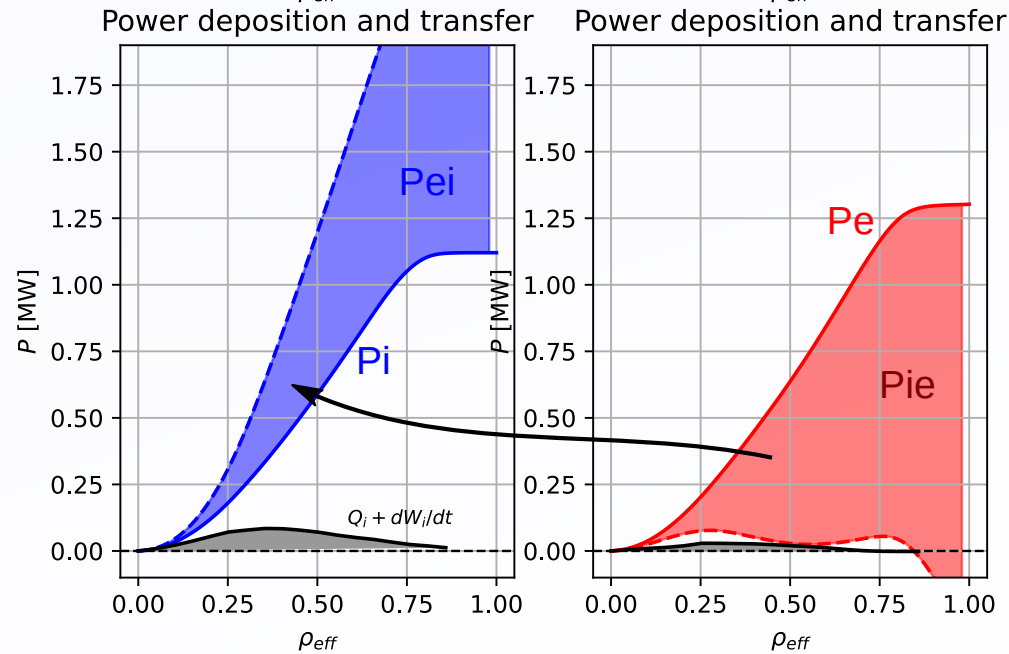
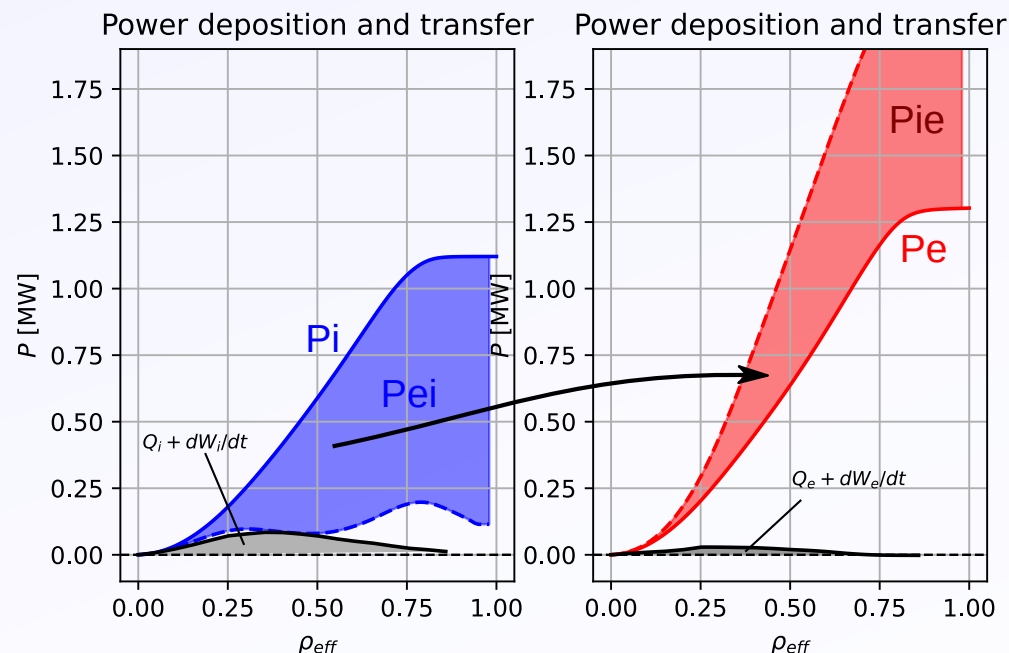
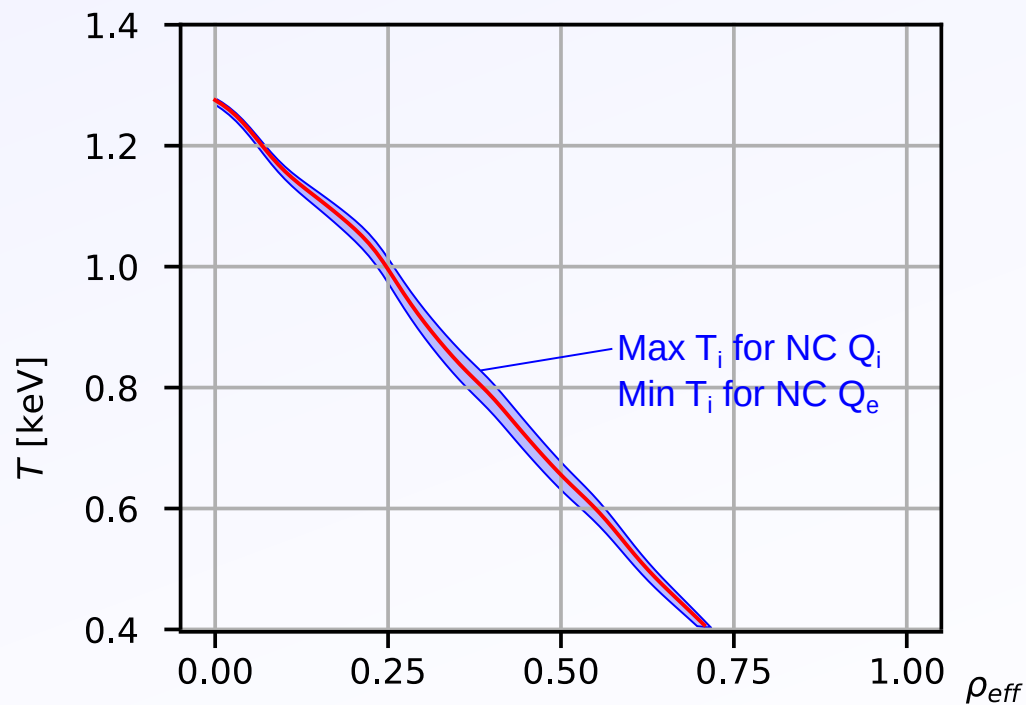
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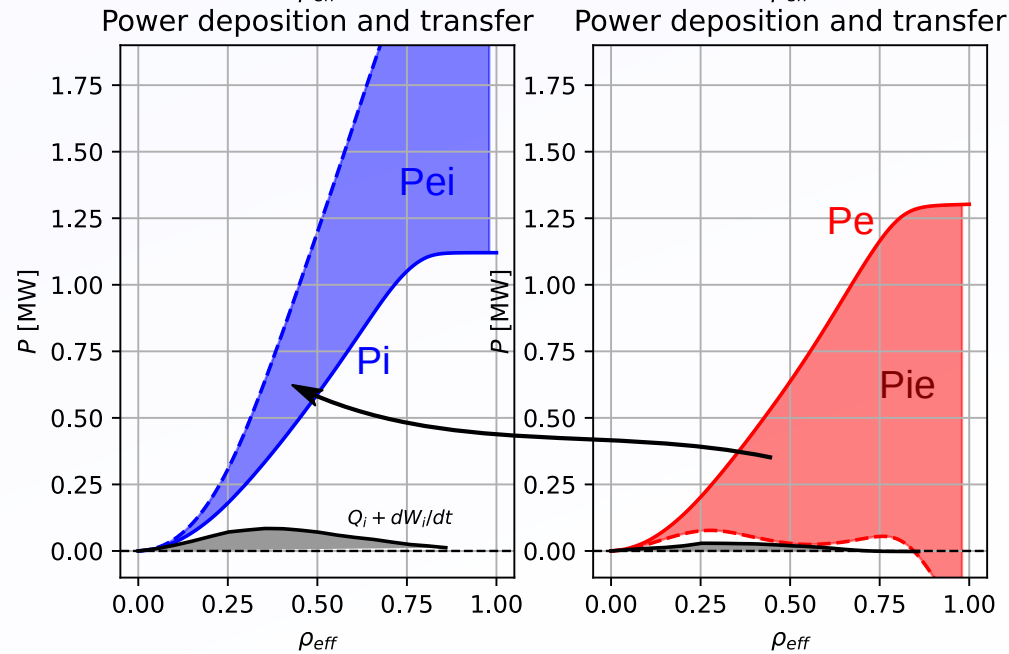
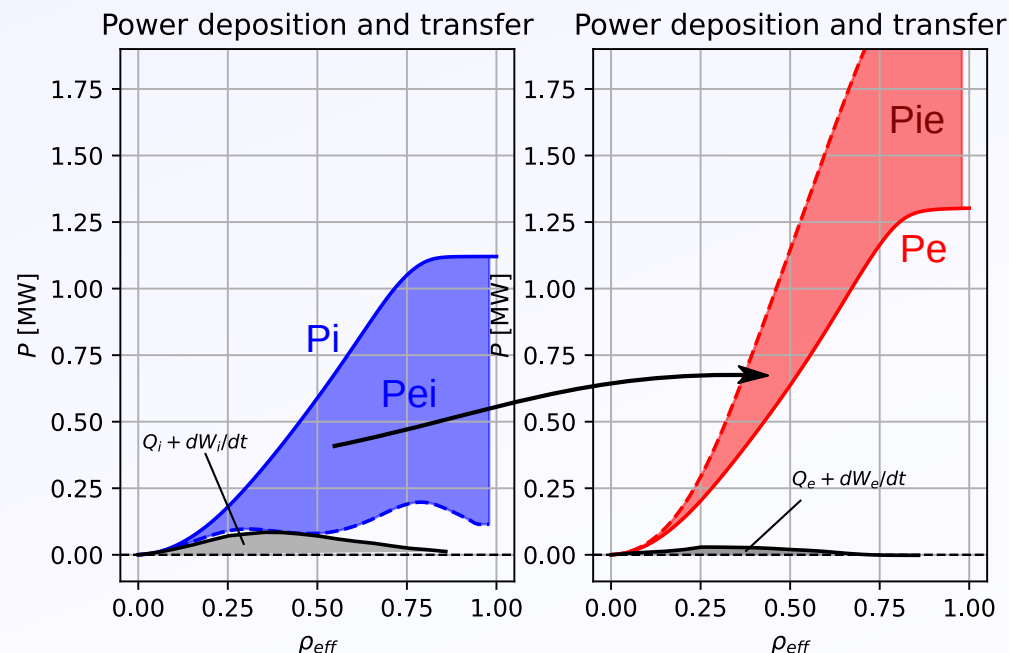
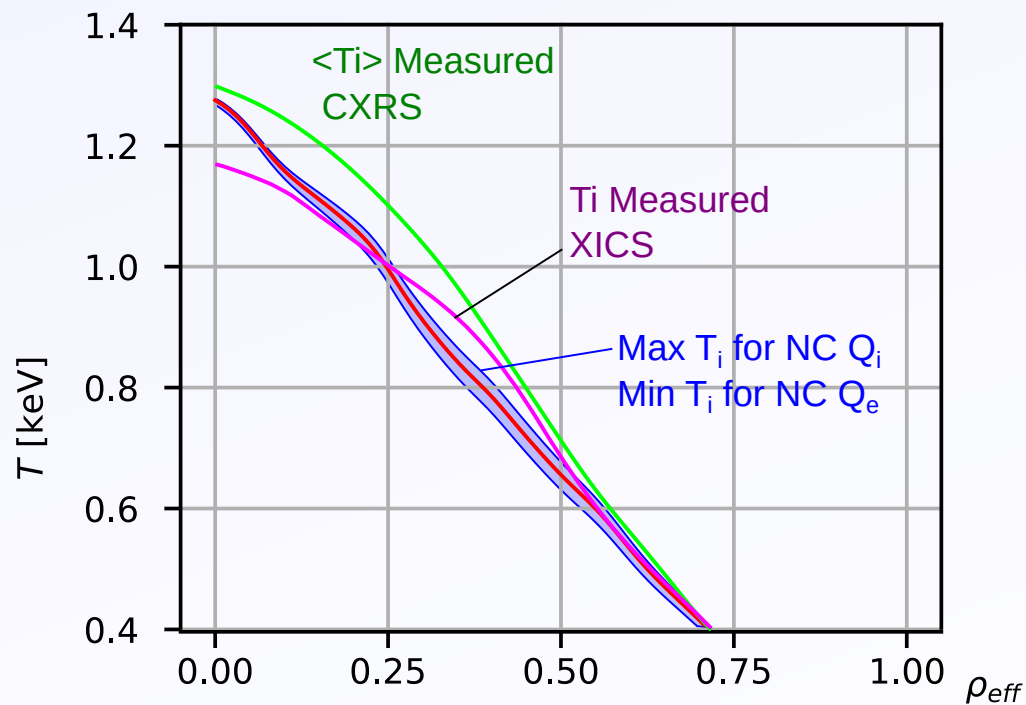
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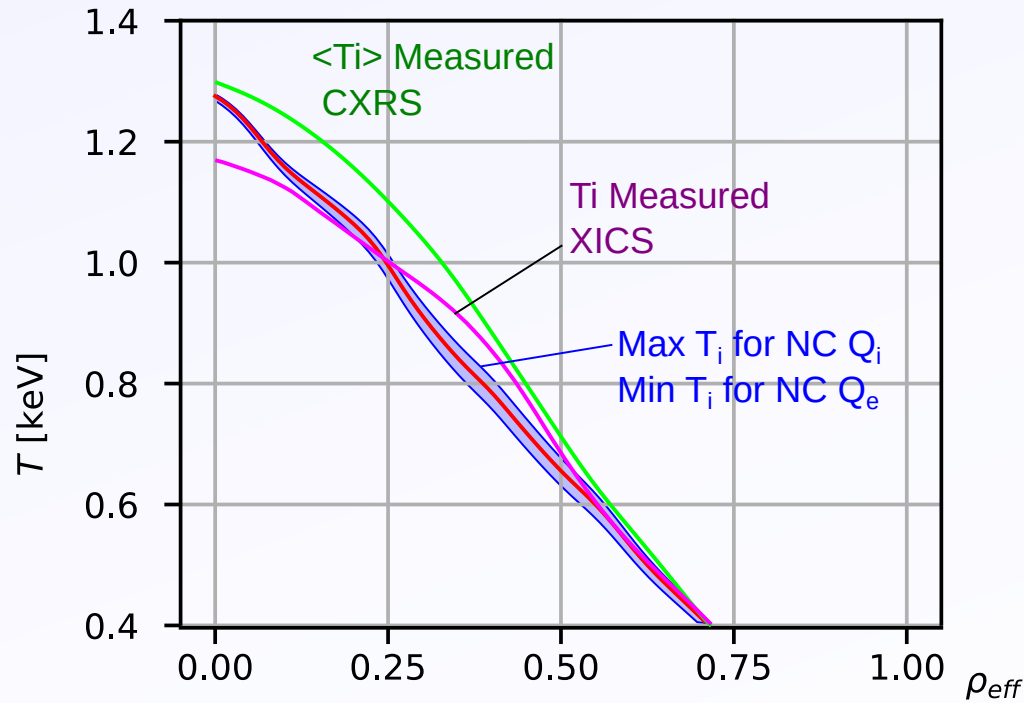


# Pure NBI - Species power balance

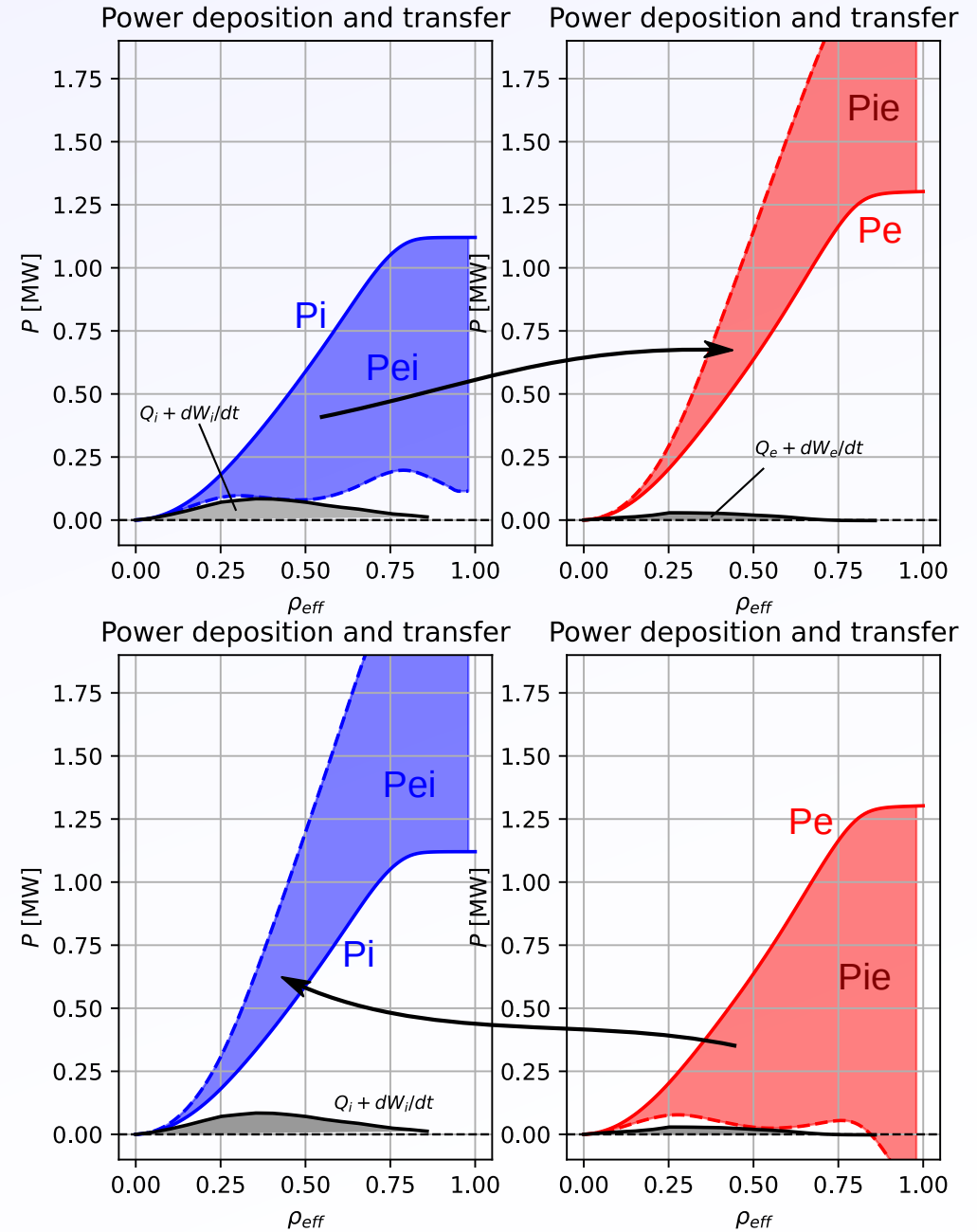
For power balance of individual species, we require the collisional power transfer  $P_{ei}$ :

$$P_{e-i} \approx 38 \cdot n_e^2 \cdot \frac{(T_e - T_i)}{T_e^{3/2}} \cdot \frac{Z}{A} \left[ \frac{kW}{m^3} \right]$$

At  $n_e \sim 10^{20} \text{ m}^{-3}$  and  $T \sim 1\text{keV}$  and integrating to mid radius:  
 $P_{e-i} \sim 2.6 \text{ MW}$  for every 100eV difference between  $T_e$  and  $T_i$ .  
 Only  $\sim 0.5\text{MW}$  is available from NBI at mid radius, so



Assumptions like  $T_e = T_i$  **are** assumptions about  $P_{ei}$  and lead to  $Q_e$  and  $Q_i$  values that **are not experimental quantities!**







## Te, Ti, Tz profiles during peaking

Can we recover  $P_{ei}$  by clever diagnostic analysis now we know  $|T_i - T_e|$  should be  $< \sim 50\text{eV}$

Temperature profiles available:

Thomson scattering  $T_e$ .

XICS Argon  $T_z$

XICS  $T_e$

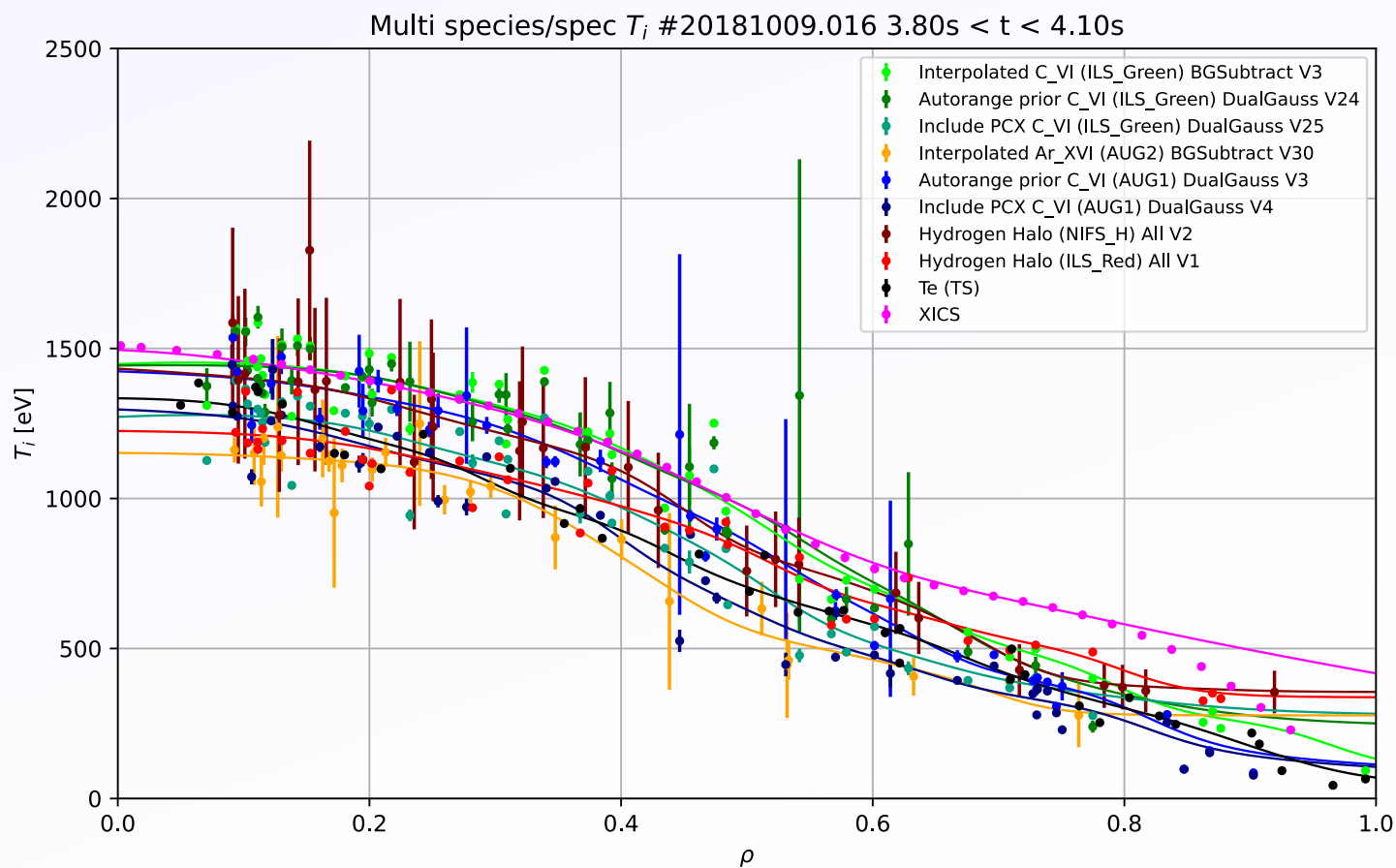
CXRS Hydrogen (Halo)  $T_i$

CXRS Carbon  $T_z$

CXRS Argon  $T_z$  (in 1 shot)

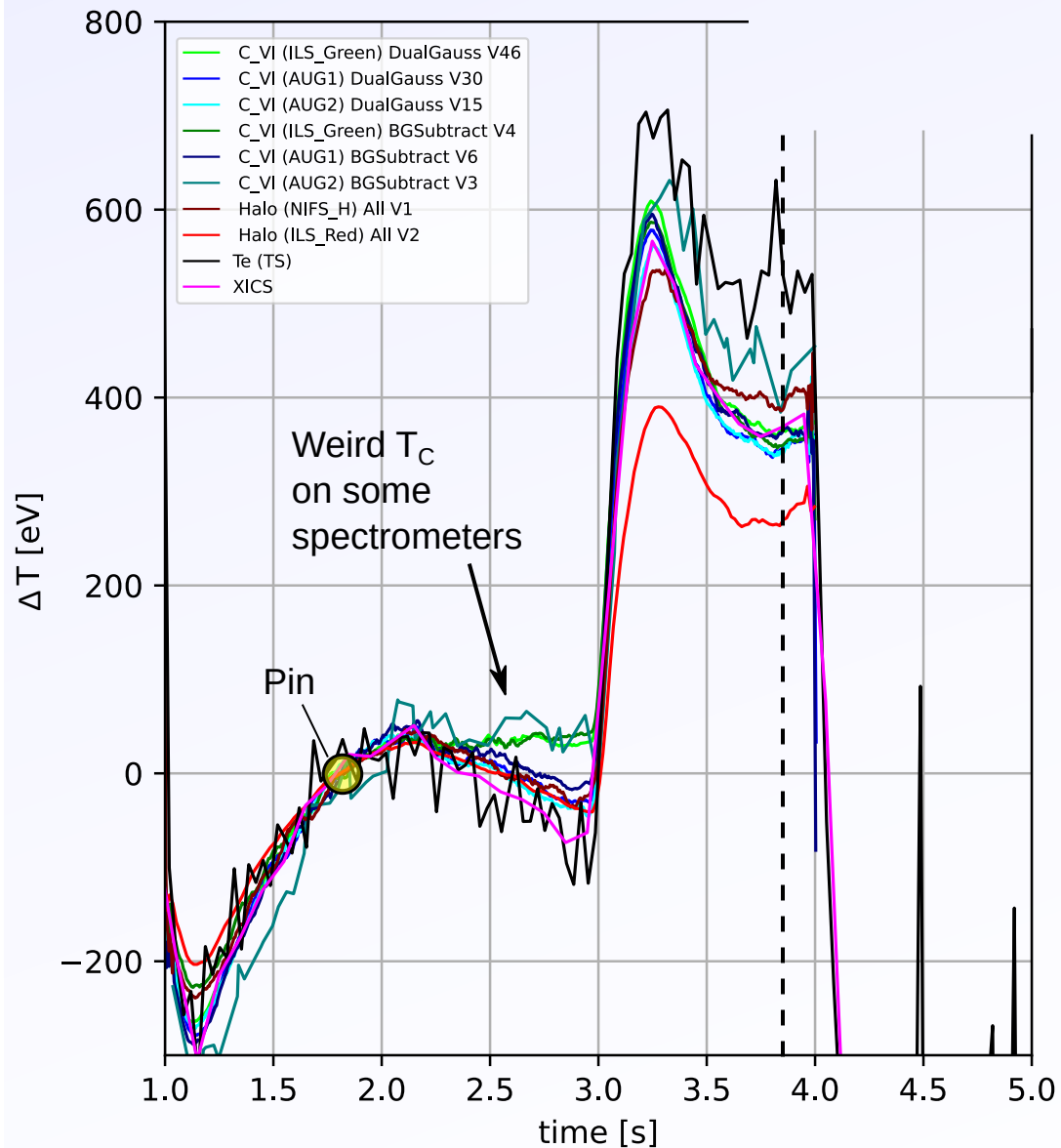
(All CXRS profiles corrected for fine structure, Zeeman and instrument function. Various methods to correct for PCX)

Generally these are mess of systematic errors:



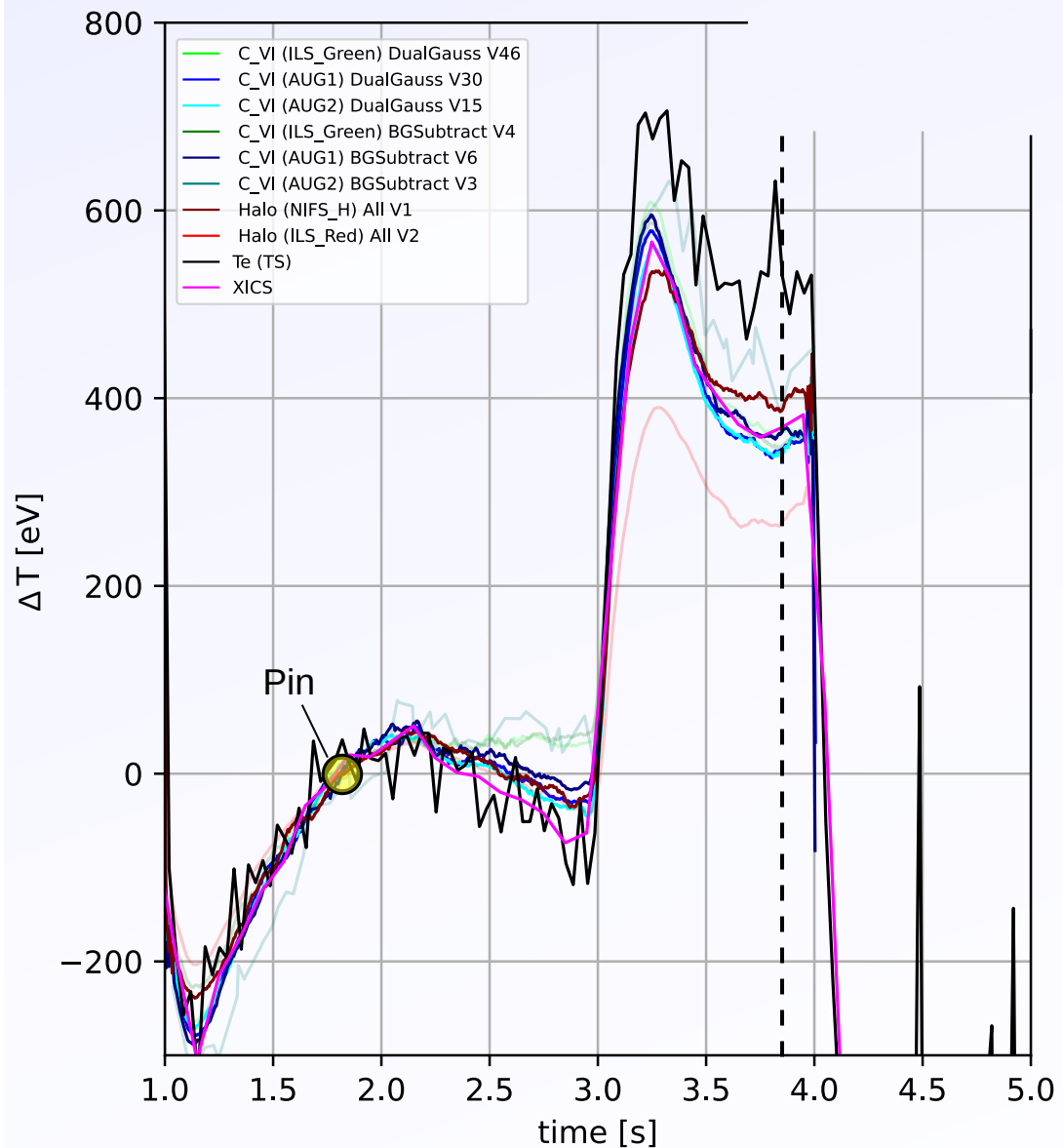


# $T_e, T_i, T_z$ profiles



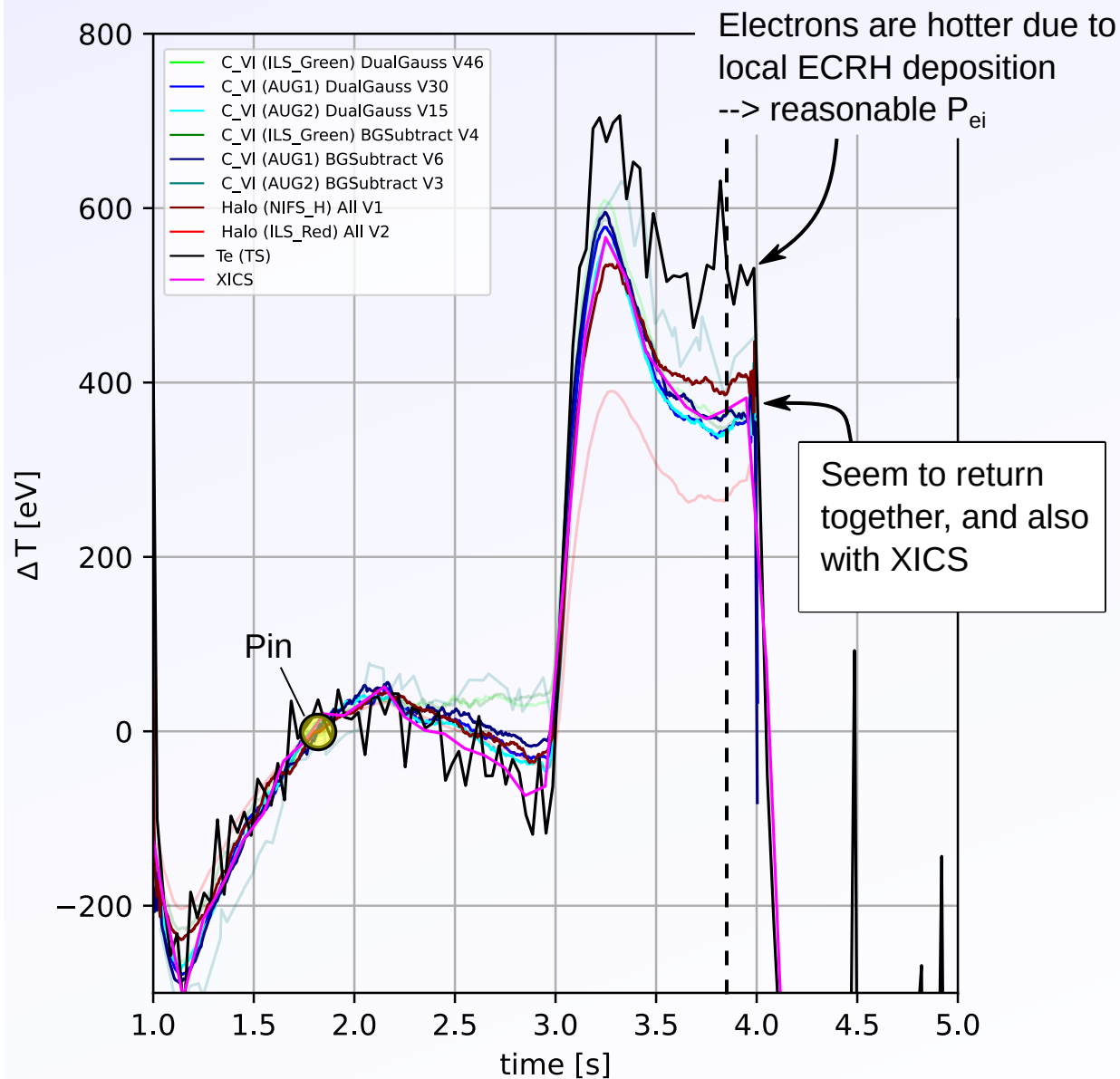


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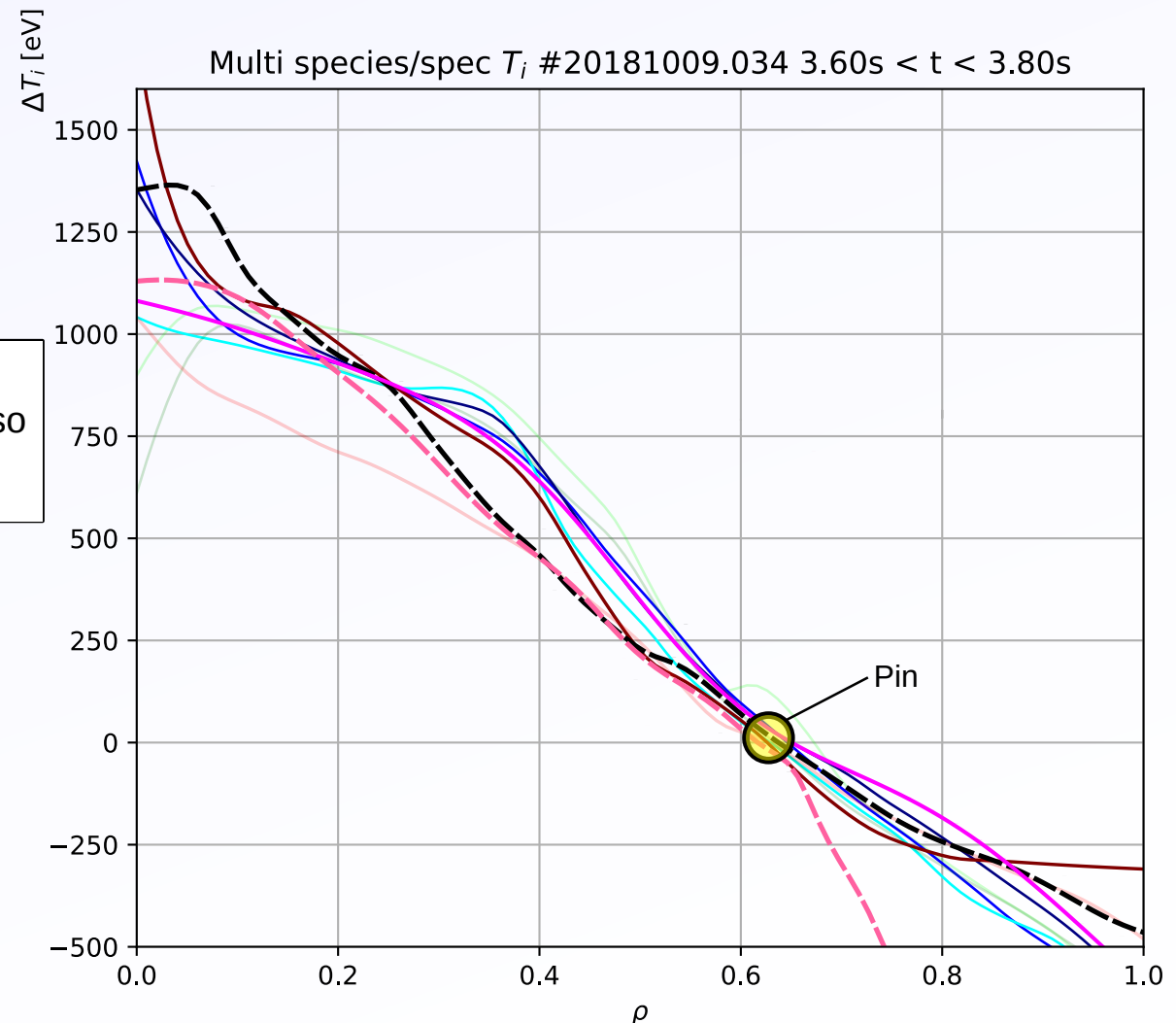
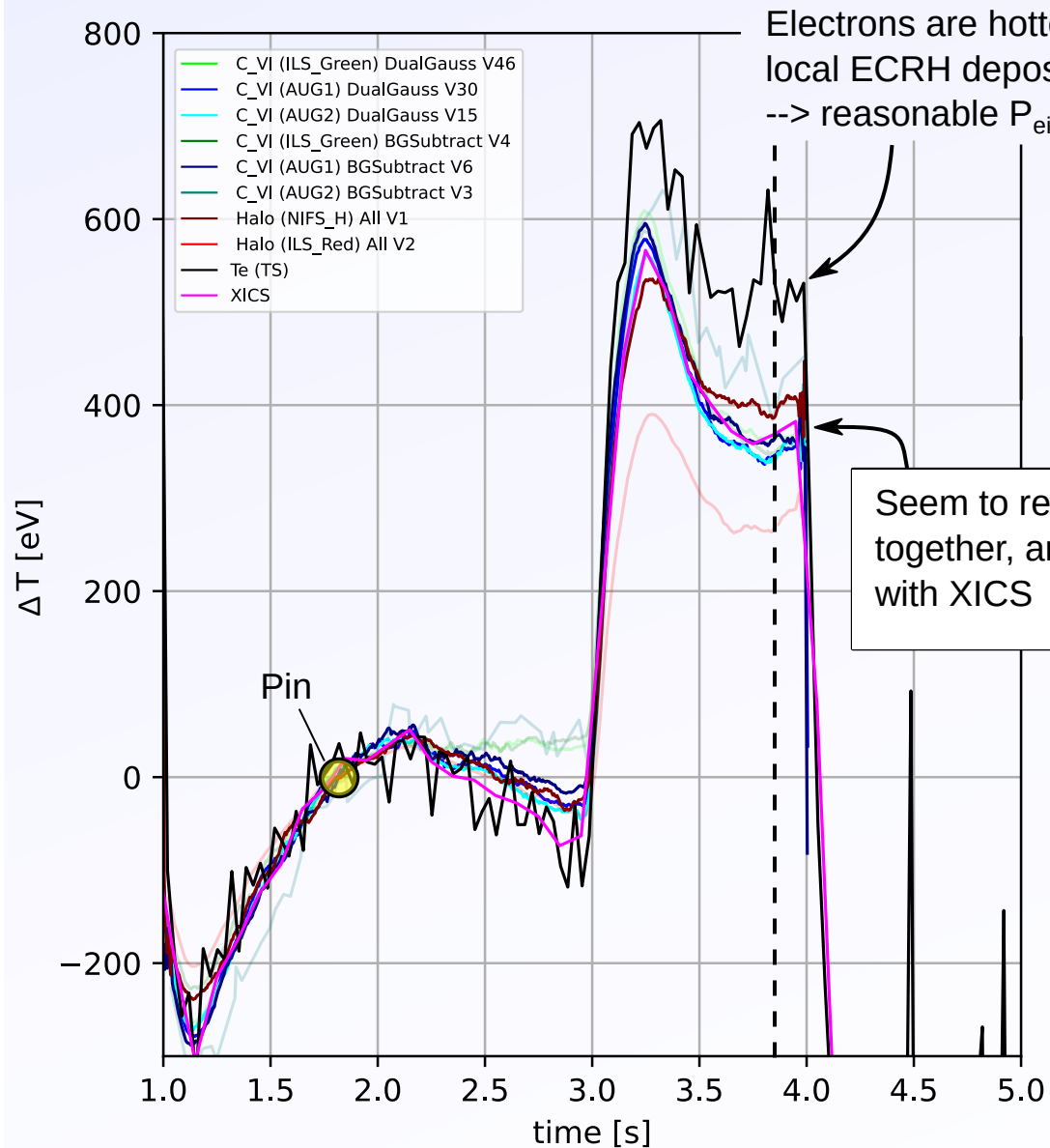


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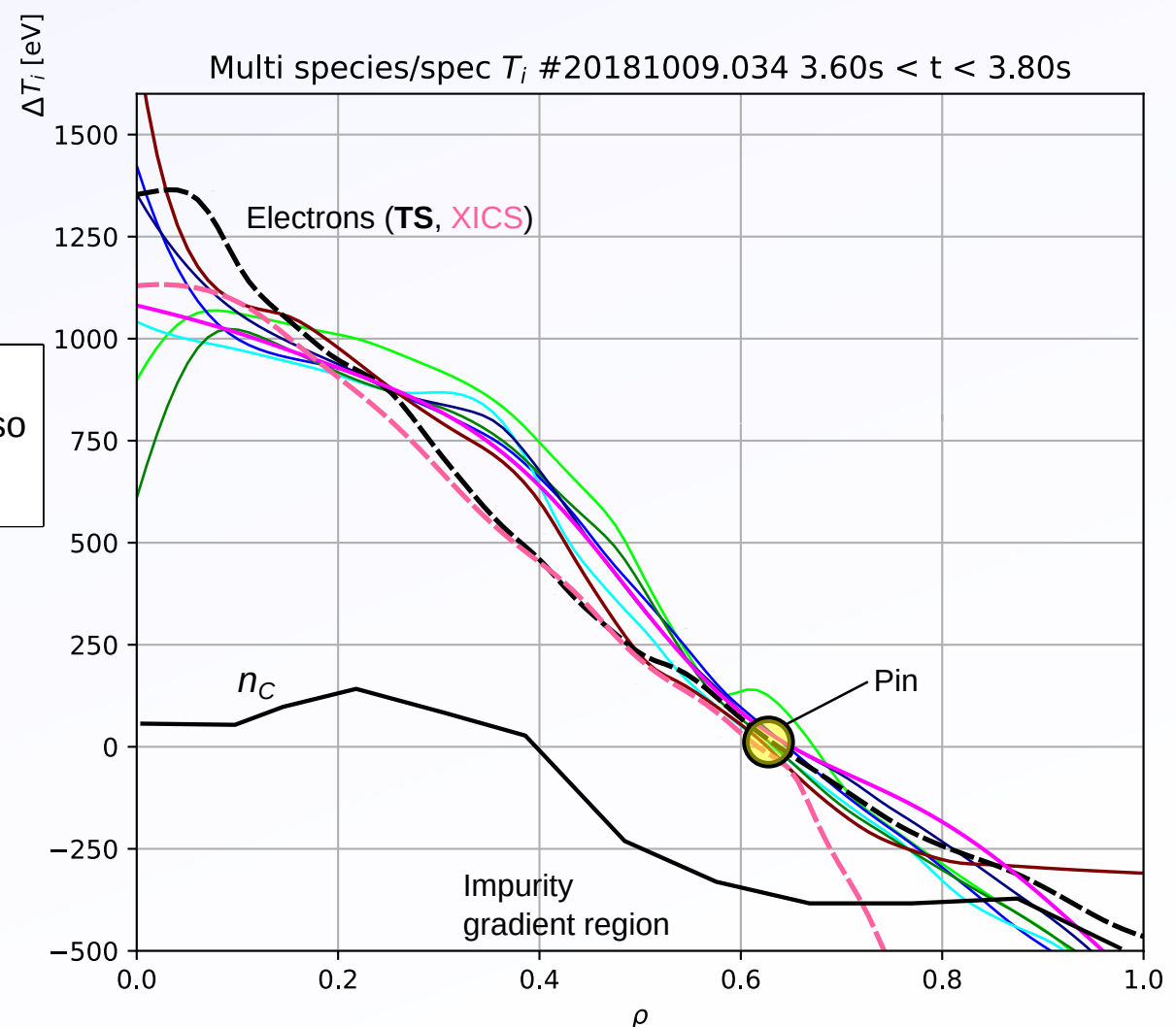
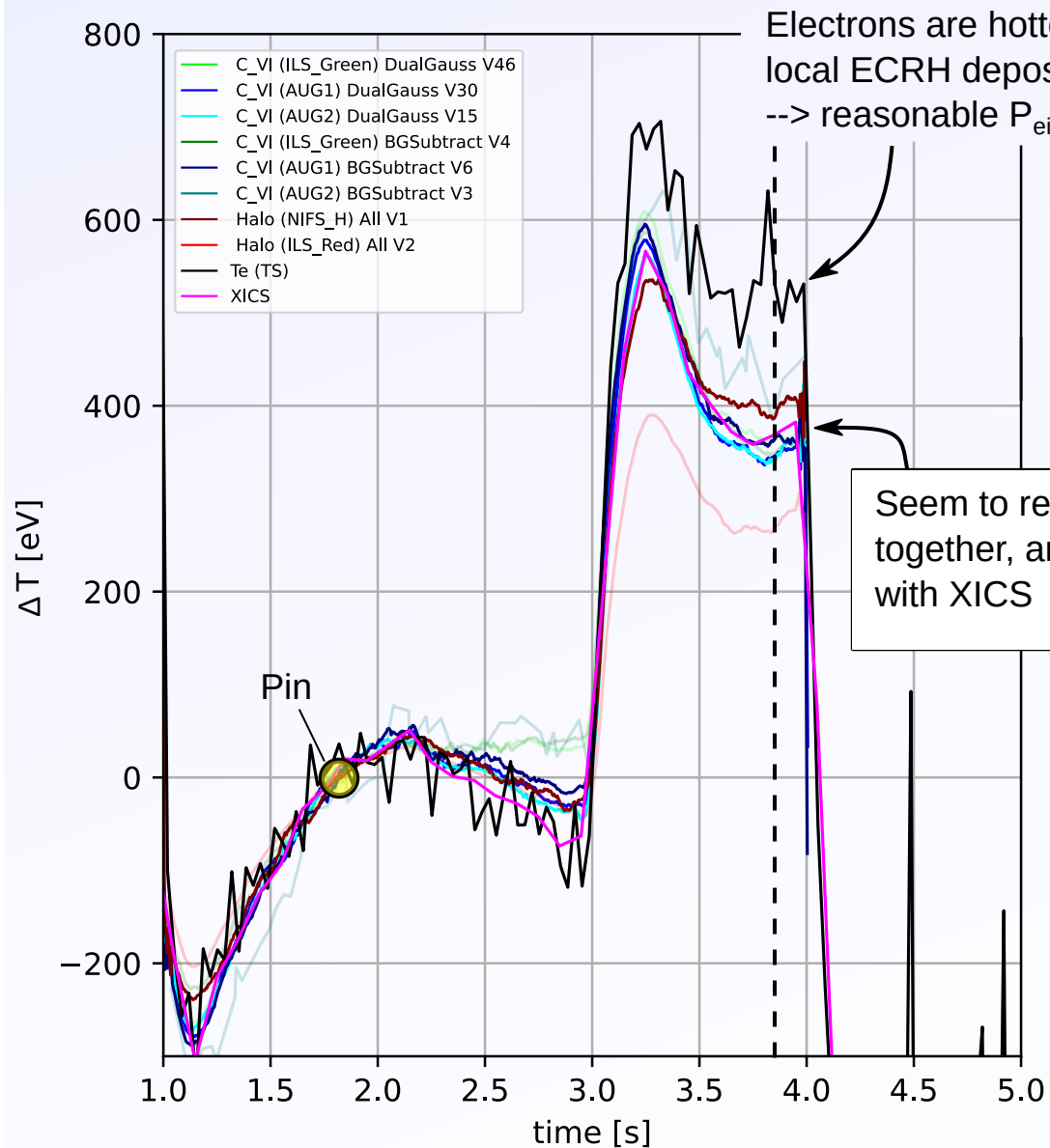


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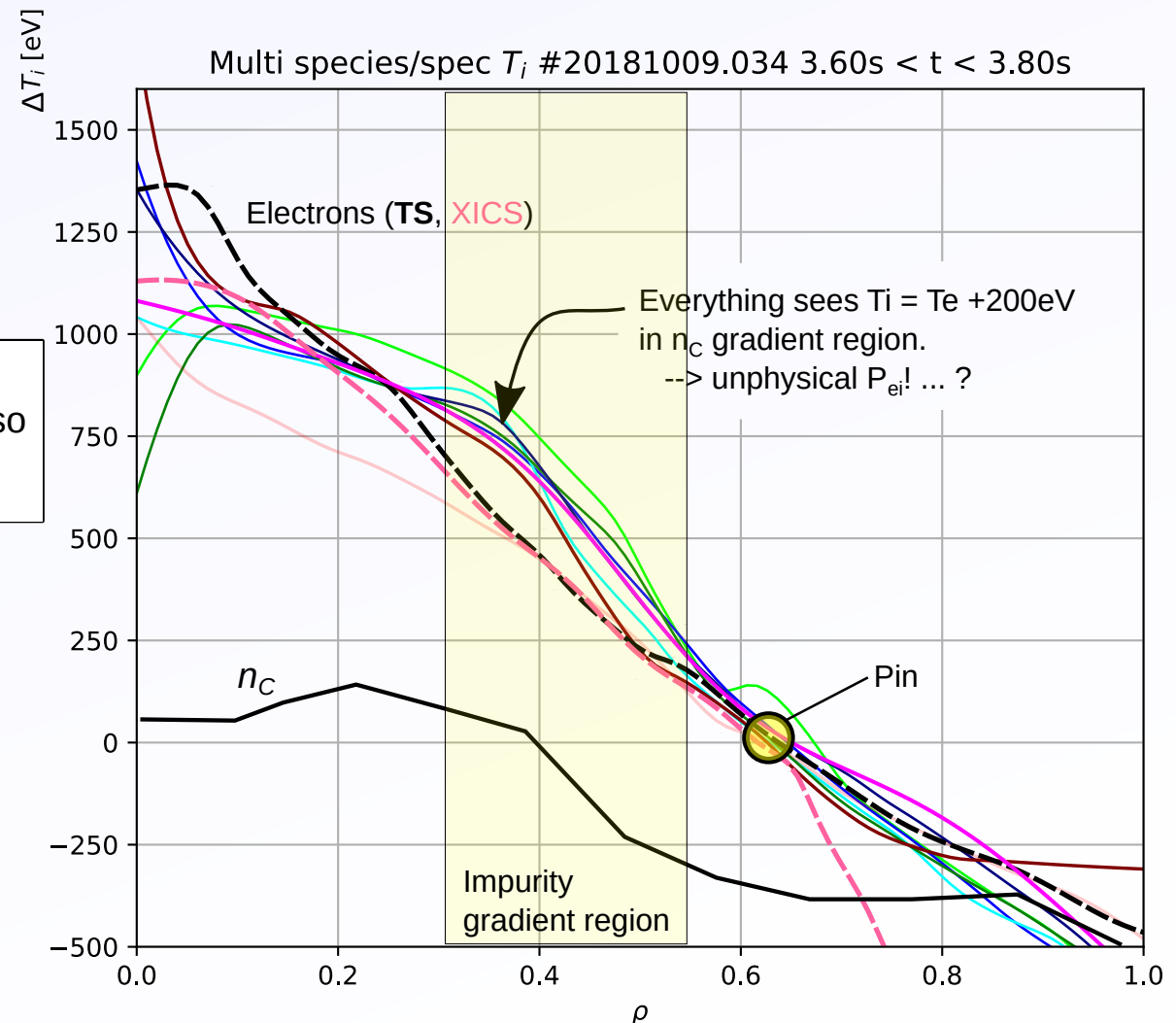
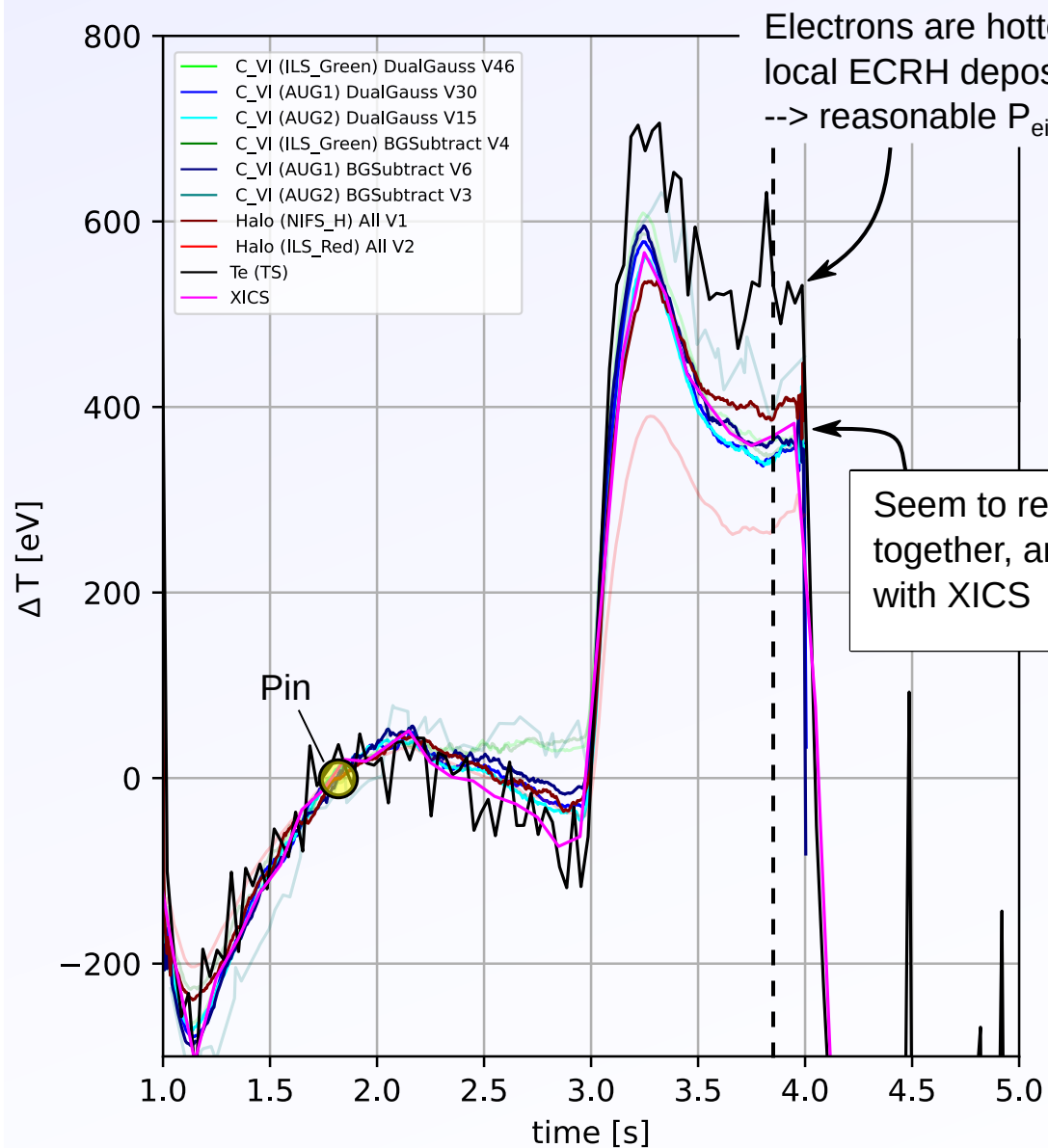
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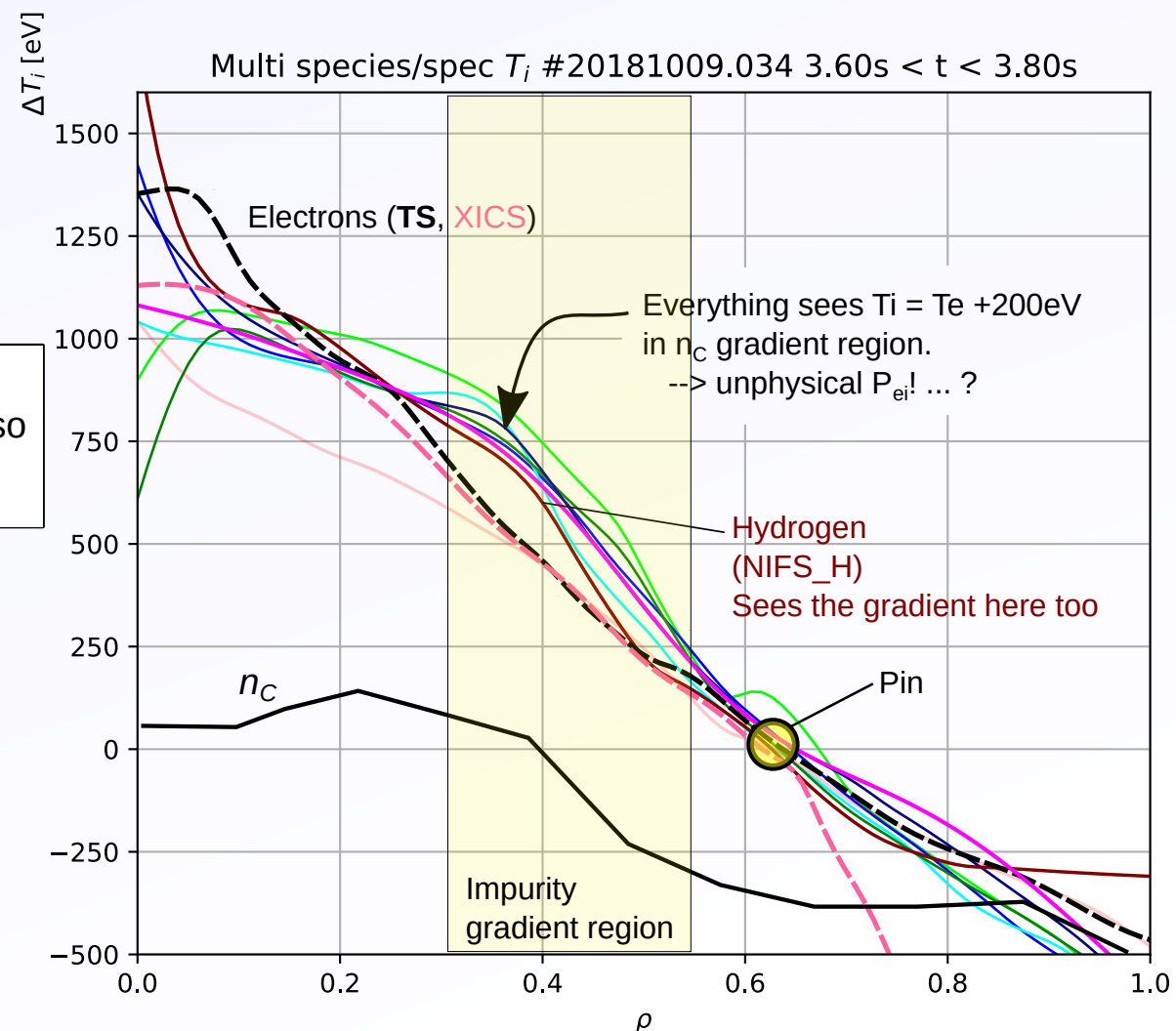
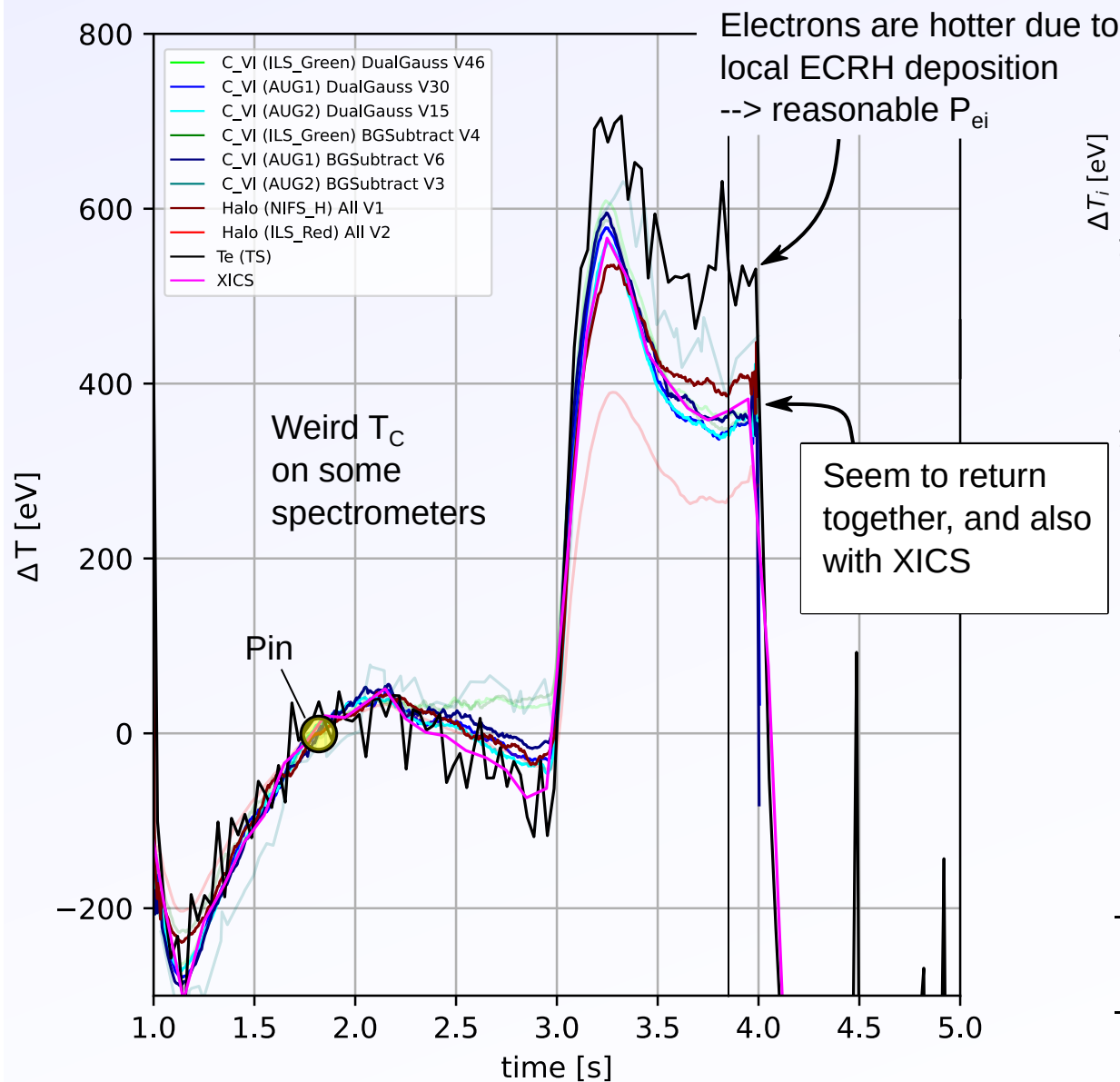
It seems like the  $T_z$  becomes much higher near very steep gradients.





$T_e, T_i, T_z$  profiles

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## Pure NBI - Species power balance

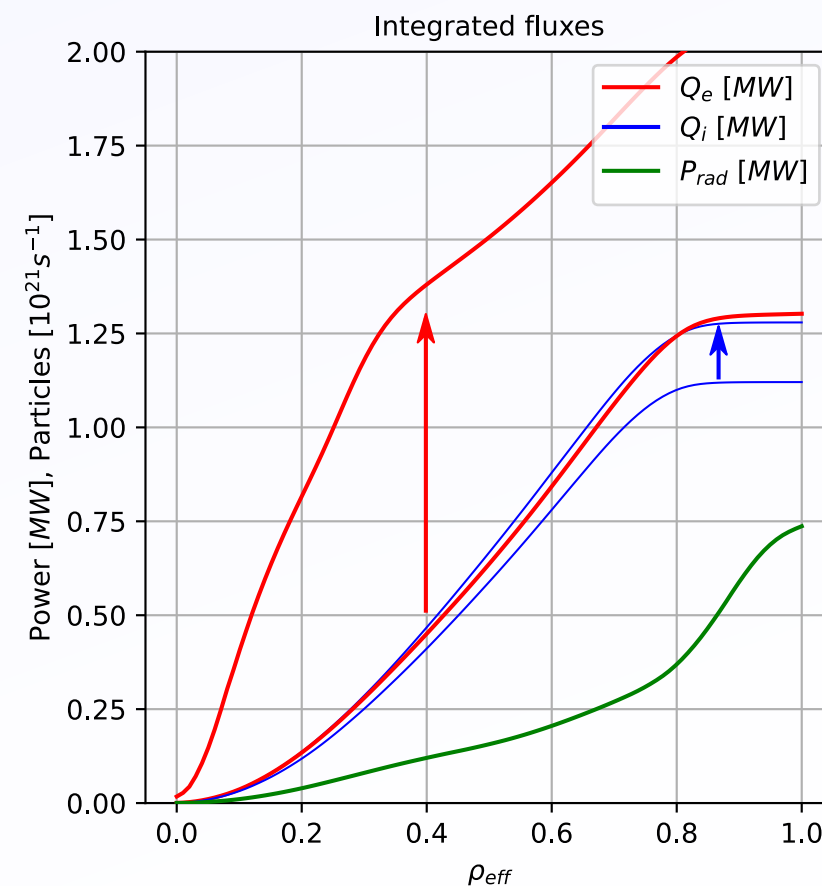
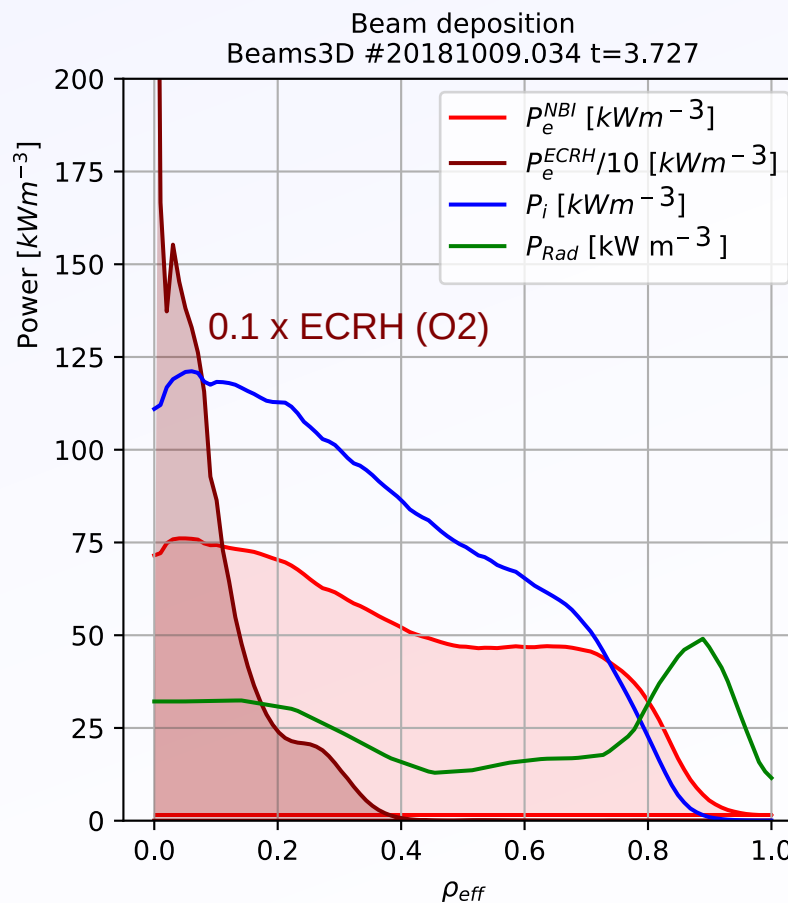
So what can we say?

- There are no believable cases where  $T_i < T_e$ , so we probably do not have neoclassical electrons.  
(This would fit with post-pellets plasmas, where we have near neoclassical ions but still very anomalous electrons)
- The ions could easily be completely neoclassical.
- There is no good reason to assume  $T_i = T_e$ . Any small differences in the heat transport would lead to differences building up radially until  $P_{ei}$  compensates it. To assume this, one would need to propose some mechanism to expect an exact  $Q_e = Q_e^{NBI} + Q_e^{ECRH}$  and  $Q_i = Q_i^{NBI}$  balance.

During the pure NBI phase, only 1.2MW of total power is available by  $\rho=0.5$ , so we can have max 45eV difference.

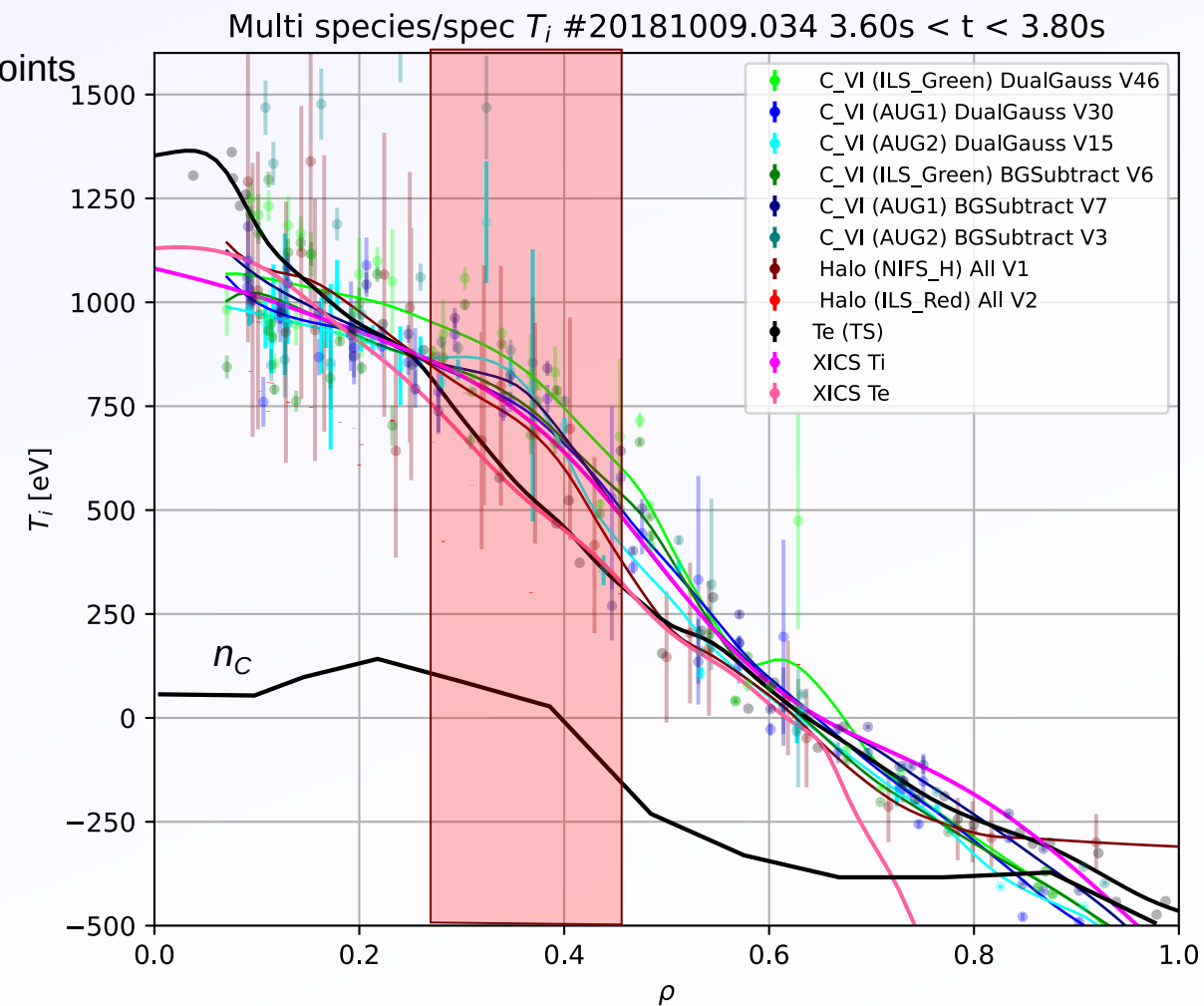
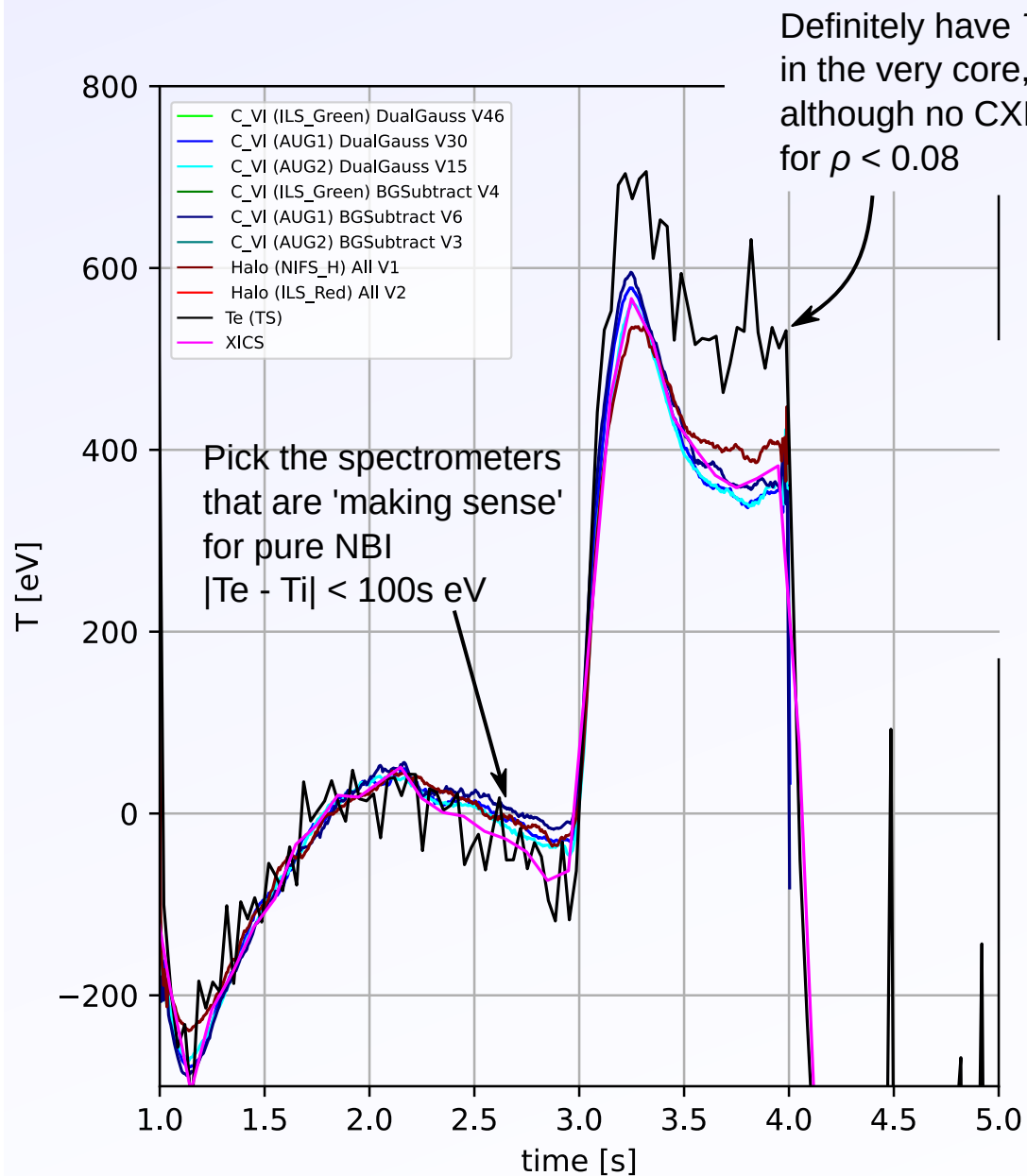
In the NBI+ECRH phase we get an additional 1MW of O2 ECRH power and can easily now have higher  $T_e$  in  $\rho < 0.2$ .

And in fact, the data tells us this...



## Pure NBI - Species power balance

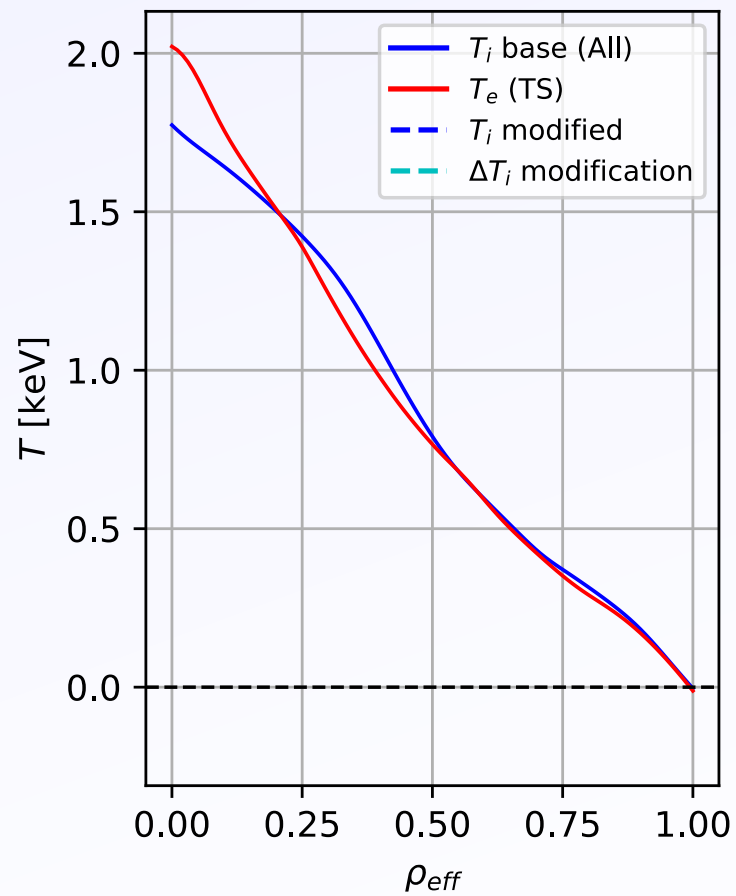
Examining the pinned measurement time traces again:



We're looking for an issue that affects both types of spectroscopy (CXRS + XICS), or something spatial (e.g. equilibrium). So we might 'mistrust' the gradient region where  $T_i \gg T_e$

## Pure NBI - Species power balance

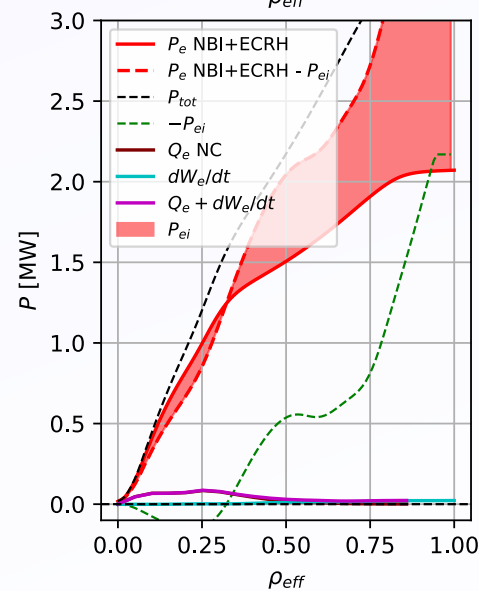
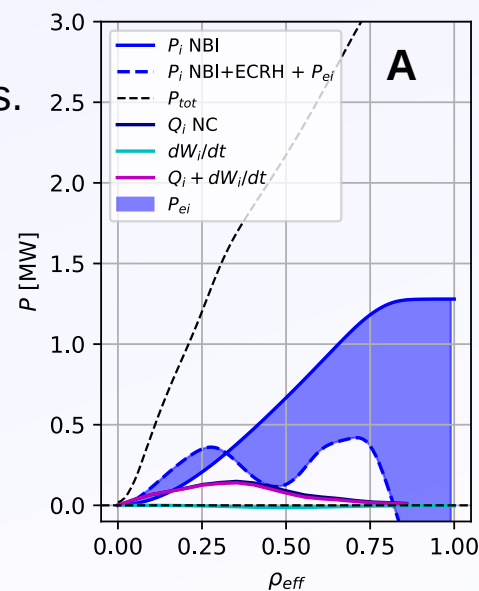
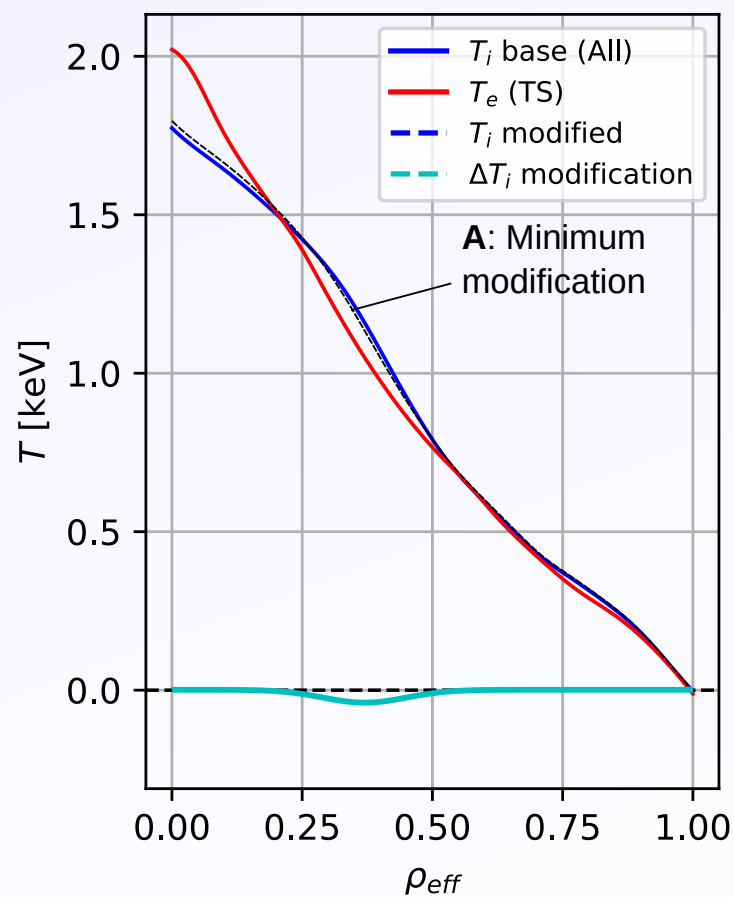
We can go back to the original data (no pinning, no adjustment) and just average everything:  
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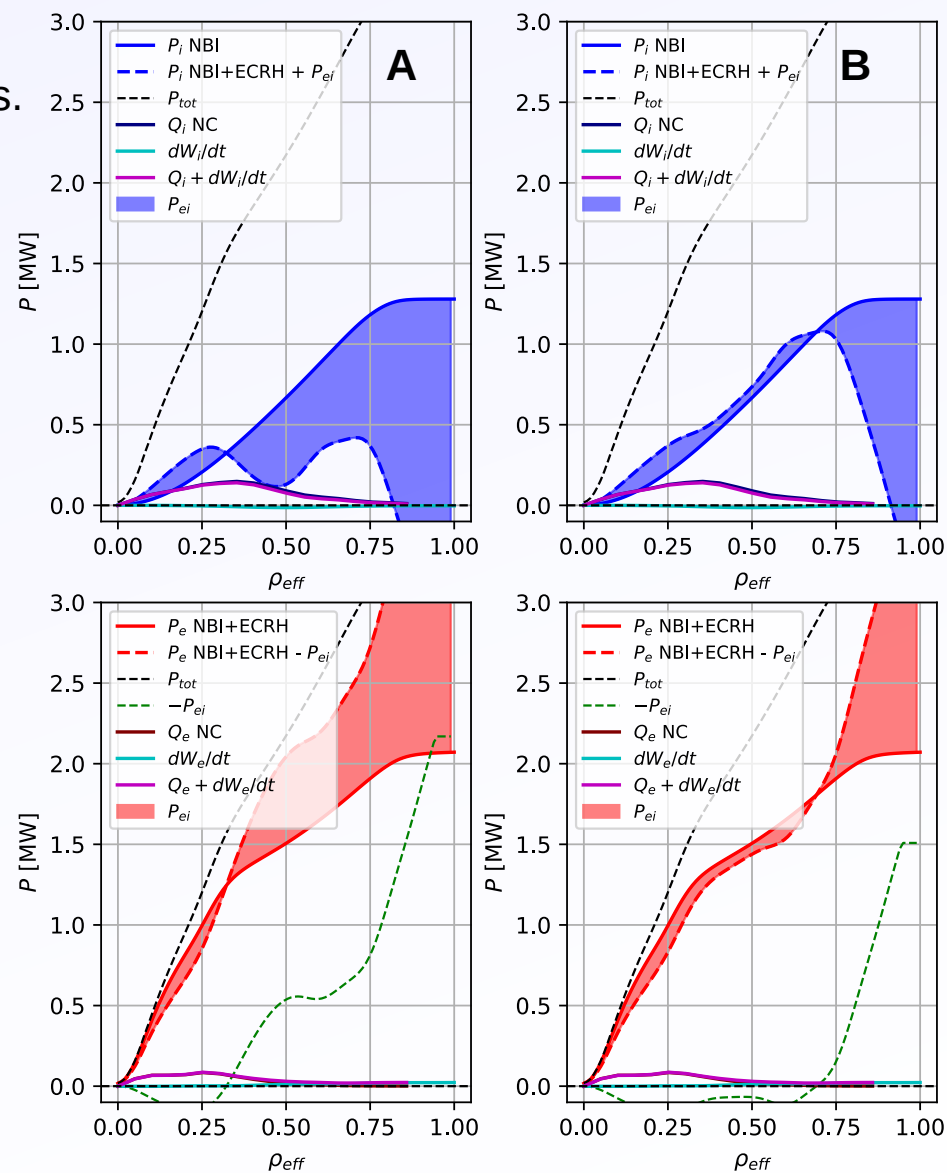
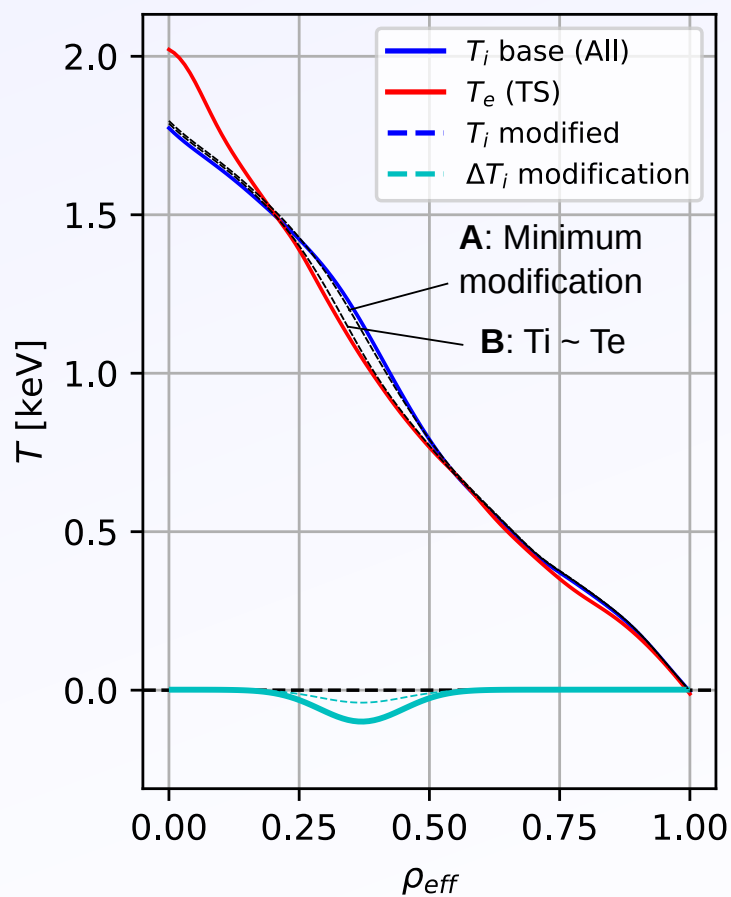
**A:** Minimum modification case *allows for a fully neoclassical transport barrier in  $Q_i$  at mid radius.*



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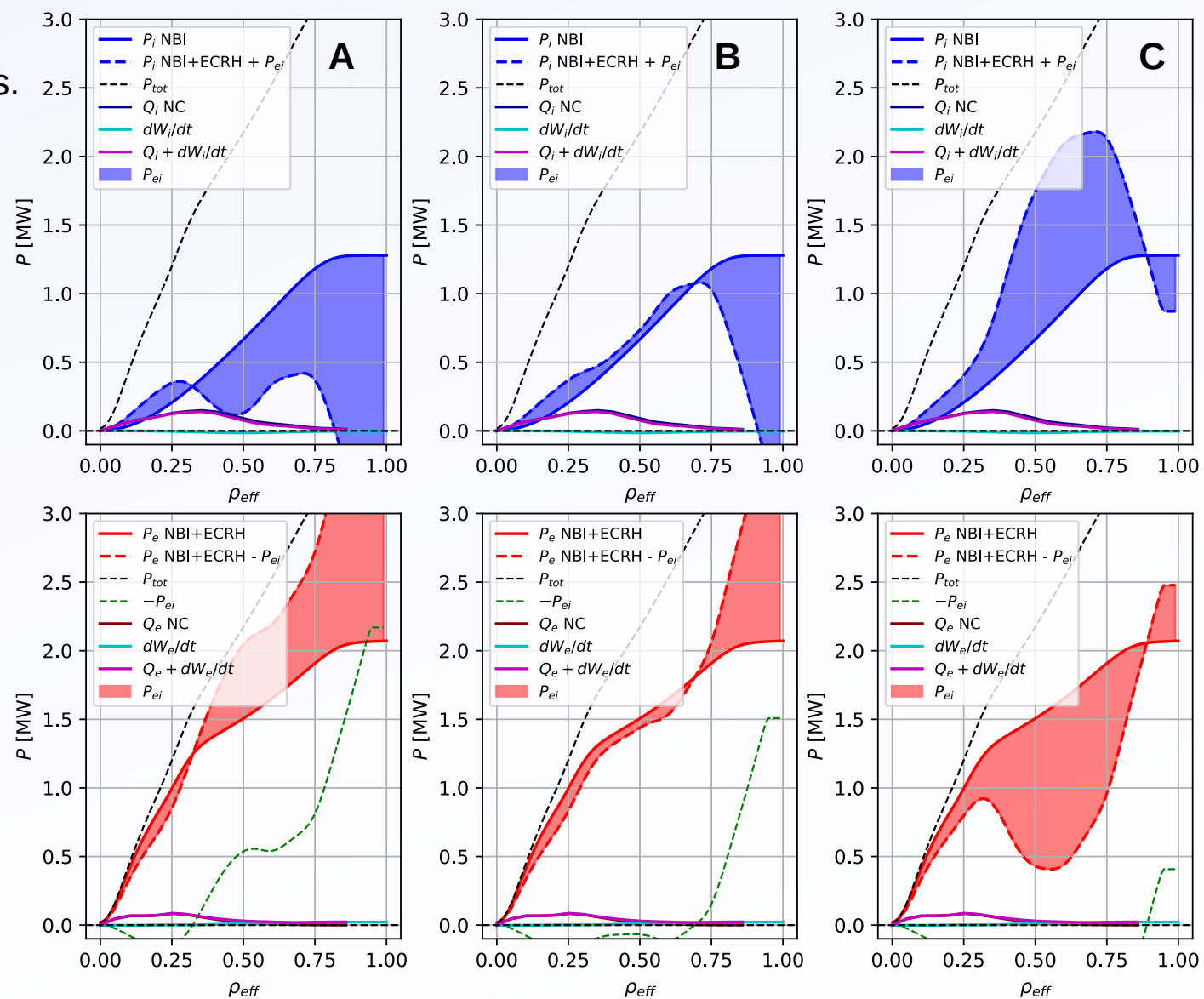
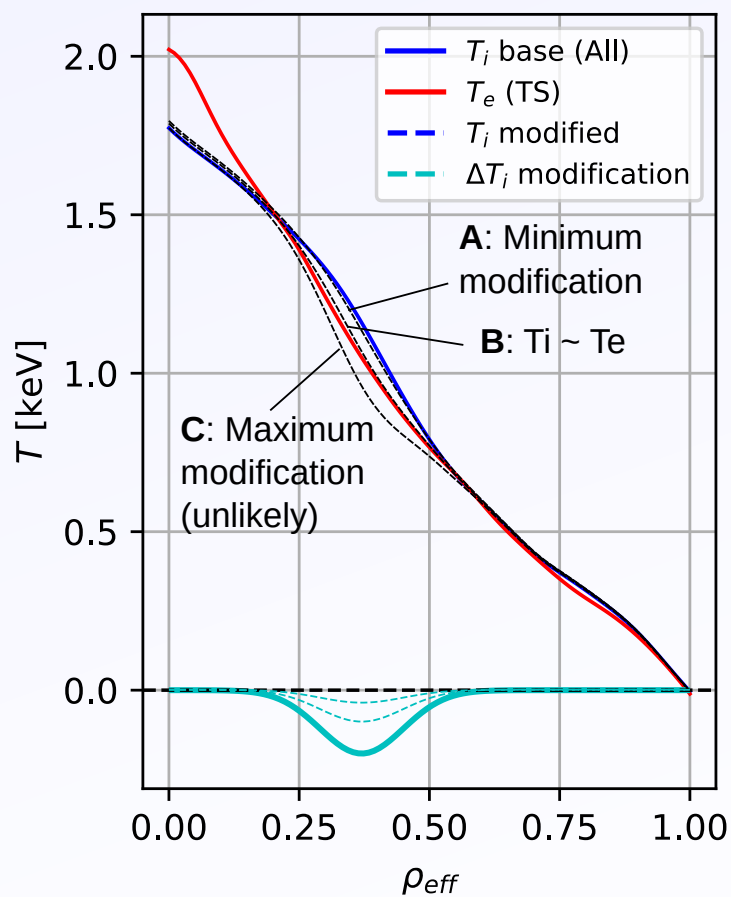
- A:** Minimum modification case **allows for** a fully neoclassical transport barrier in  $Q_i$  at mid radius.
- B:**  $T_i \sim T_e$  is also a possibility.



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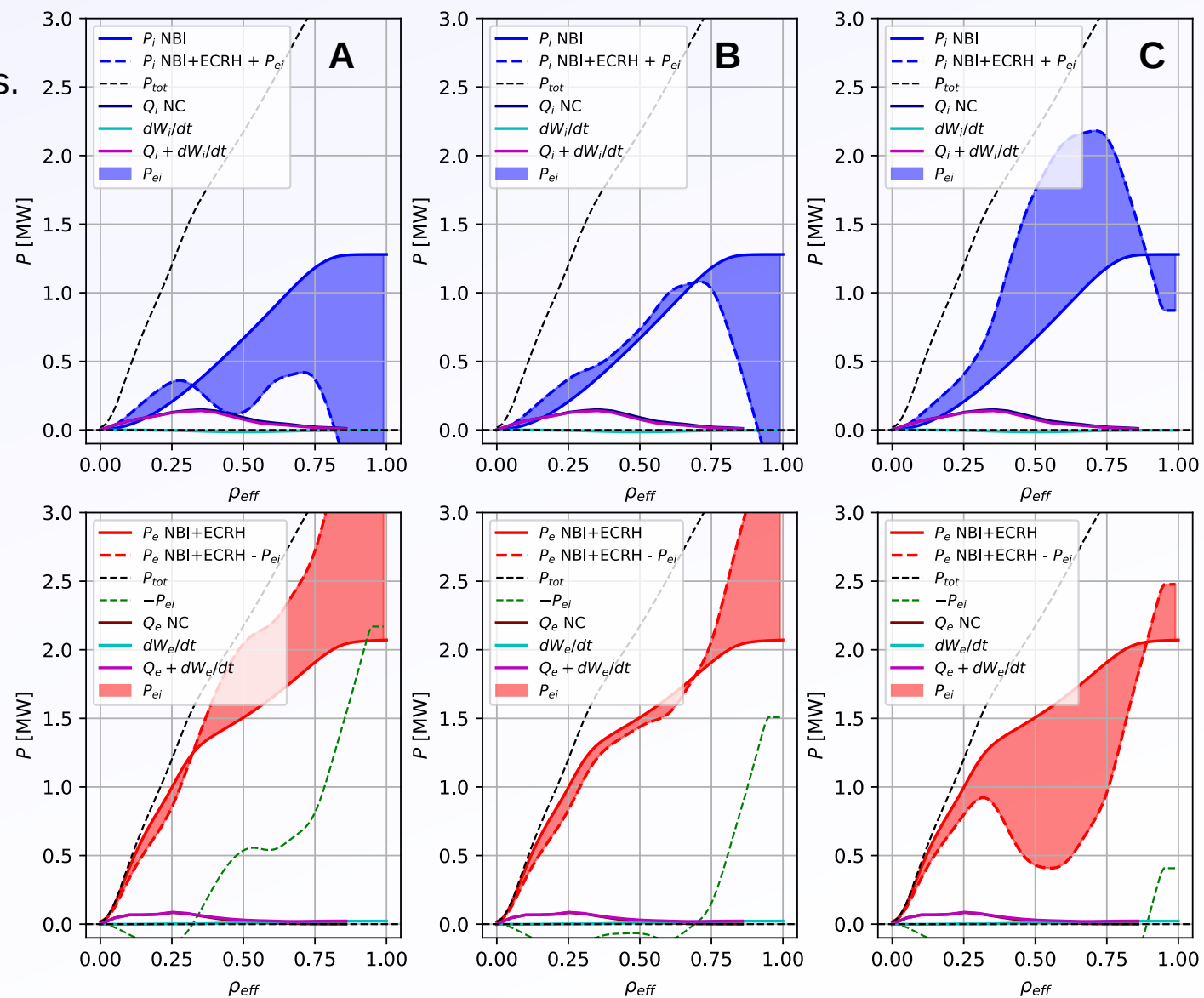
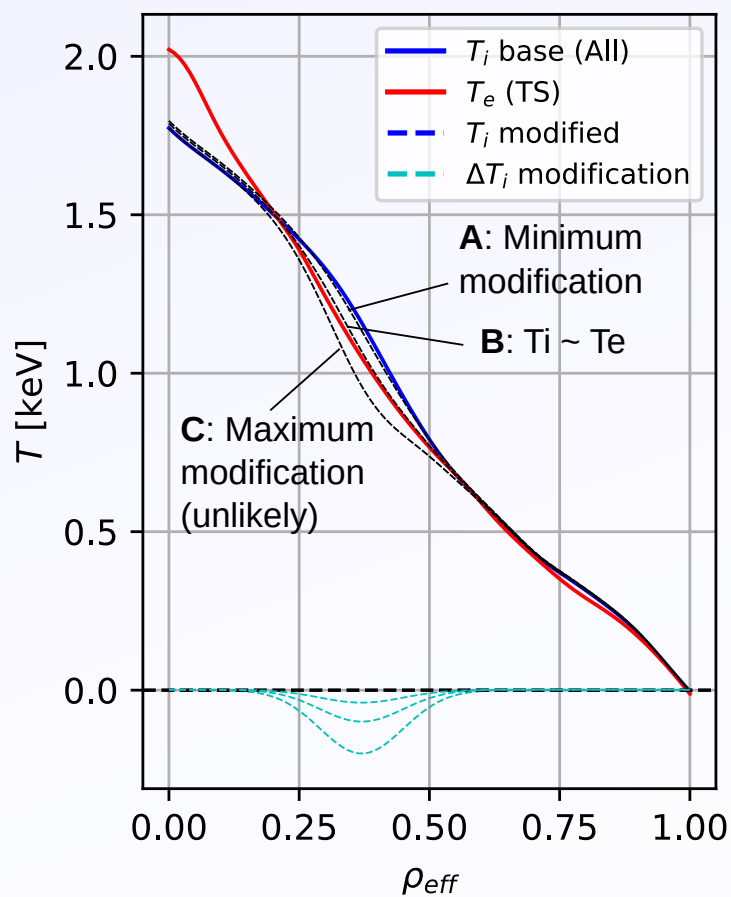
- A:** Minimum modification case *allows for* a fully neoclassical transport barrier in  $Q_i$  at mid radius.
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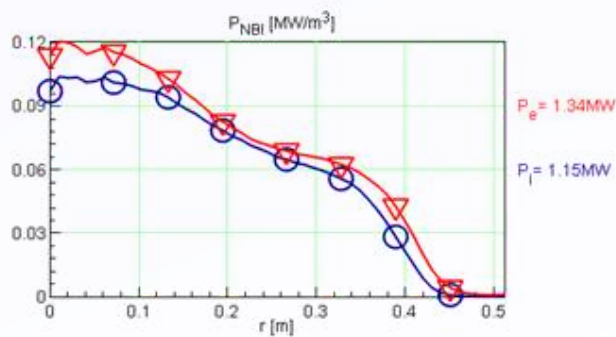
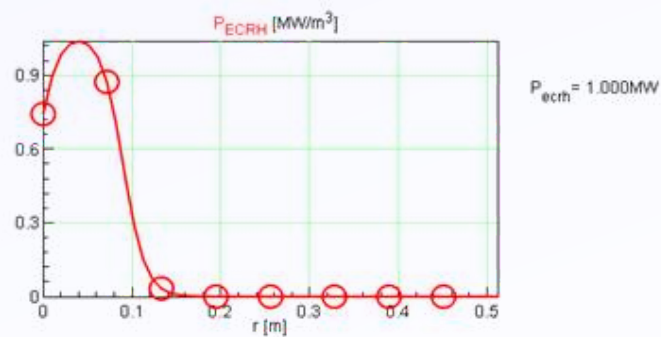
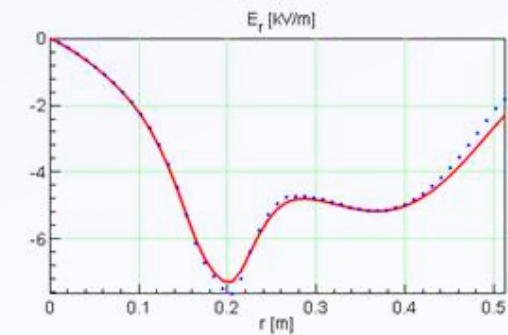
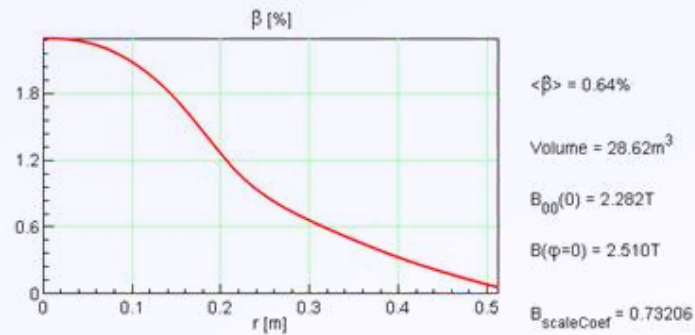
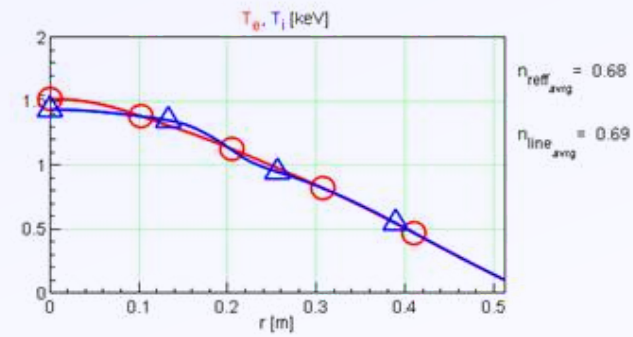
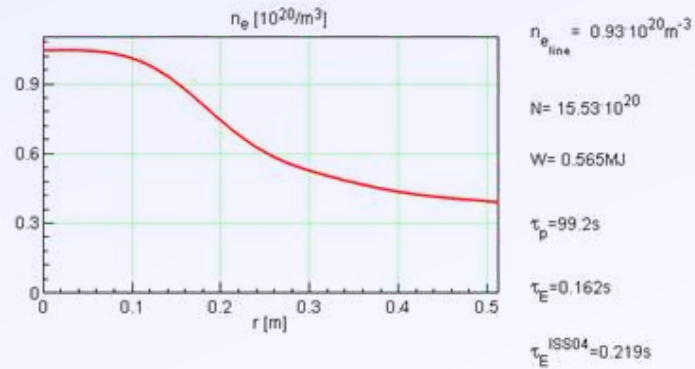




## Pure NBI - Species power balance

- We can construct this situation in NTSS:

[Beurskens]

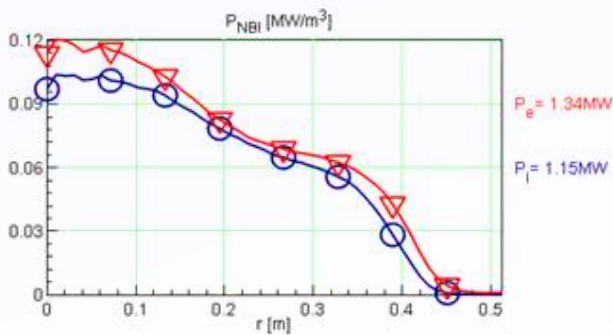
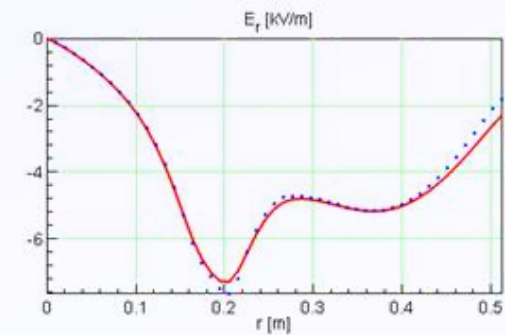
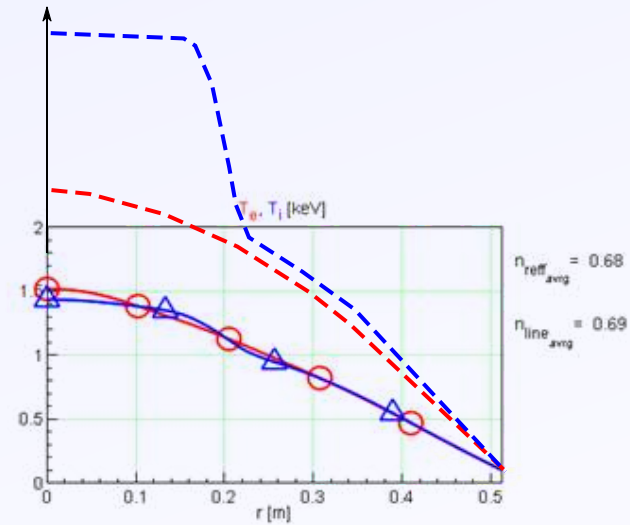
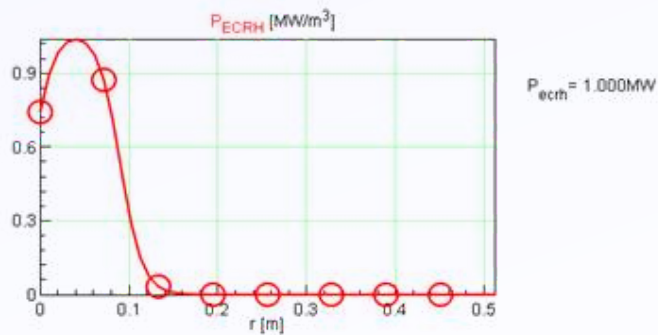
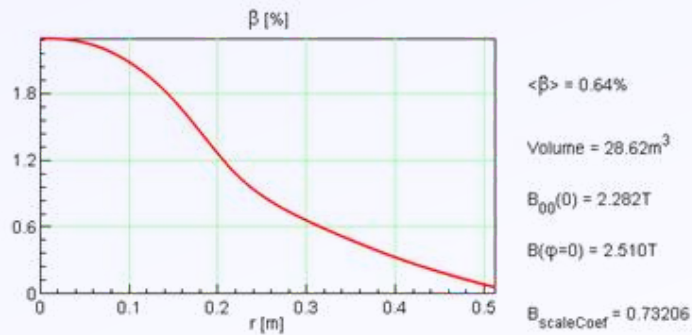
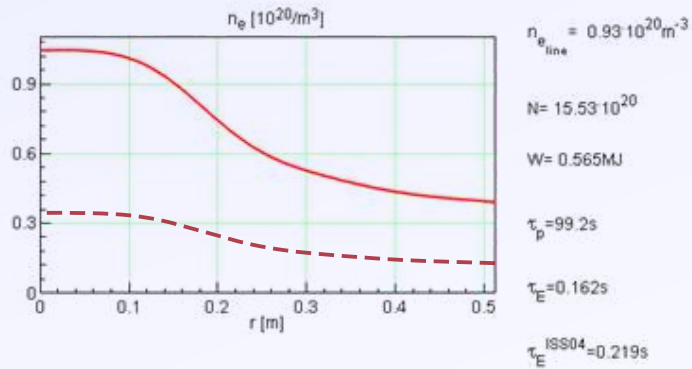




# Pure NBI - Species power balance

- We can construct this situation in NTSS:

[Beurskens]



At low density such a barrier would be very significant, but we would not get the  $a/L_{ne}$  required to create it.

## Pure NBI - Species power balance

- Intermediate conclusions from profile analysis:

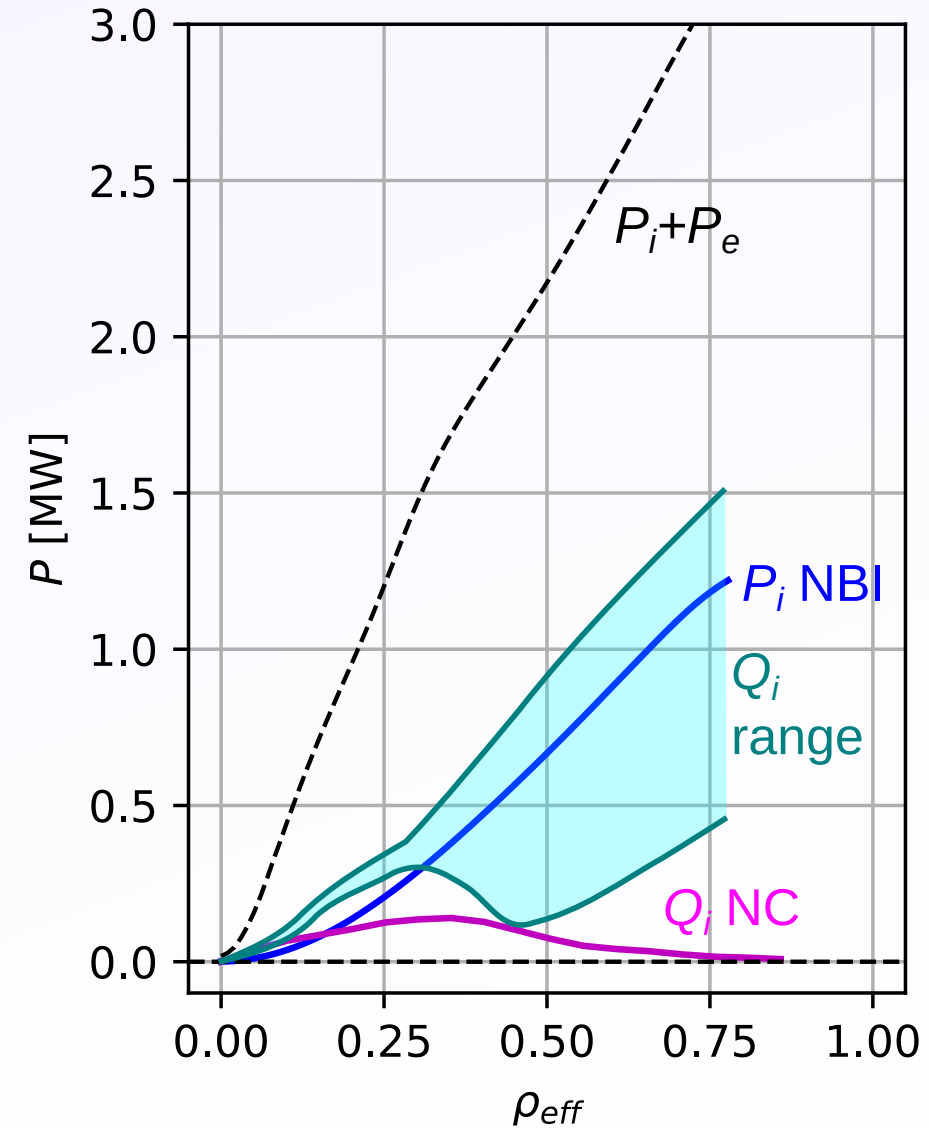
1) We can not separate  $Q_i$  and  $Q_e$  at high collisionality without improvements to the CXRS, XICS and TS analysis!  
It needs  $\sim 50\text{eV}$  accuracy, which is hard (but not impossible).

2)  $Q_i$  at  $\rho = 0.5$  is somewhere between NC value and  $Q_i^{\text{NBI}}$ .  
It is unlikely to have taken a large fraction of the ECRH power.

3) Fully suppressed ion turbulence barrier is very possible at  $\rho = 0.5$ ,  
conincident with the apparent particle transport barrier.

However, this is not useful, since all power is transferred to  
electrons, so that  $T_i = T_e$ .

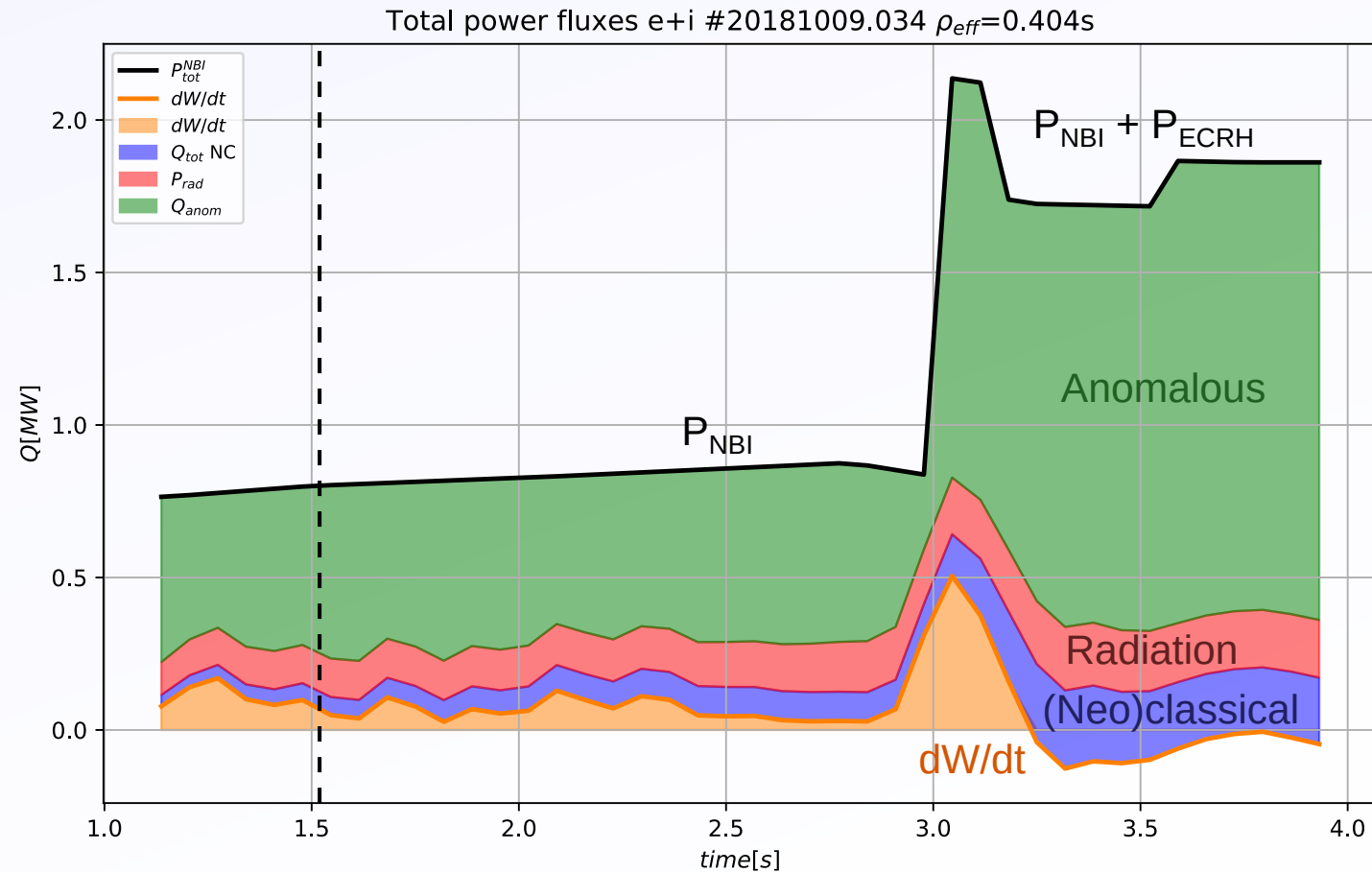
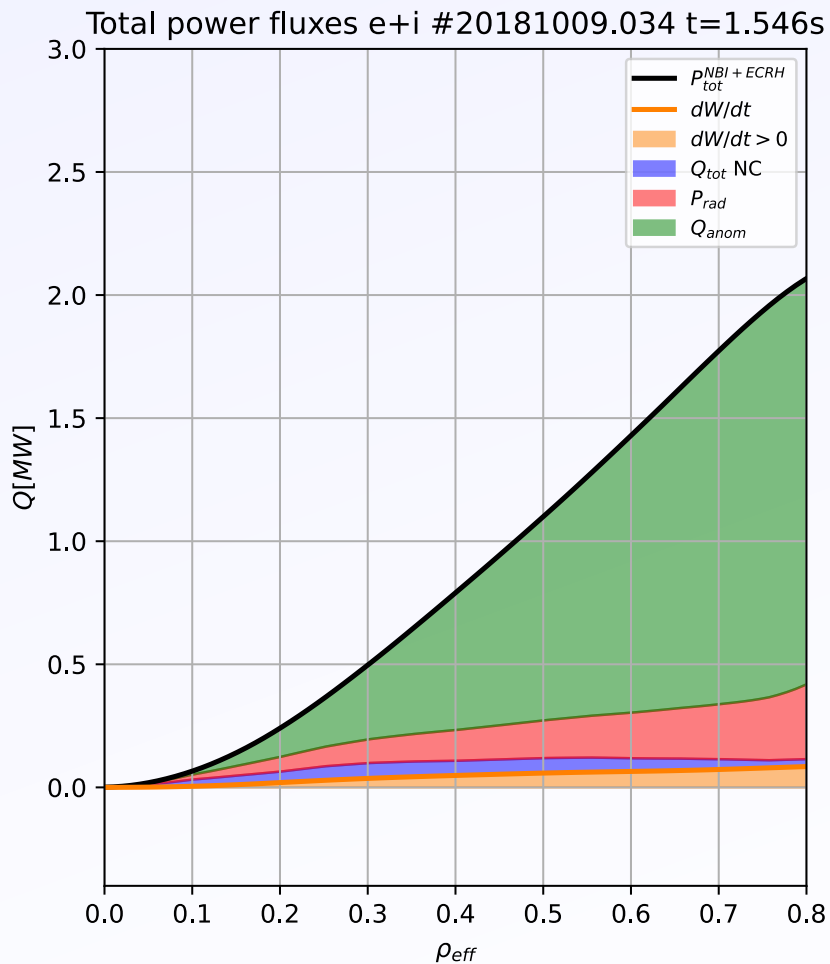
--> In high collisionality plasmas, the species with fastest heat transport  
completely determines both temperatures and stored energy.



## Total power balance

We can ignore  $P_{ei}$  and examine the balance of  $Q_{total}$ :

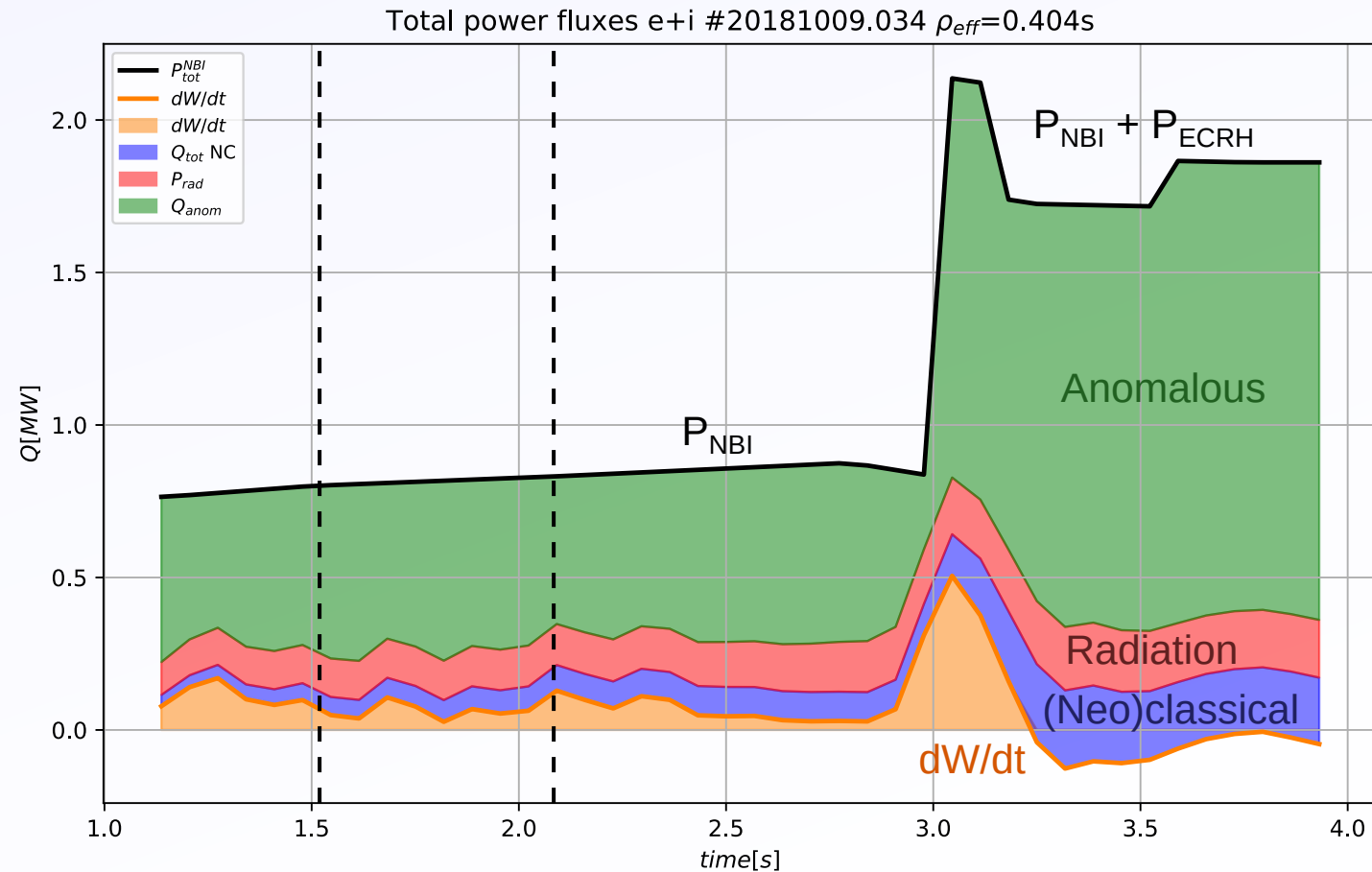
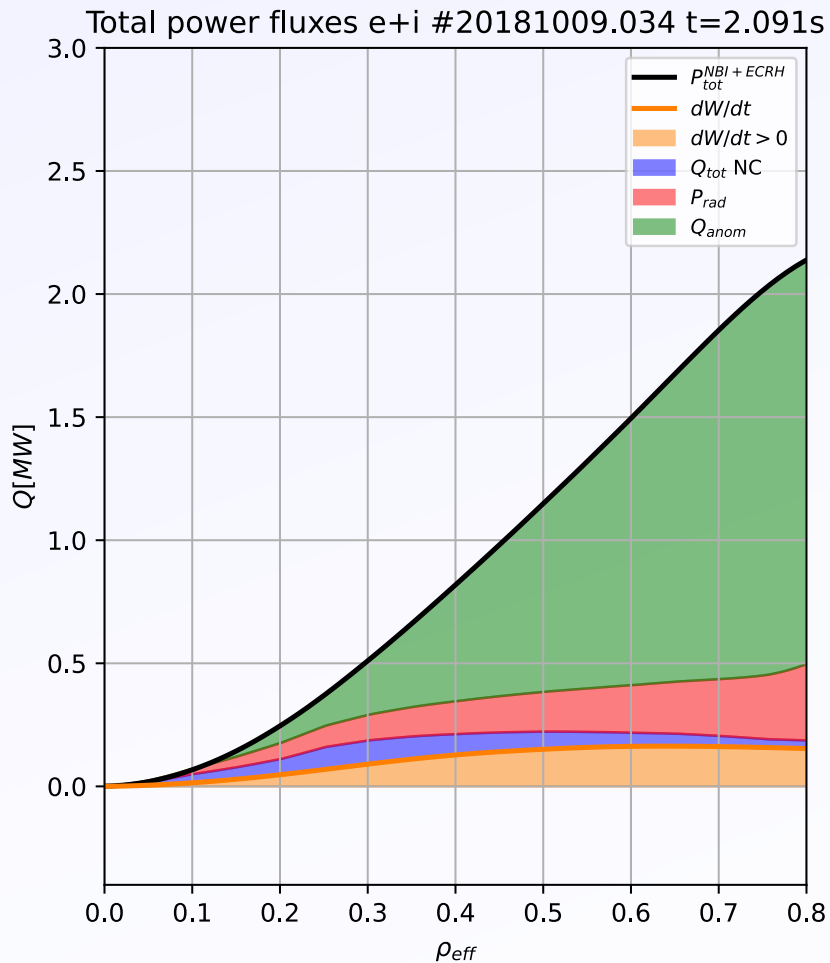
- Dominated by anomolous transport at all time points for  $\rho \geq 0.3$ .
- Both radiation and (neo)classical transport are small but significant (classical  $\sim 20\%$  x neoclassical).
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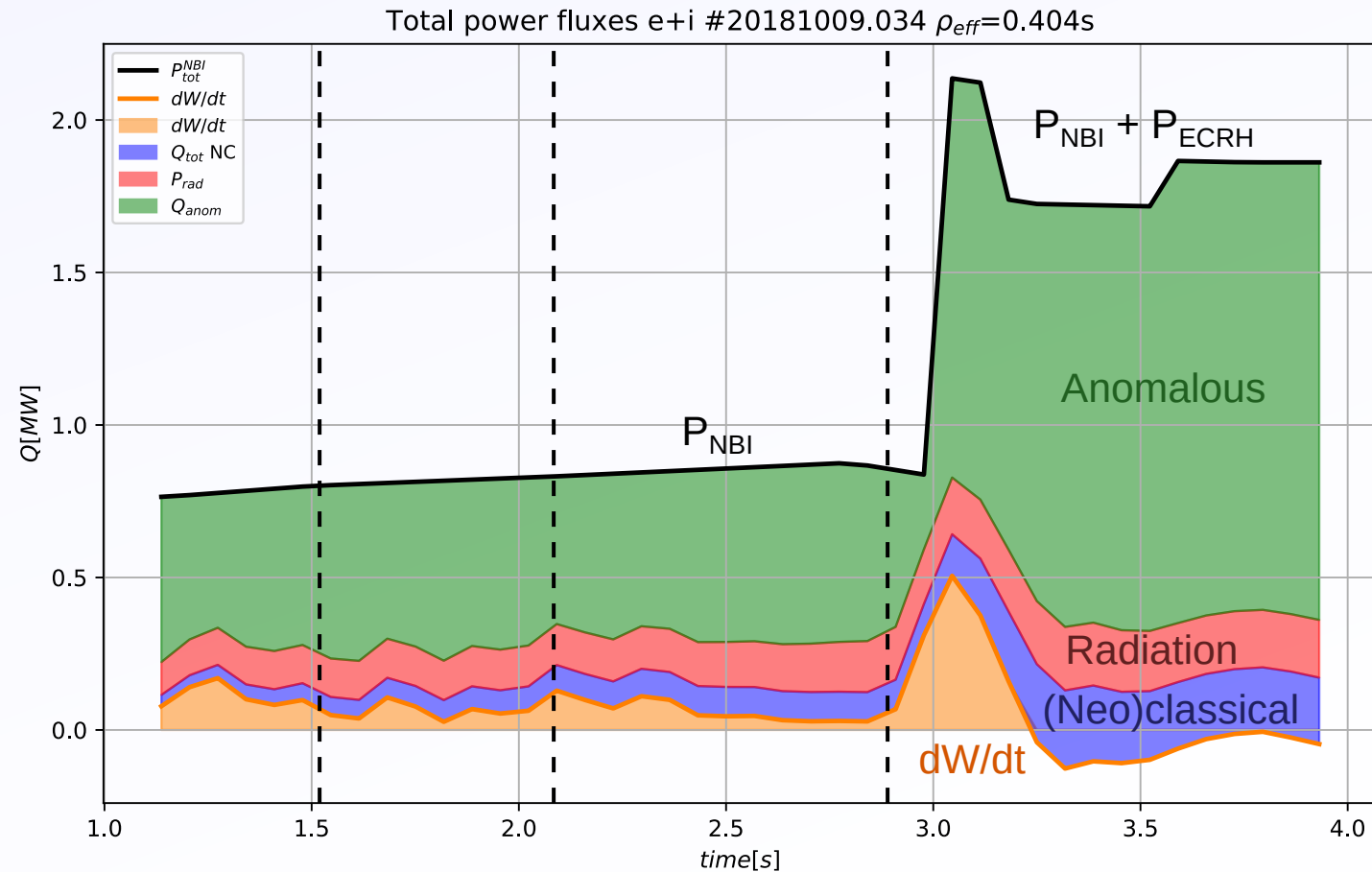
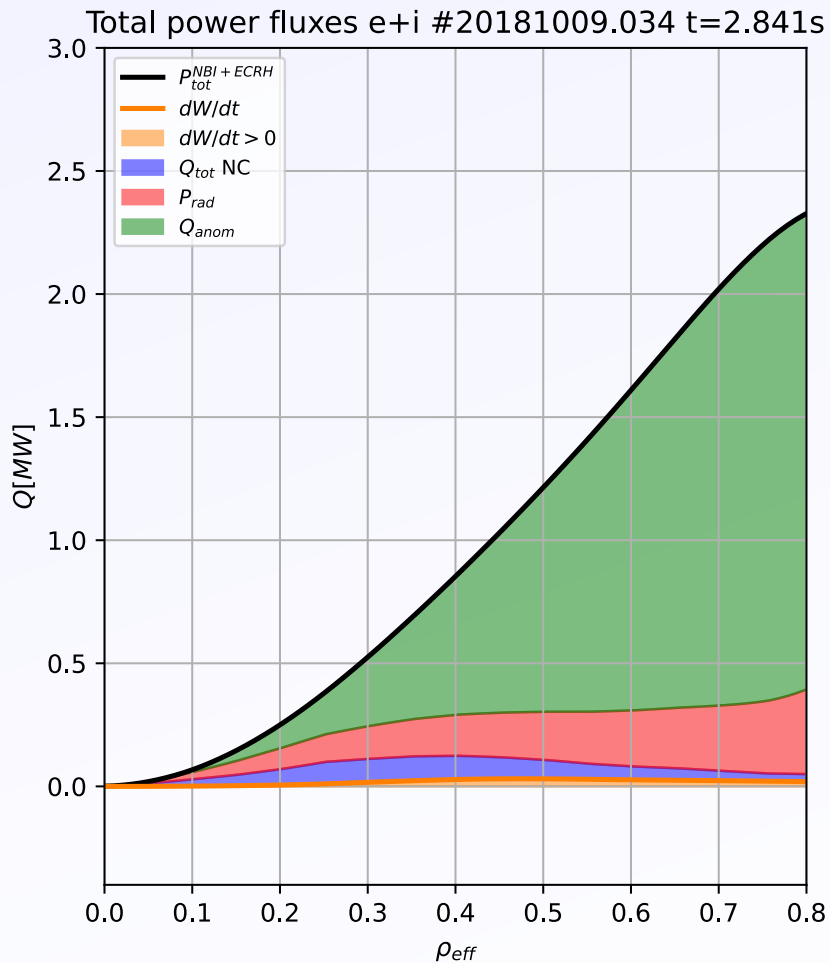
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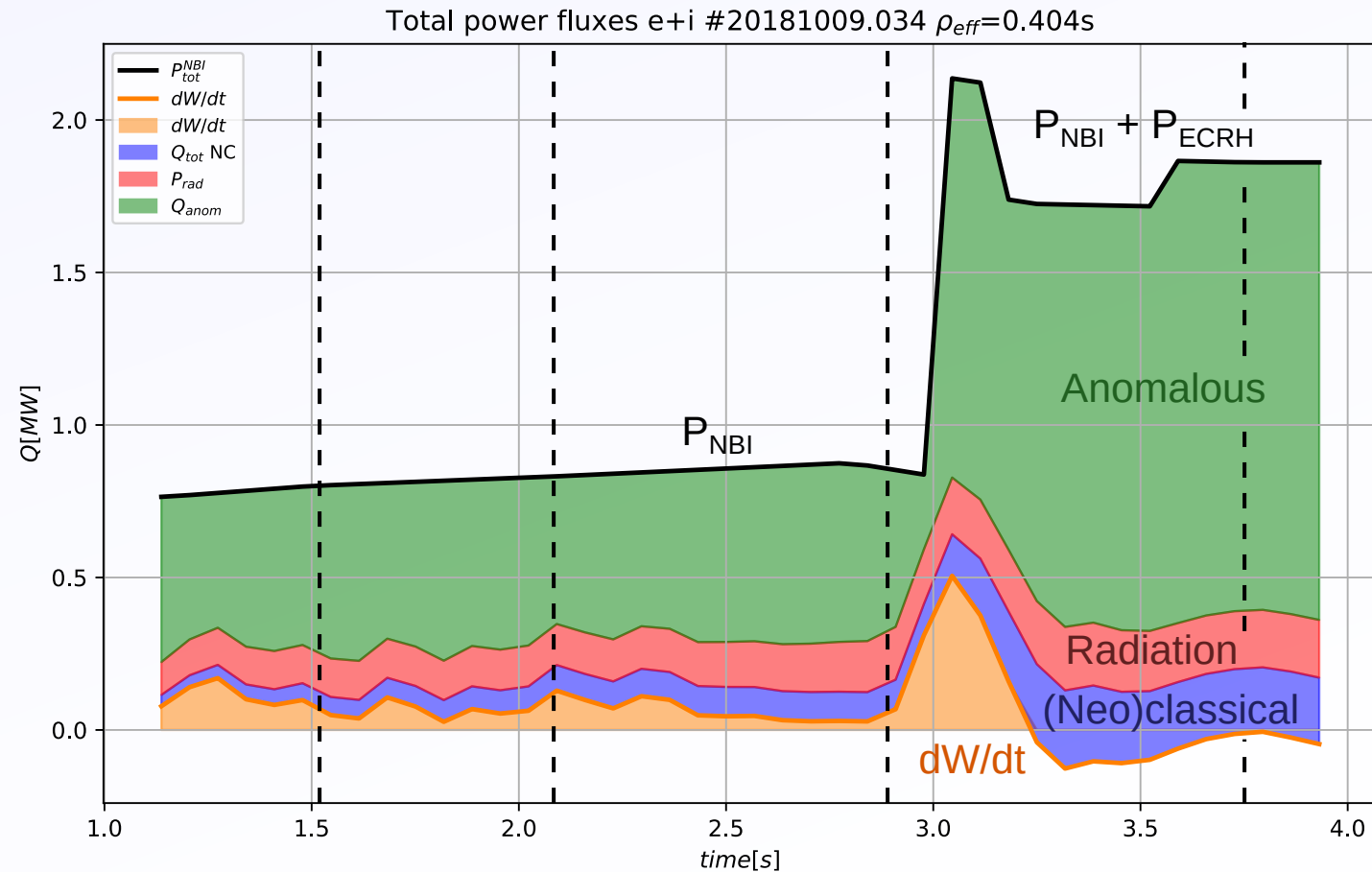
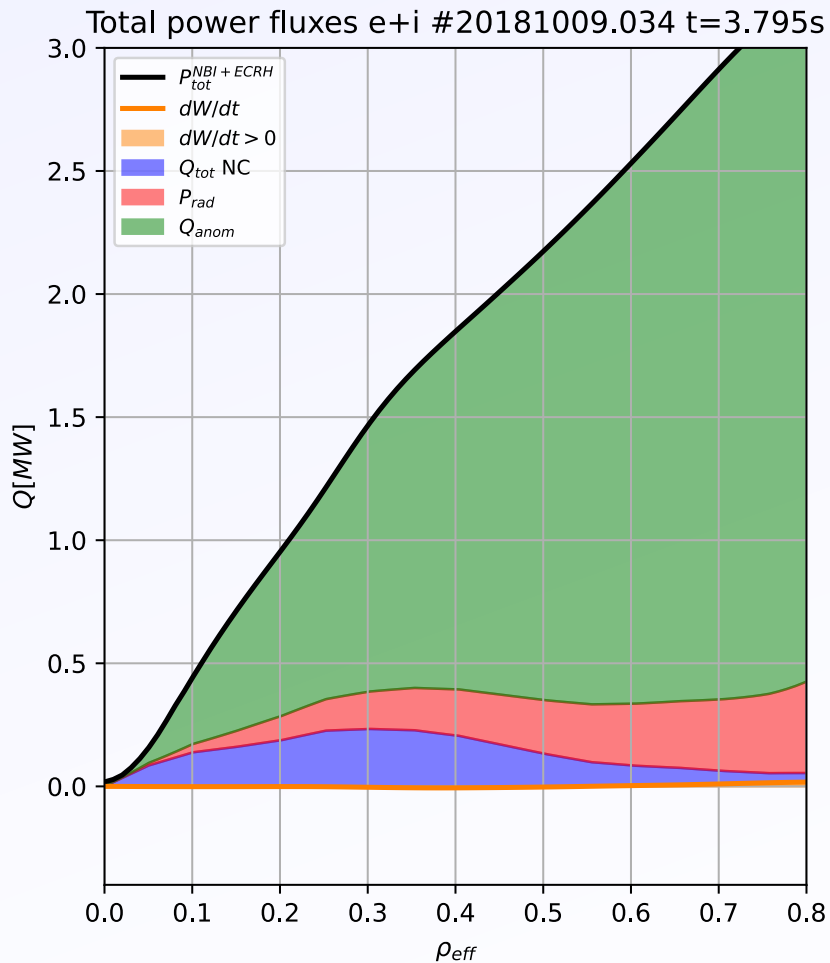




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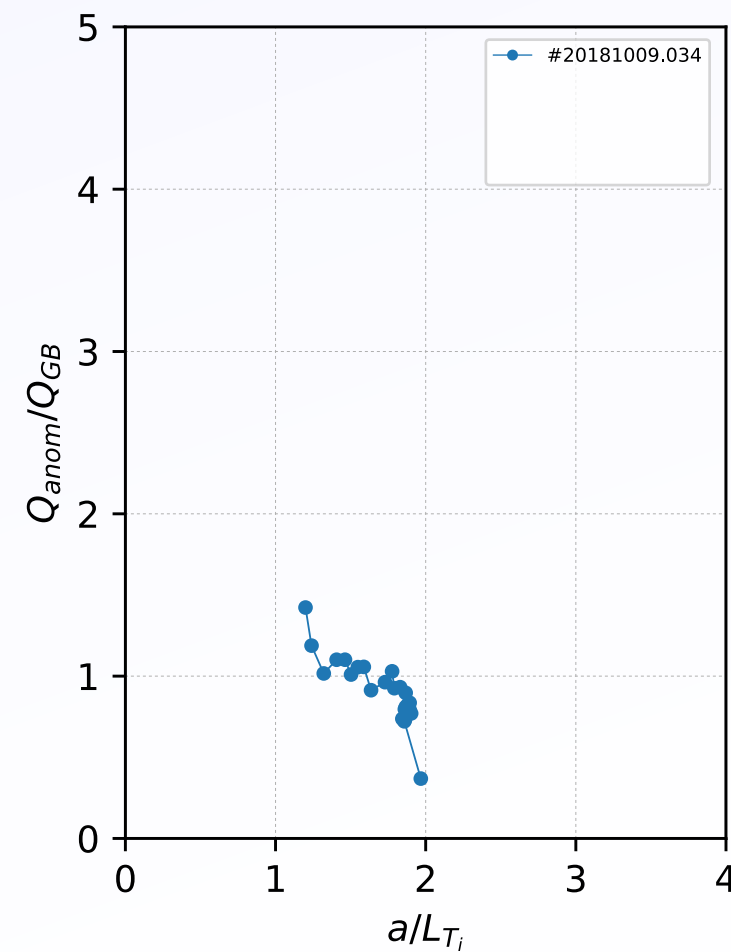
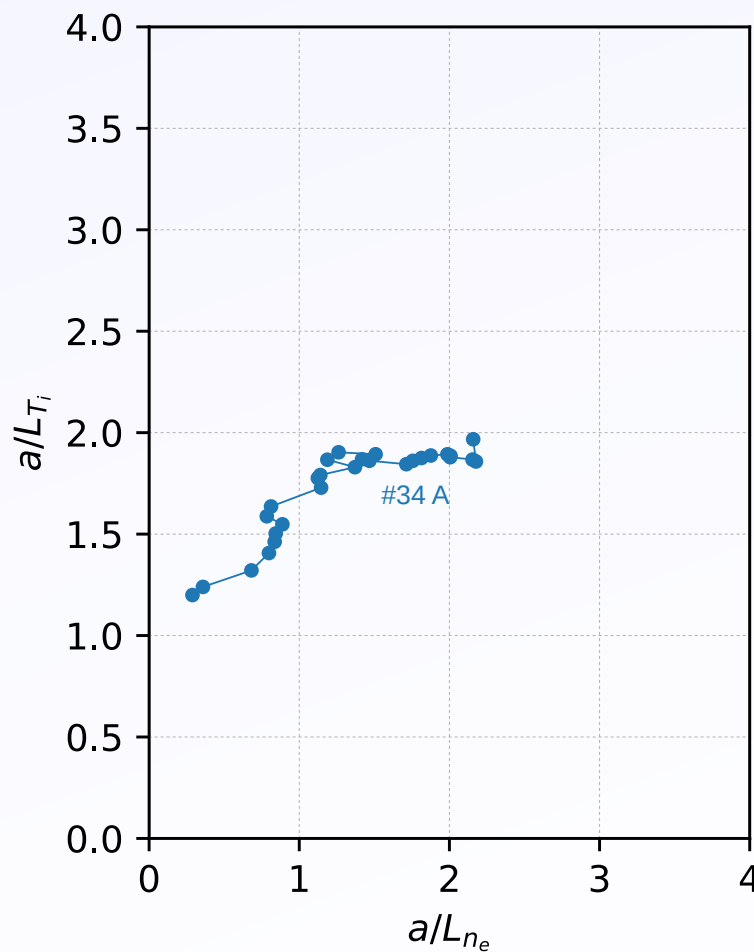
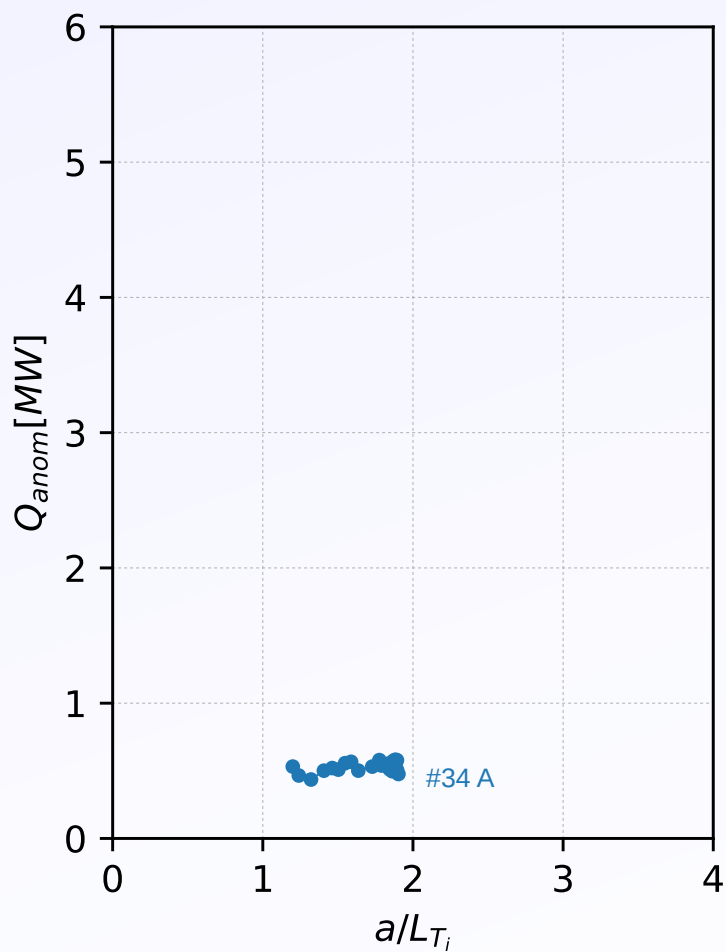




# Total power flux

Over multiple shots, a pattern emerges:

#34: **A)** Pure NBI phase builds up density gradient.

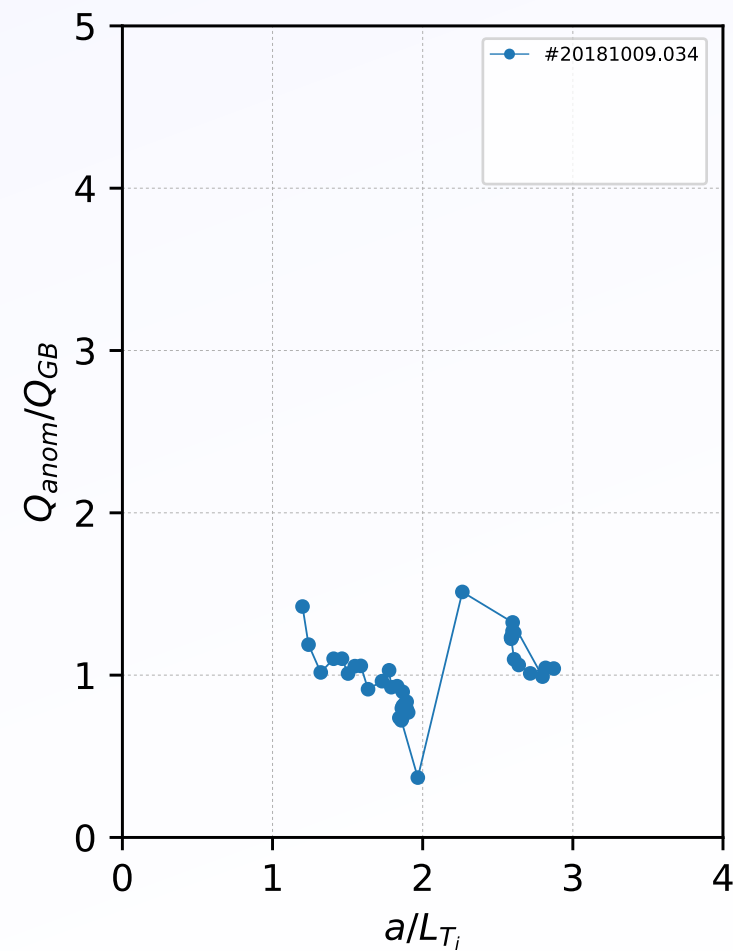
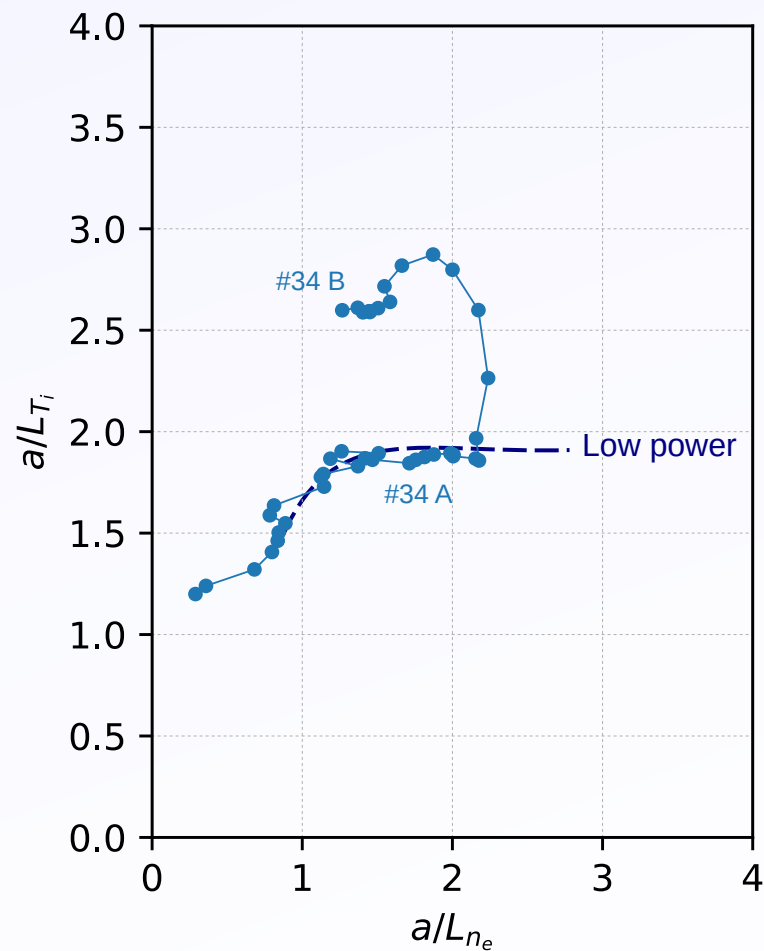
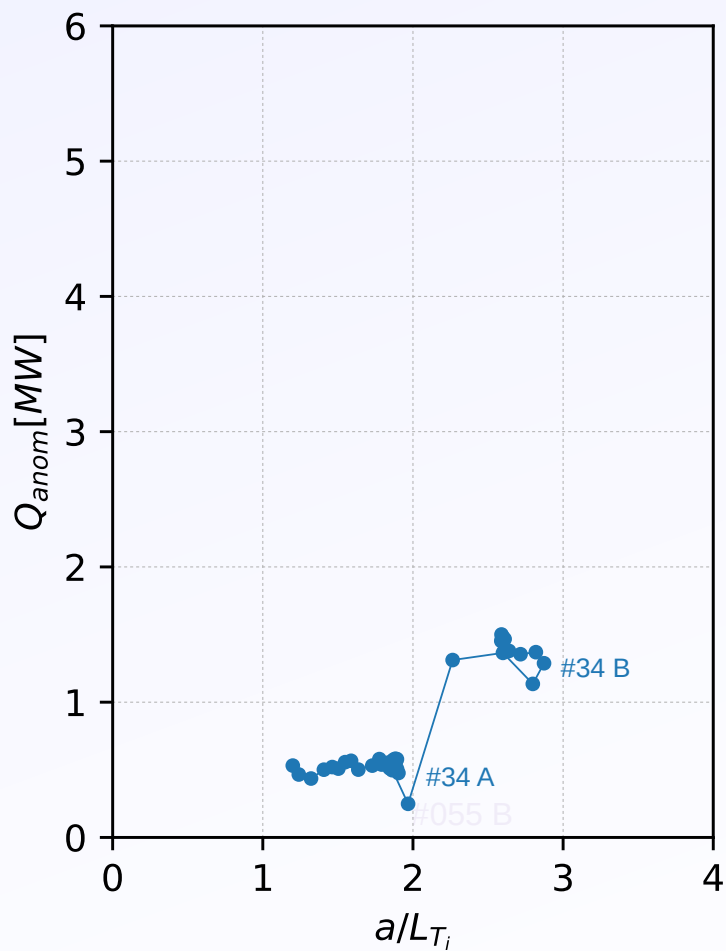




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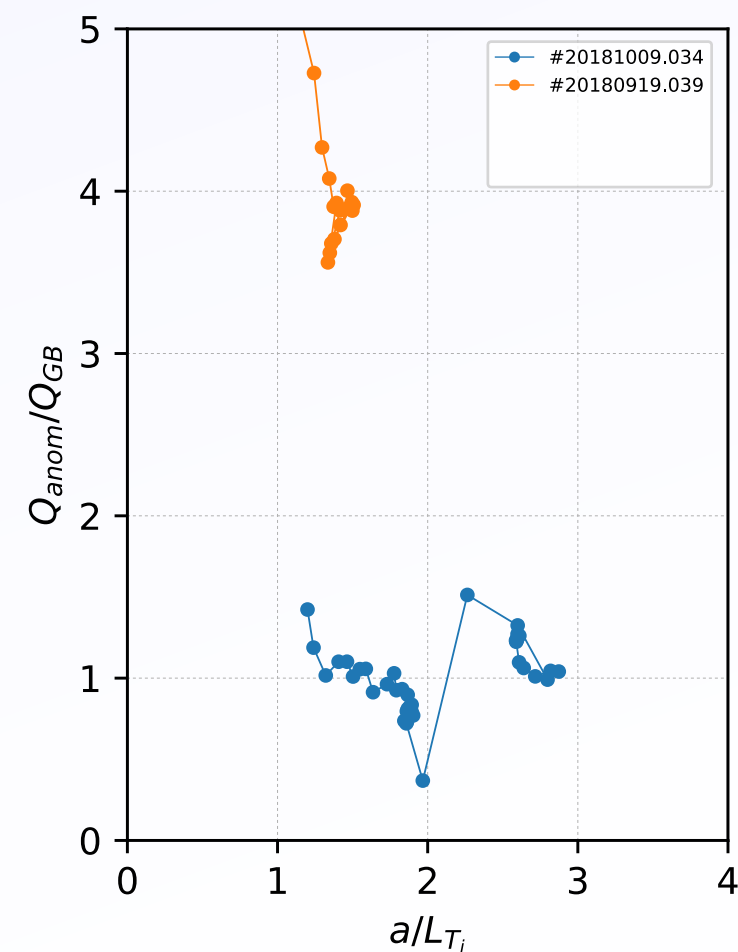
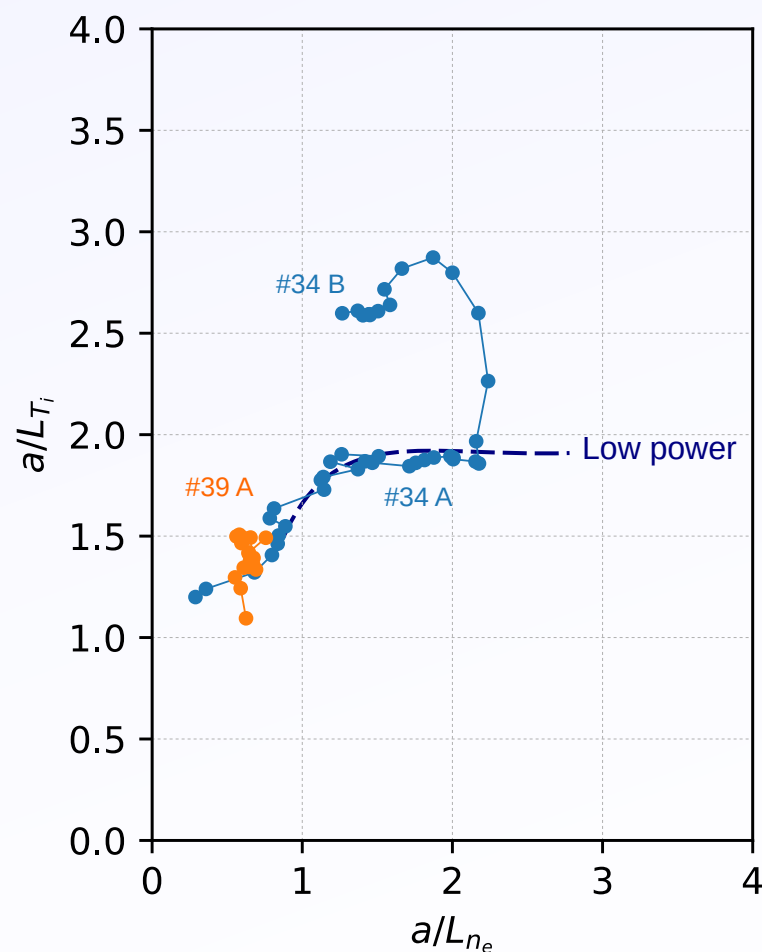
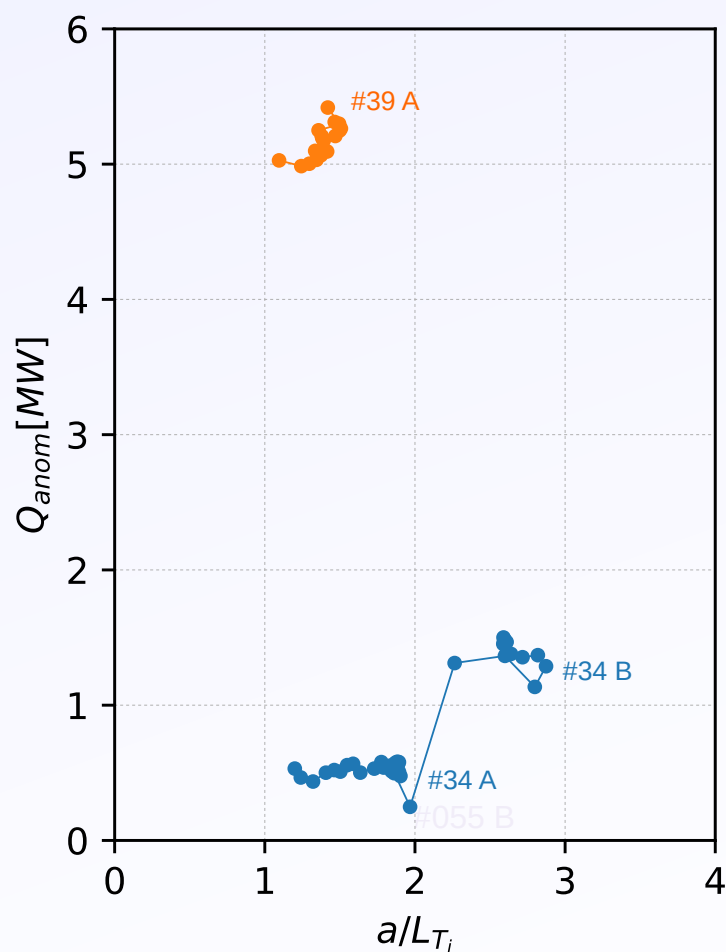


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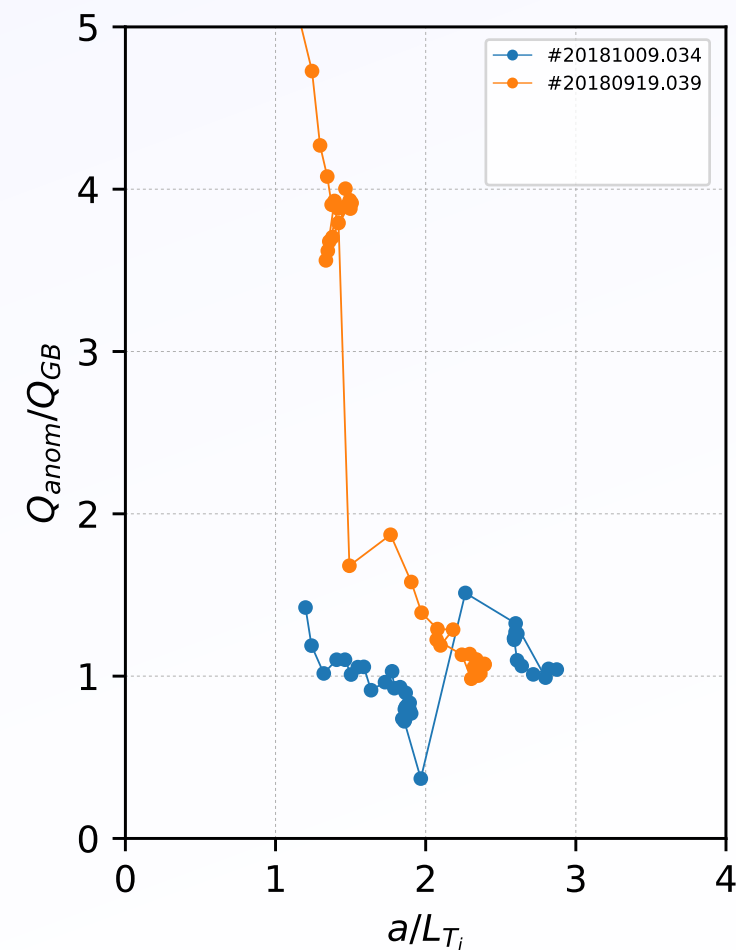
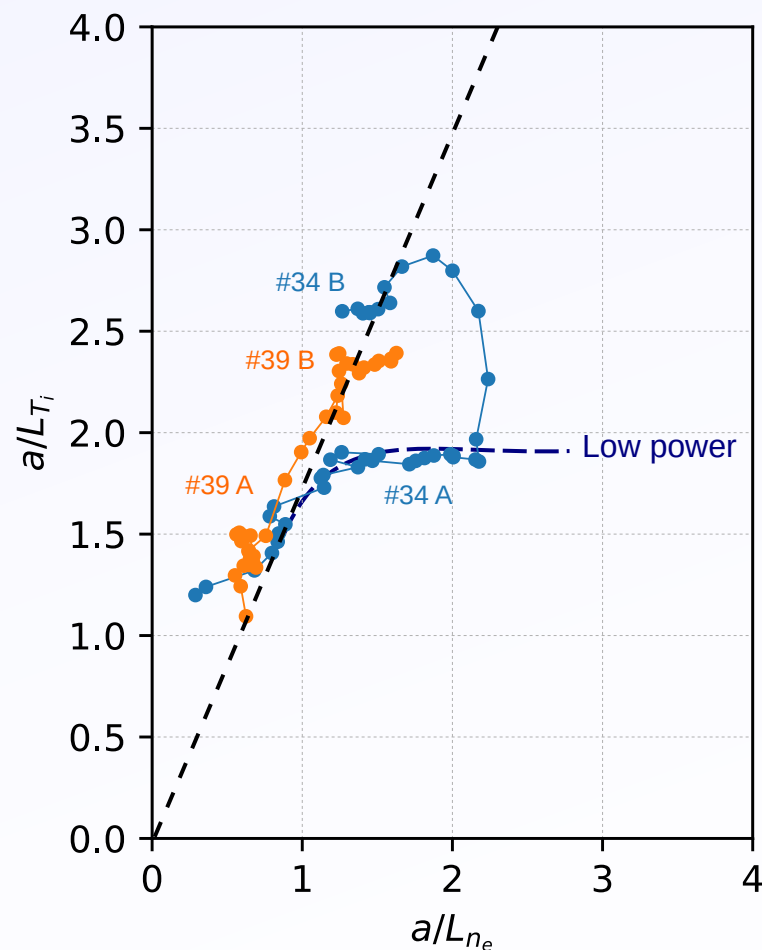
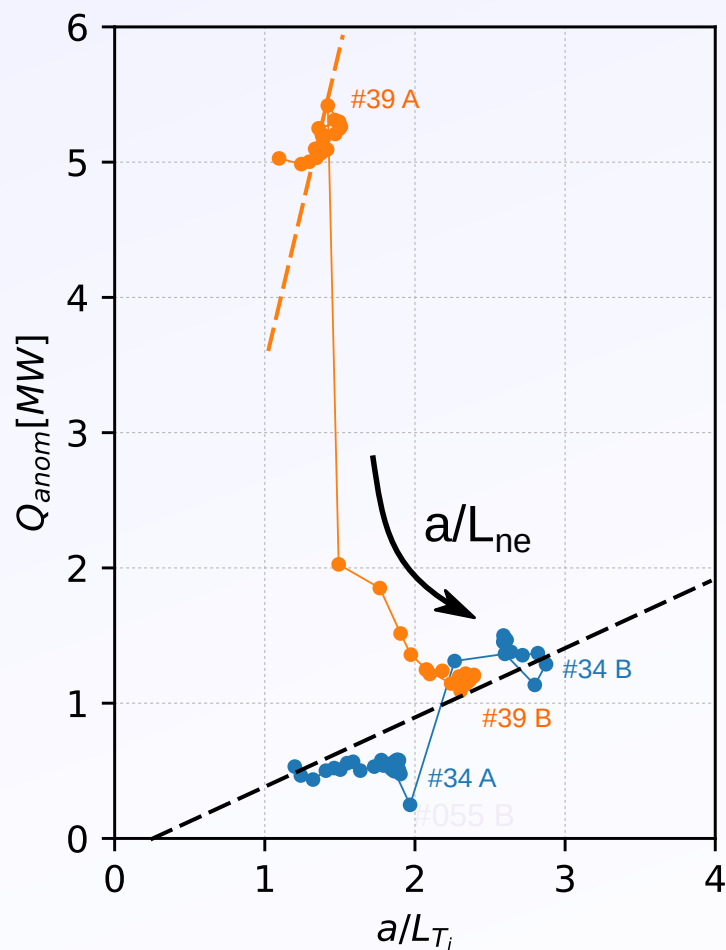


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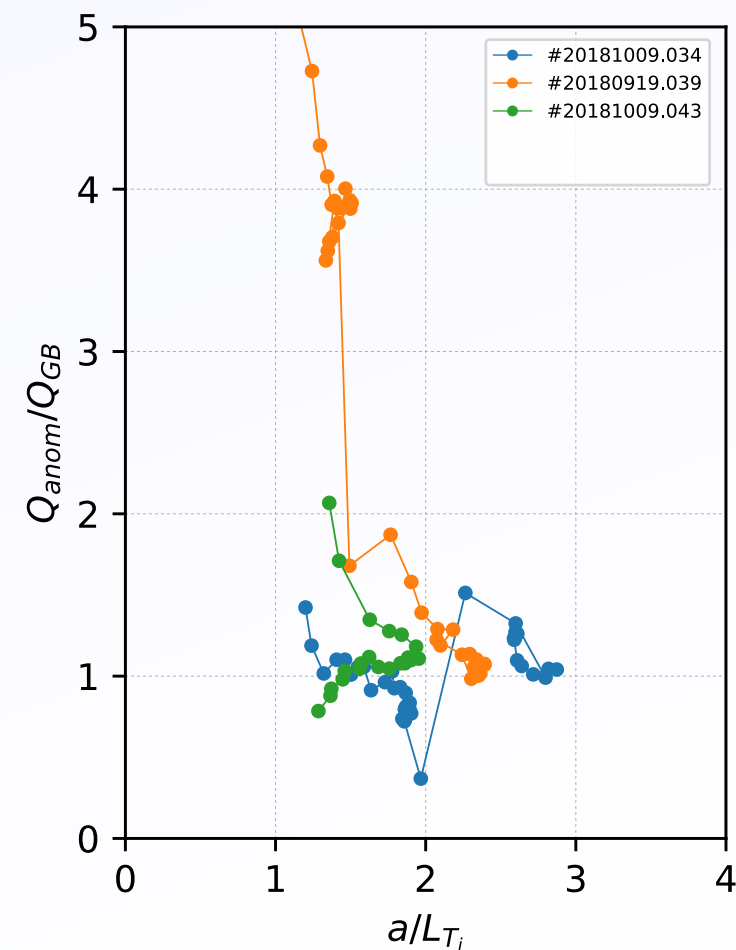
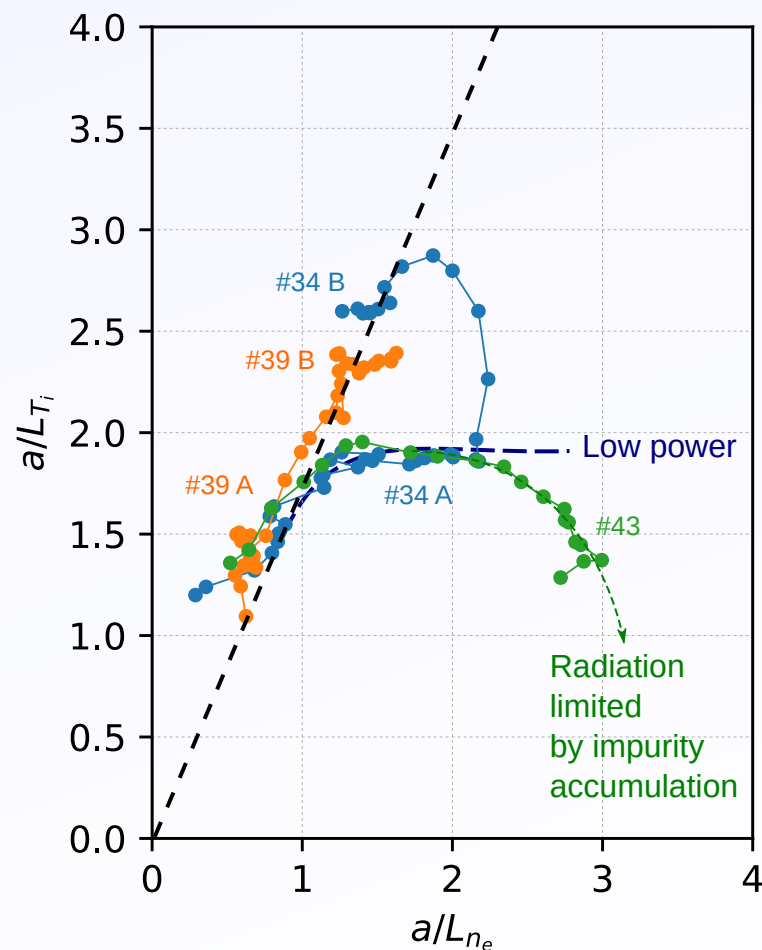
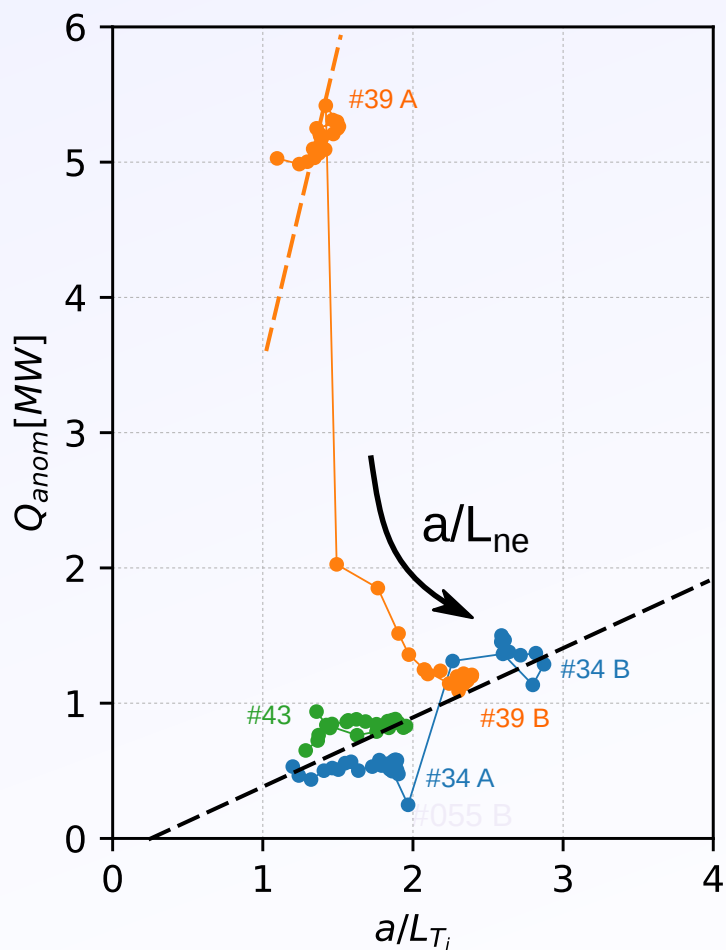




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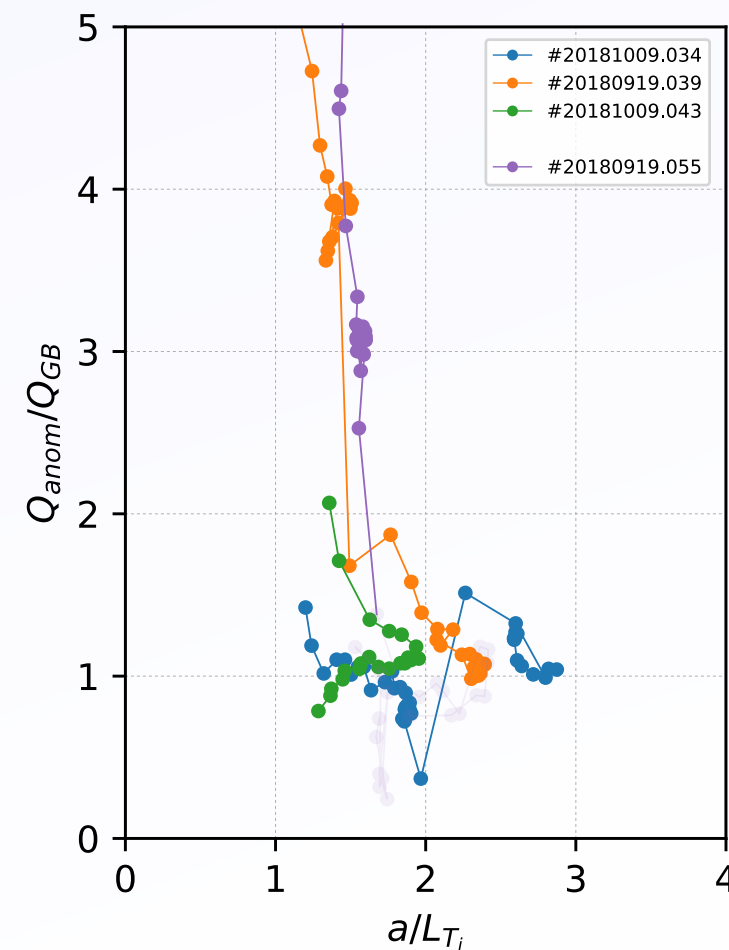
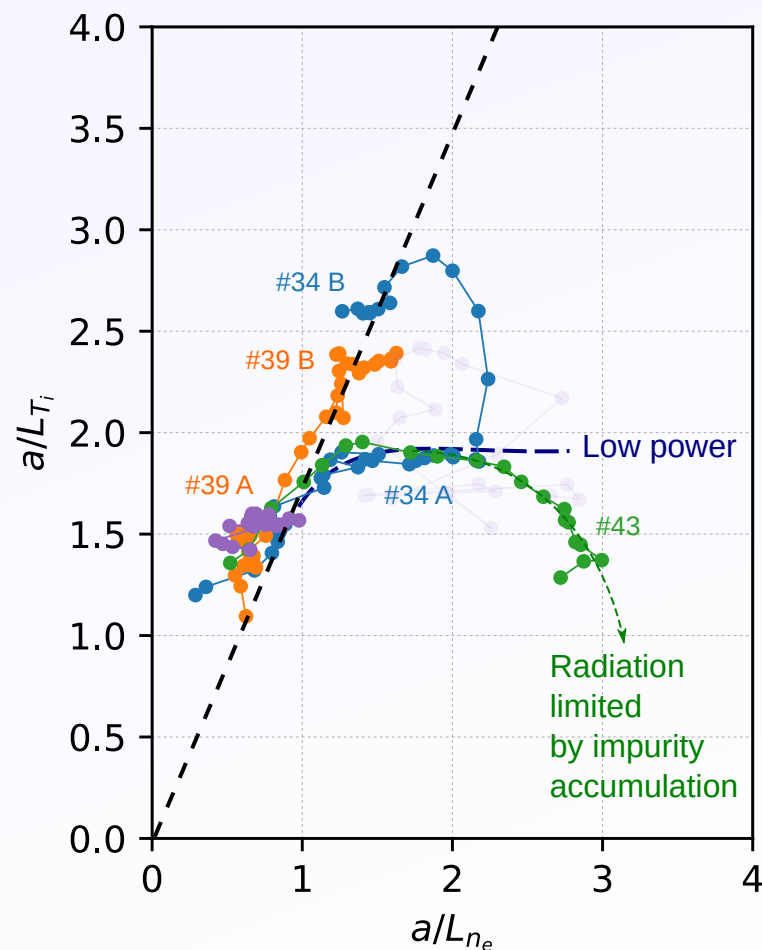
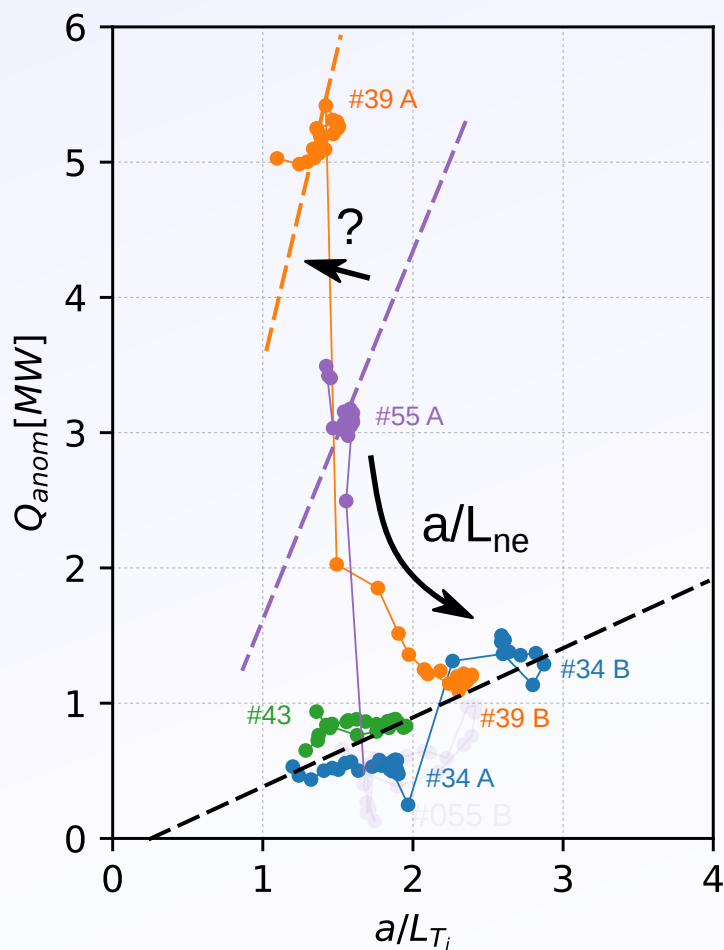
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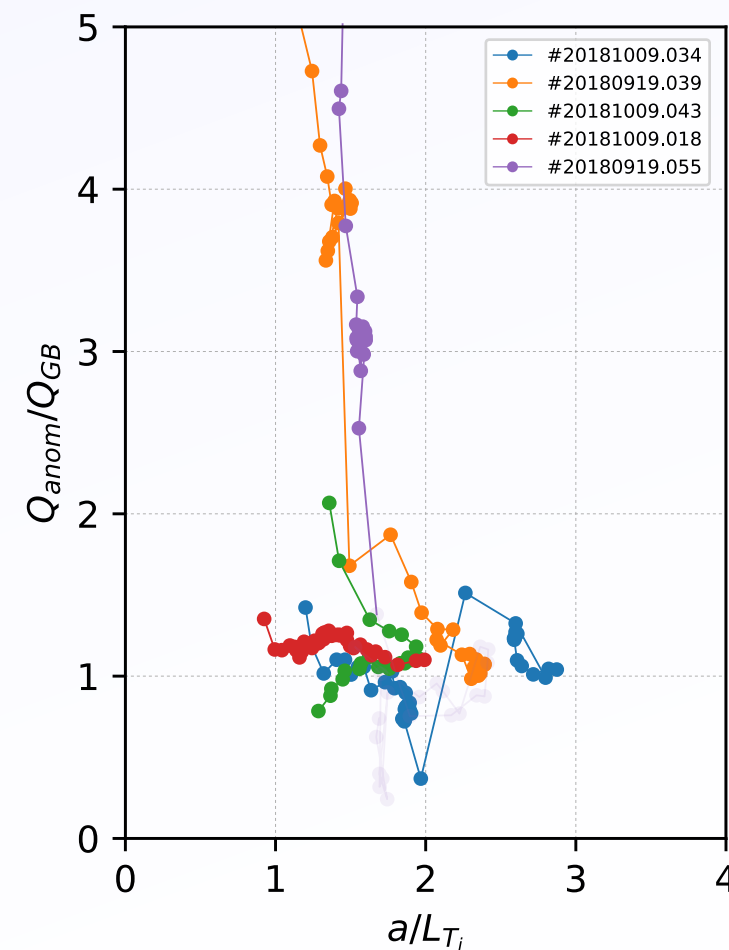
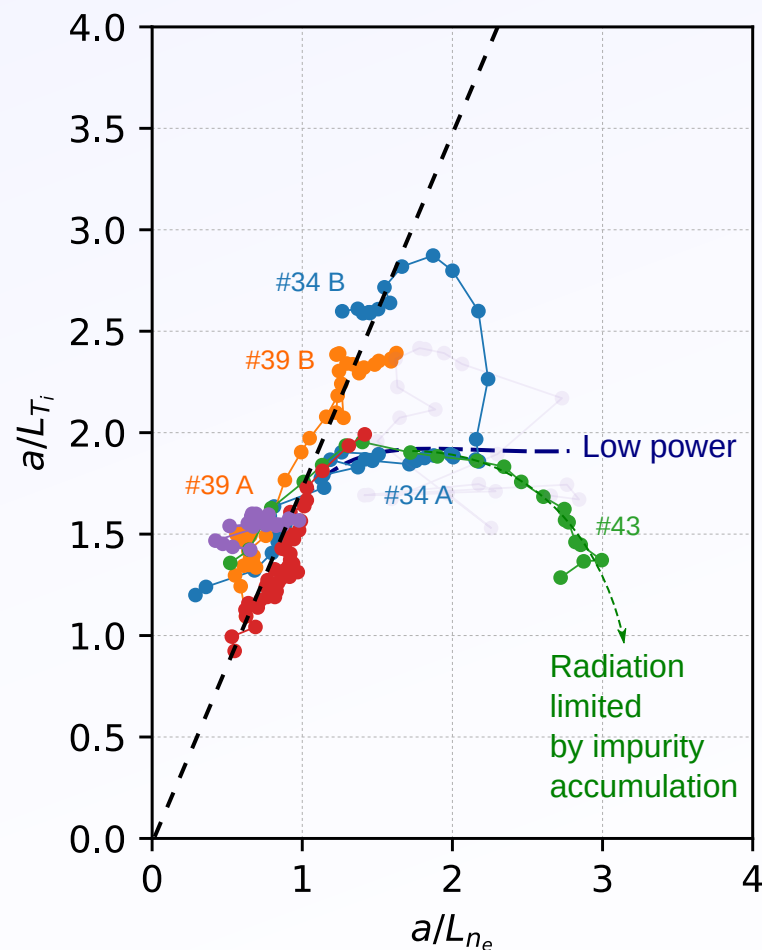
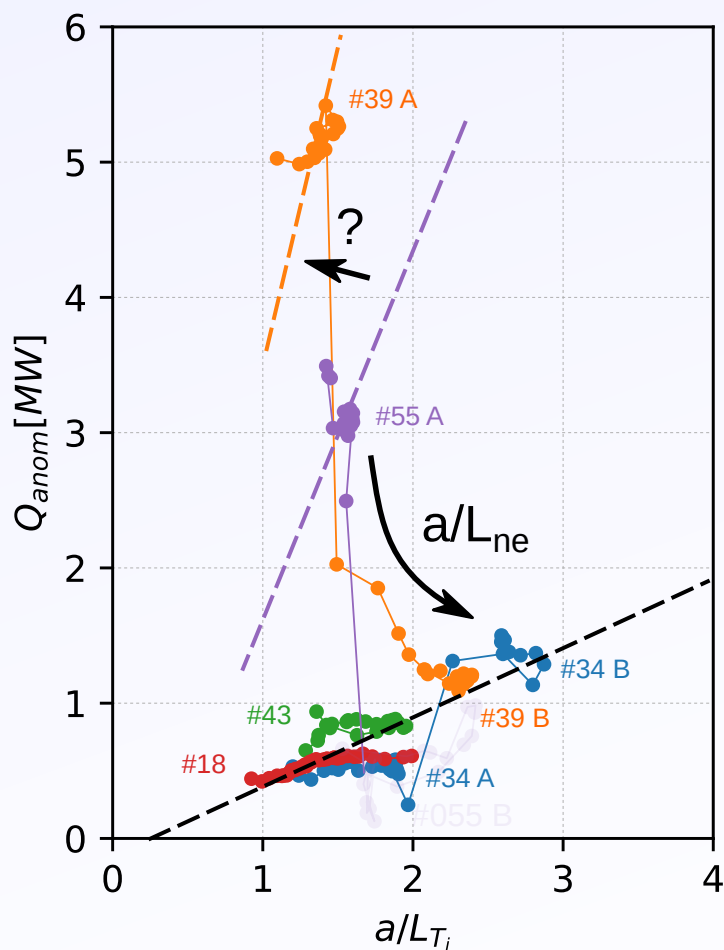
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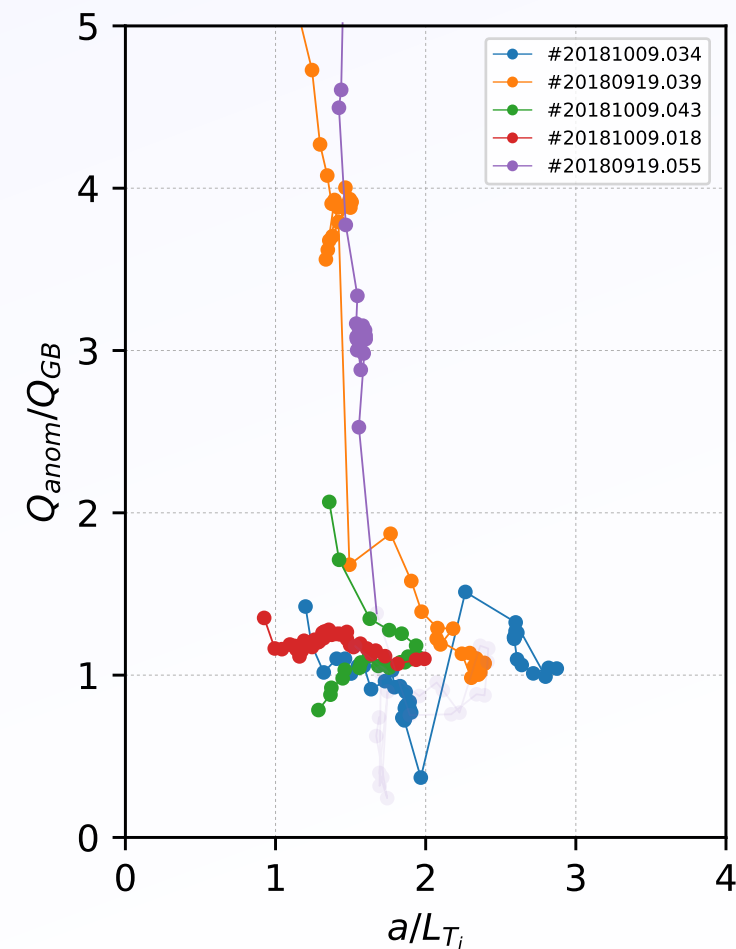
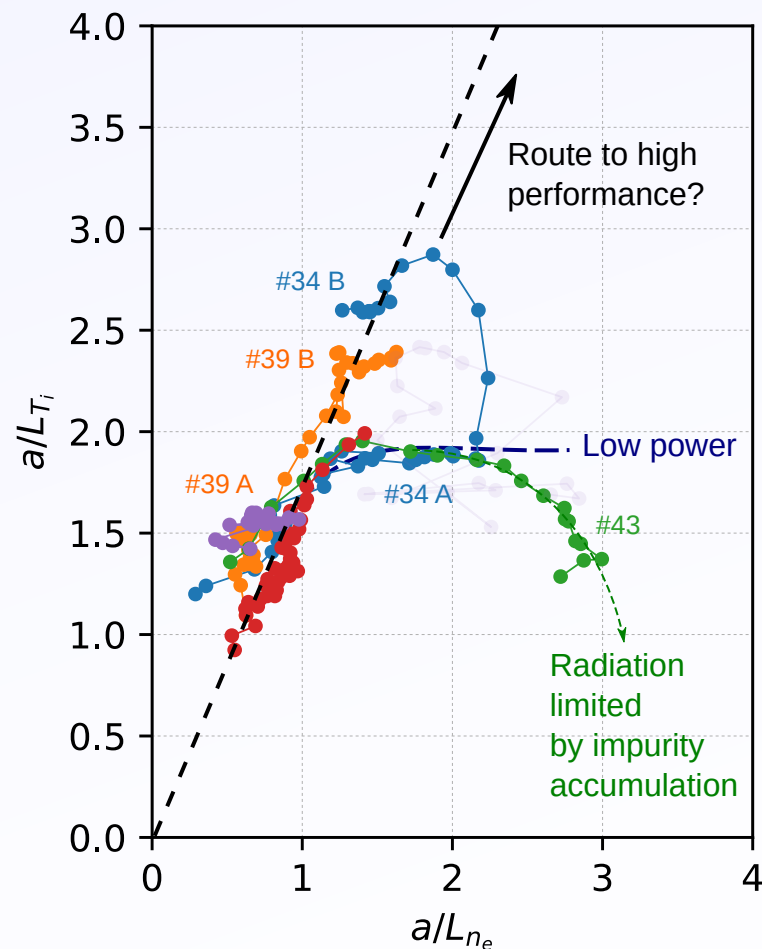
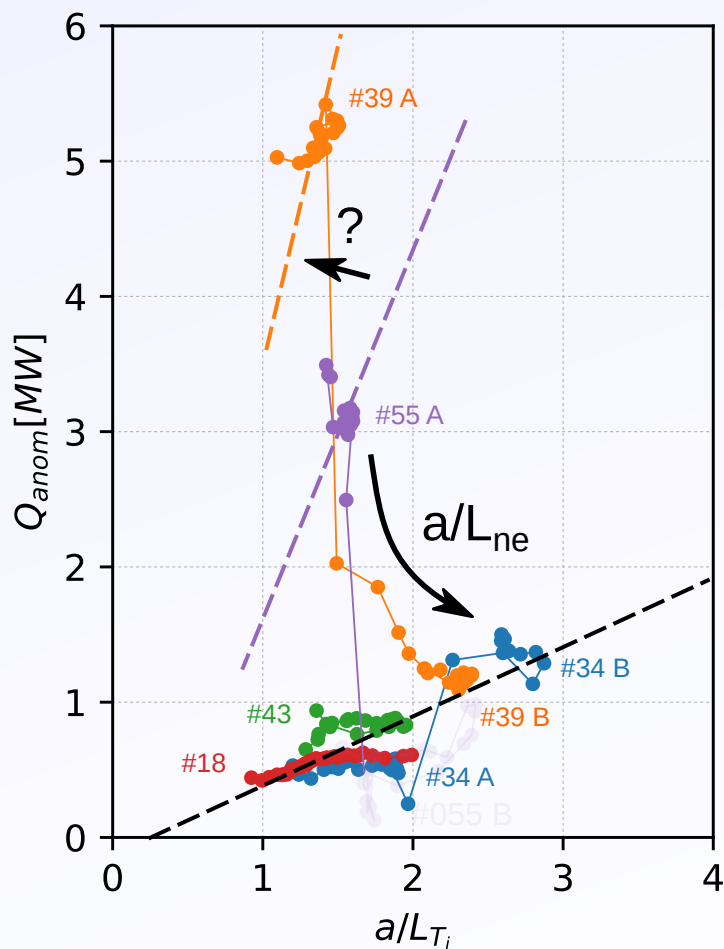




## Total power flux

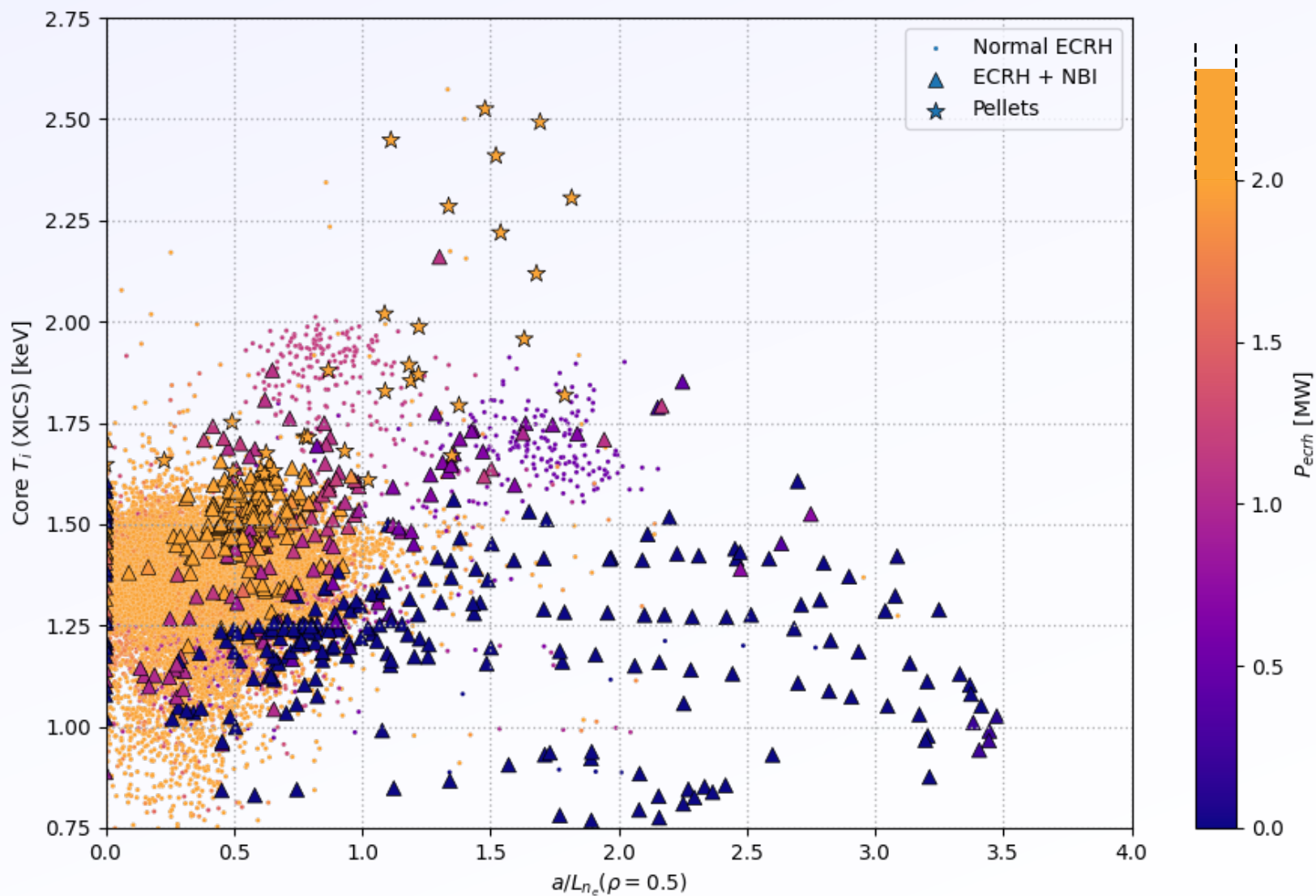
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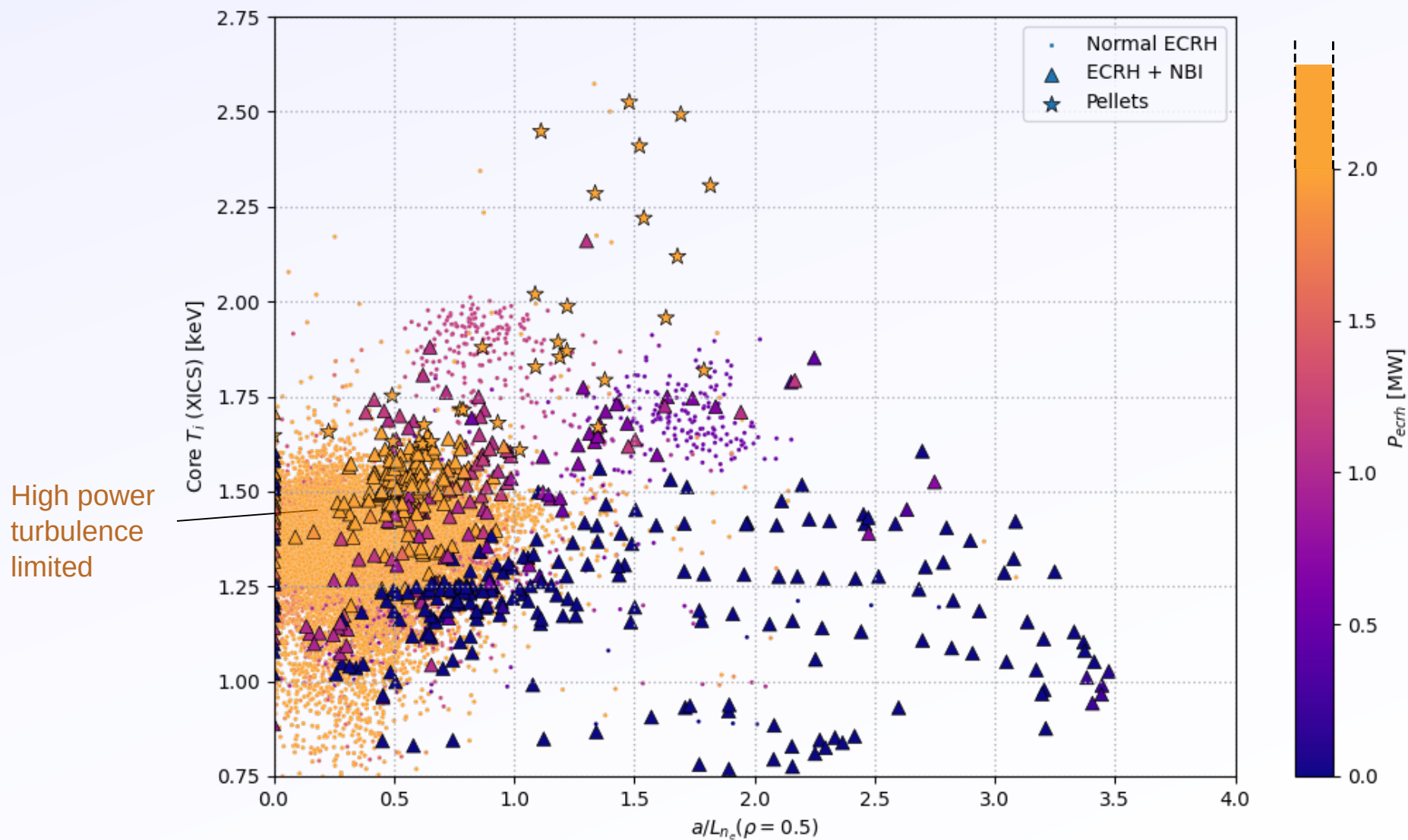
## Confinement vs density gradient

This can also be seen in the global view:



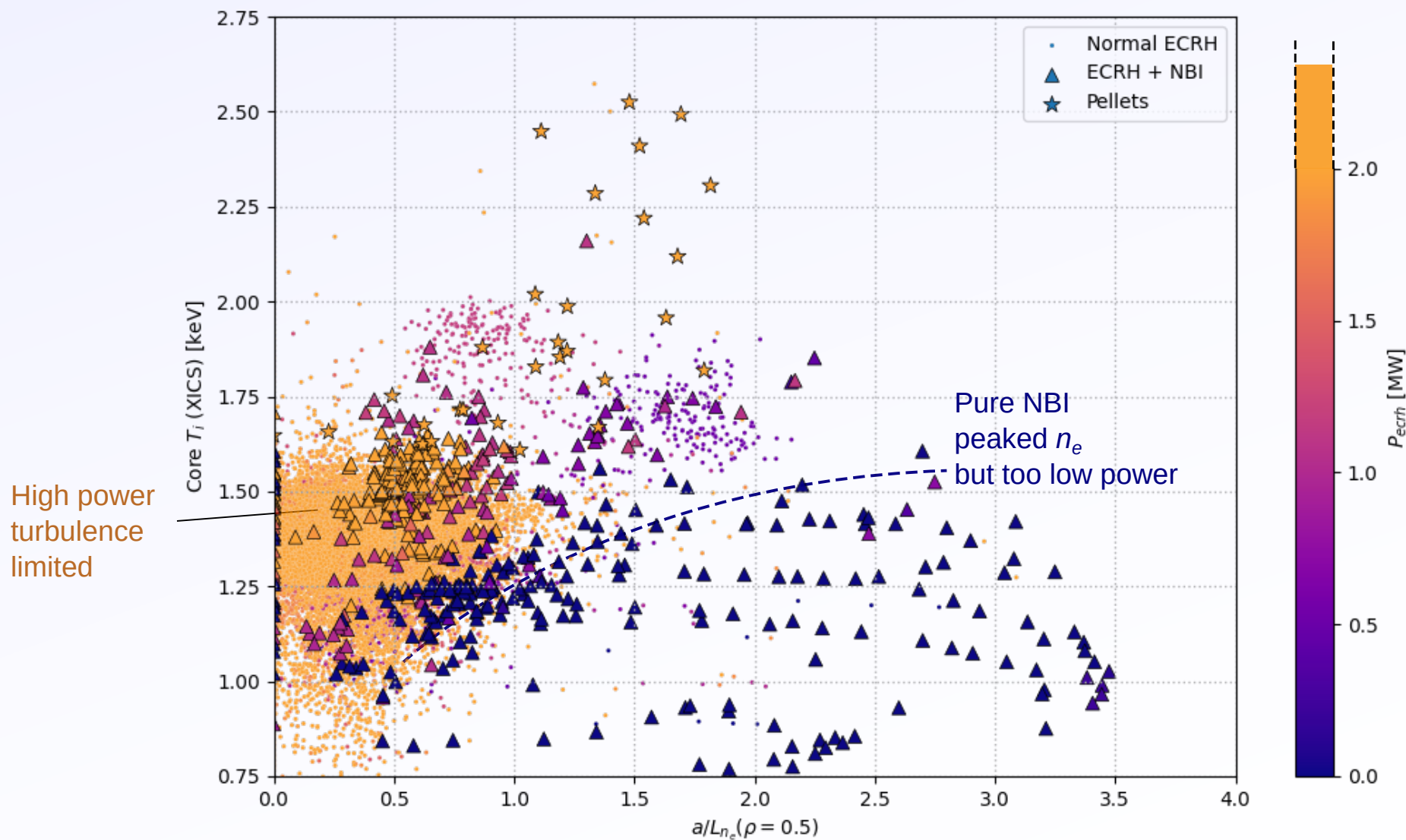
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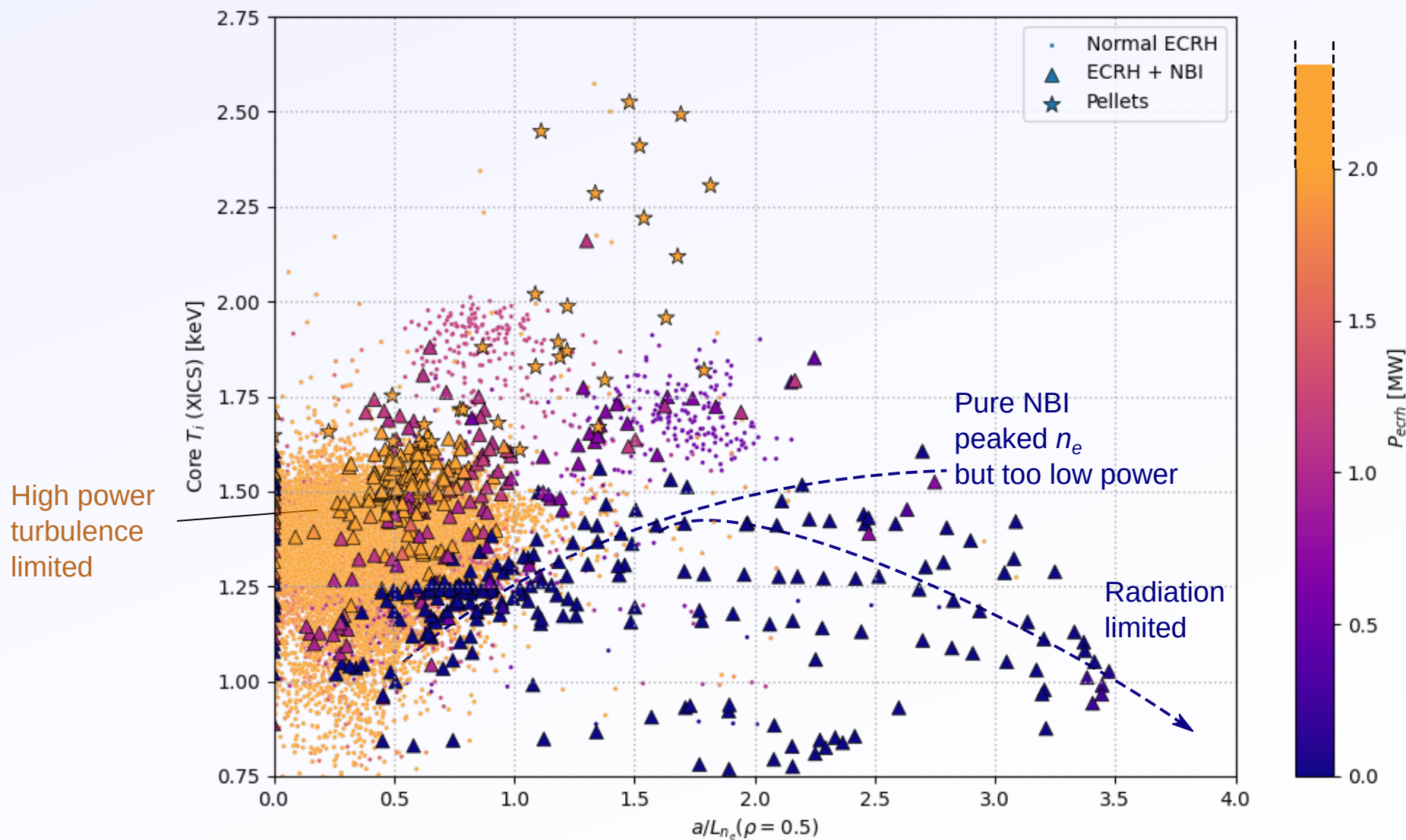
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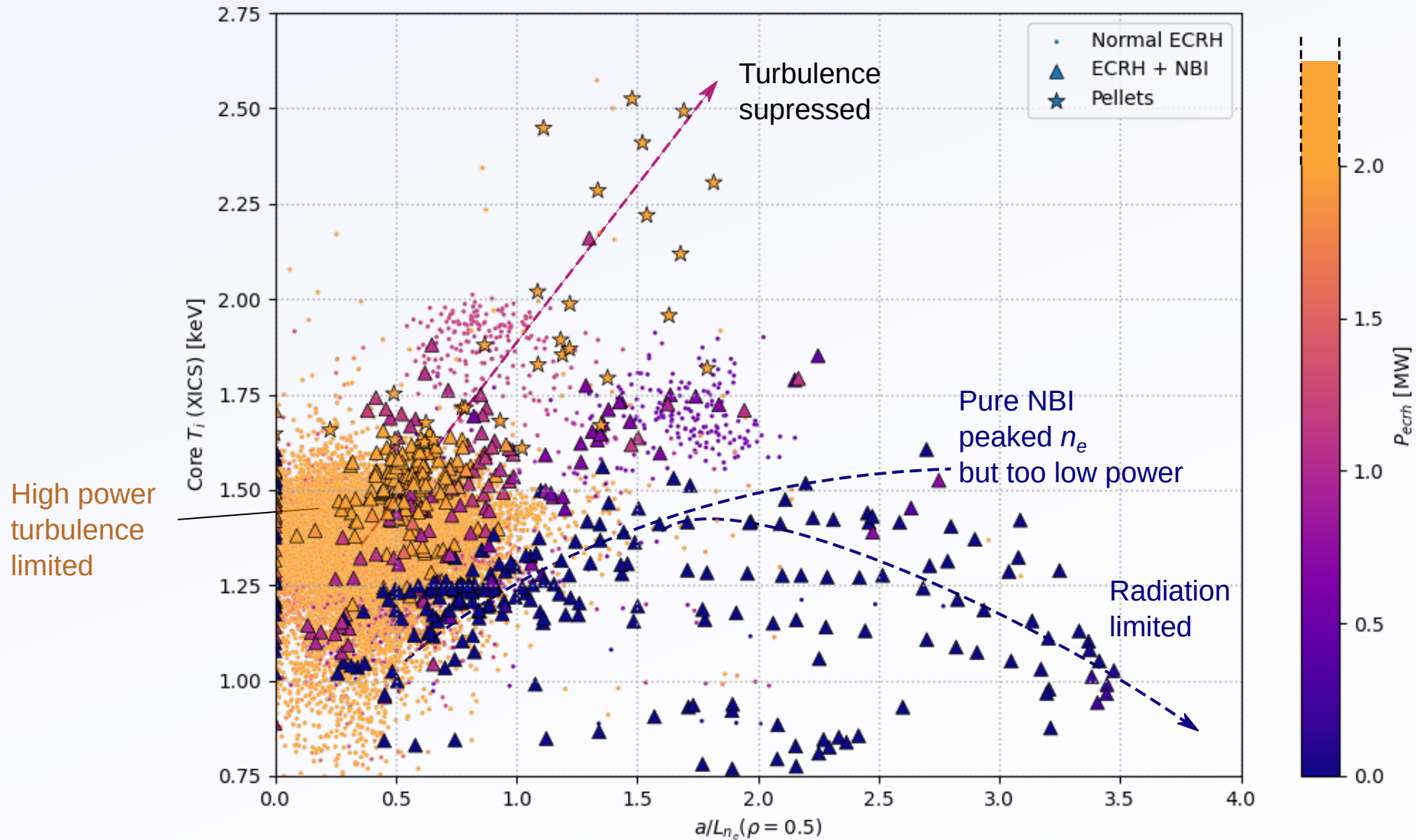




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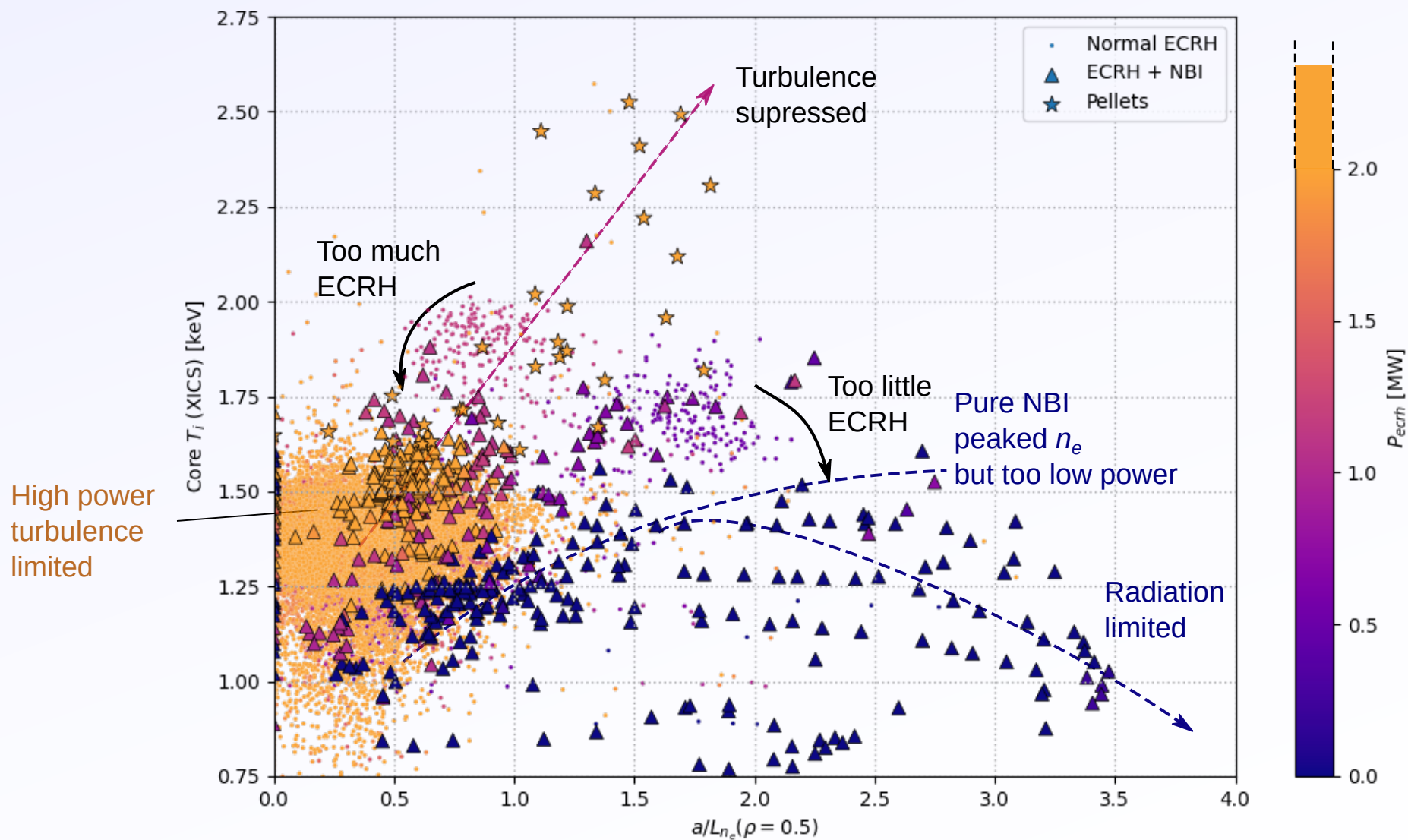
- All higher  $a/L_{ne}$  discharges with a little ECRH move up towards post-pellets HP plasmas.



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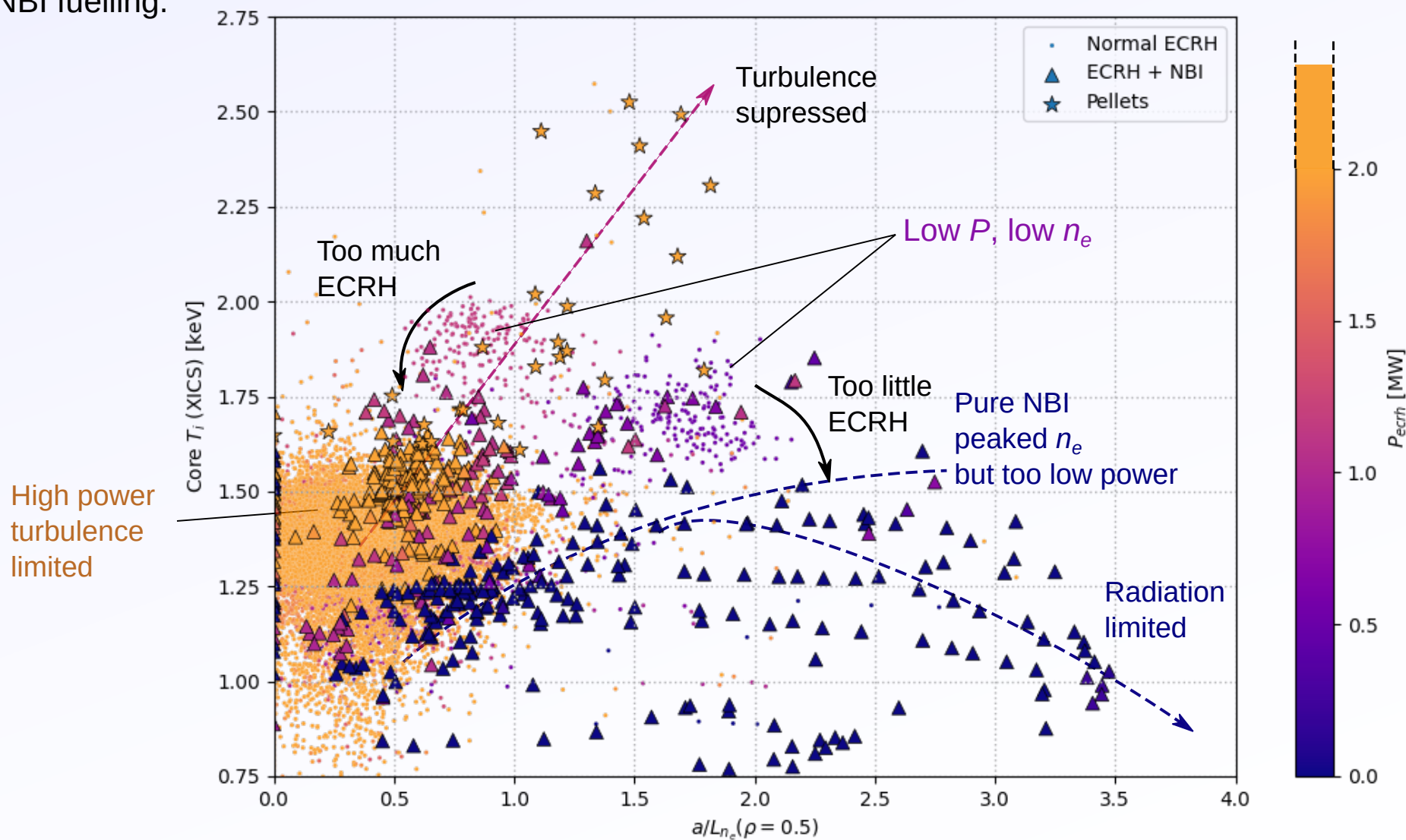




## Confinement vs density gradient

This can also be seen in the global view:

- All higher  $a/L_{ne}$  discharges with a little ECRH move up towards post-pellets HP plasmas.
- Low  $P$ , low  $n_e$  shots without NBI show the same behaviour - density peaking most likely a transport effect and not from NBI fuelling.



## Summary

### Particle transport (in pure NBI only):

- Low net particle flux initially gives slow rise of core density.
- At  $a/L_{ne} \sim 0.8$  anomalous particle flux at/inside  $\rho=0.5$  reduces dramatically --> density peaking + impurity accumulation.

### Heat transport:

- Heat transport at low  $a/L_{ne}$  is consistent with high stiffness in ECRH-only plasmas.
- Heat transport at high  $a/L_{ne}$  is consistent with gyro-Bohm scaling.
- Pure NBI plasmas are limited by the input power --> More power initially gives higher  $T_i$ .
- In pure NBI plasmas, the radiation from impurity accumulation eventually kills the plasma.
- Too much power (at least with ECRH  $> \sim 1.2$  MW) reduces  $a/L_{ne}$  and heat transport degrades dramatically.
- There could be a strong  $Q_i$  barrier at mid-radius... interesting, but probably not very useful.  
--> Should we invest resources to measure it?

### General:

- It *might* be possible to slowly increase the NBI and ECRH power together, such that  $\eta_i = (a/L_{Ti} / a/L_{ne}) \sim 1.75$  is maintained and to follow this path *towards* the post-pellet plasma performance.
- Is it most critically important to understand when the extra ECRH power decreases  $a/L_{ne}$ .  
--> **Study the turbulent particle transport!**



## Open questions

- What causes the low particle flux at  $\rho < 0.5$  in (some) NBI discharges? Is this really density gradient?
- Is there really an ITB in  $Q_i$ ? (Although we probably shouldn't care)
- Why does high ECRH increase the core particle flux?
- What is the 'right amount' of ECRH to flush out impurities and control density rise?
  - If the ECRH needed to control impurities is already enough to lower  $a/Lne$ , then we cannot win.
- What happens when we add more NBI?
  - If the particle fluxes do not increase: Add more ECRH, but this ok, because density gradient will remain. Great!
  - If particles fluxes increase: No way to add power without losing density gradient.

Study other things: beam current, momentum .... etc

Experiments for OP2:

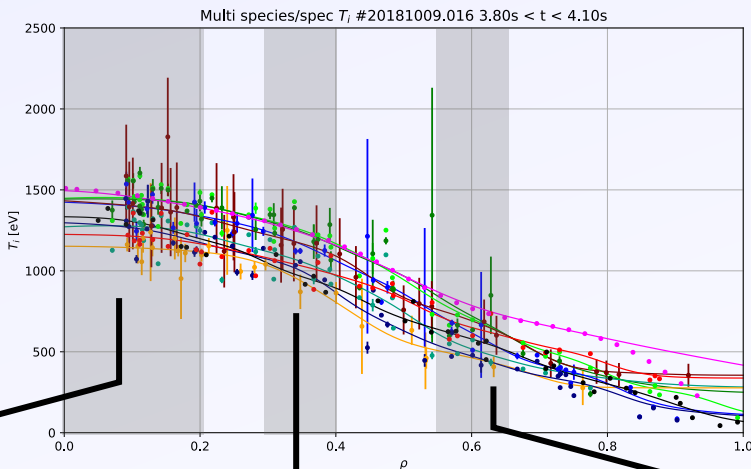
Fine ECRH power steps at several NBI power steps to empirically map:

- Density peaking and flattening with ECRH
- Impurity expulsion
- Profile stiffness at higher ECRH power and behaviour on the border.

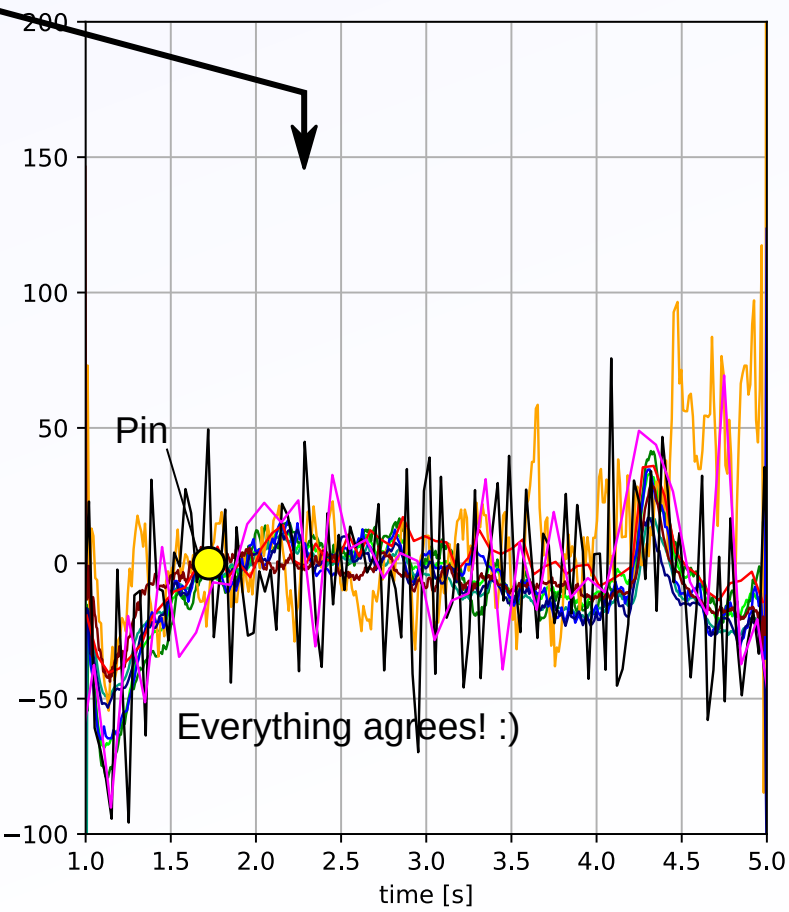
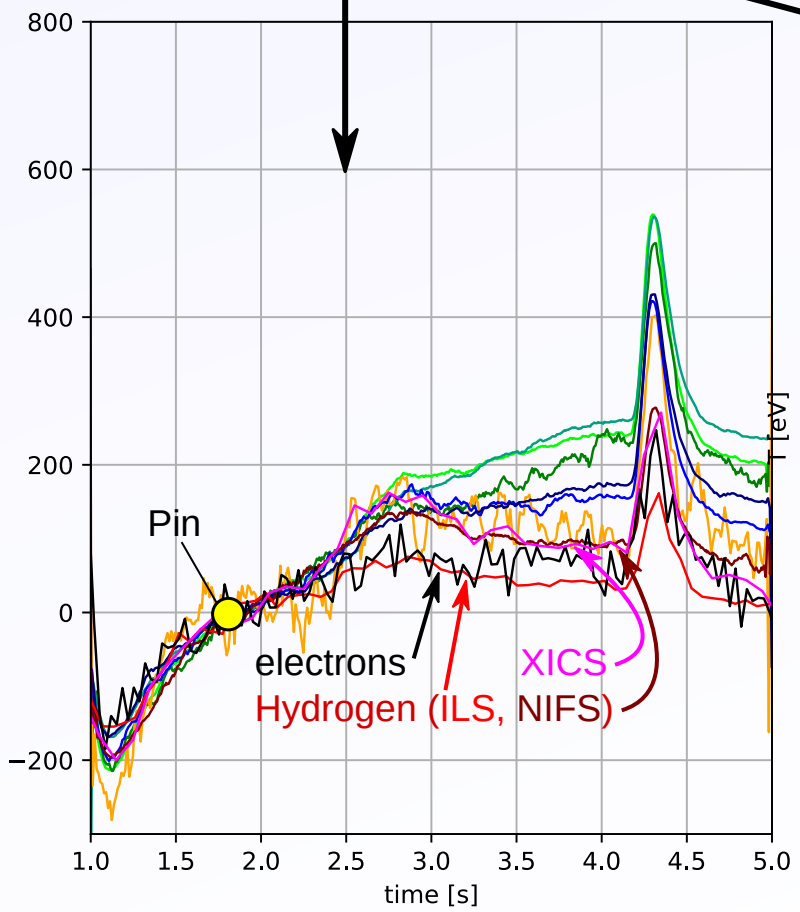
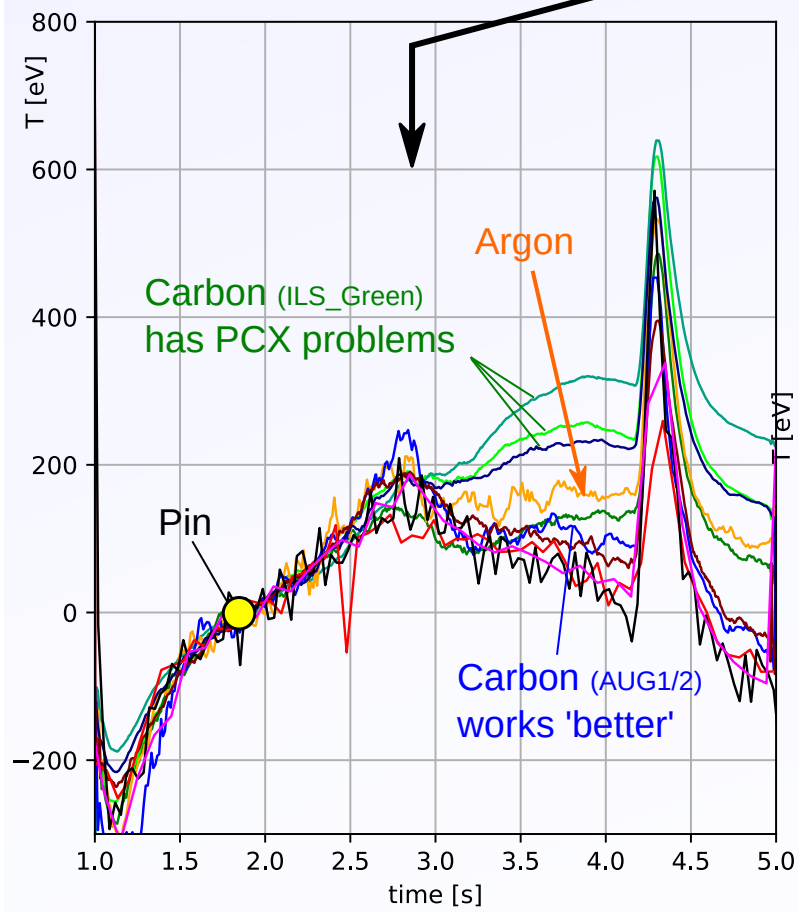


# Te, Ti, Tz profiles

To see what's going on, we can select a region and pin together at one point in time:



- Interpolated C\_VI (ILS\_Green) BGSubtract V3
- Autorange prior C\_VI (ILS\_Green) DualGauss V24
- Include PCX C\_VI (ILS\_Green) DualGauss V25
- Interpolated Ar\_XVI (AUG2) BGSubtract V30
- Autorange prior C\_VI (AUG1) DualGauss V3
- Include PCX C\_VI (AUG1) DualGauss V4
- Hydrogen Halo (NIFS\_H) All V2
- Hydrogen Halo (ILS\_Red) All V1
- Te (TS)
- XICS

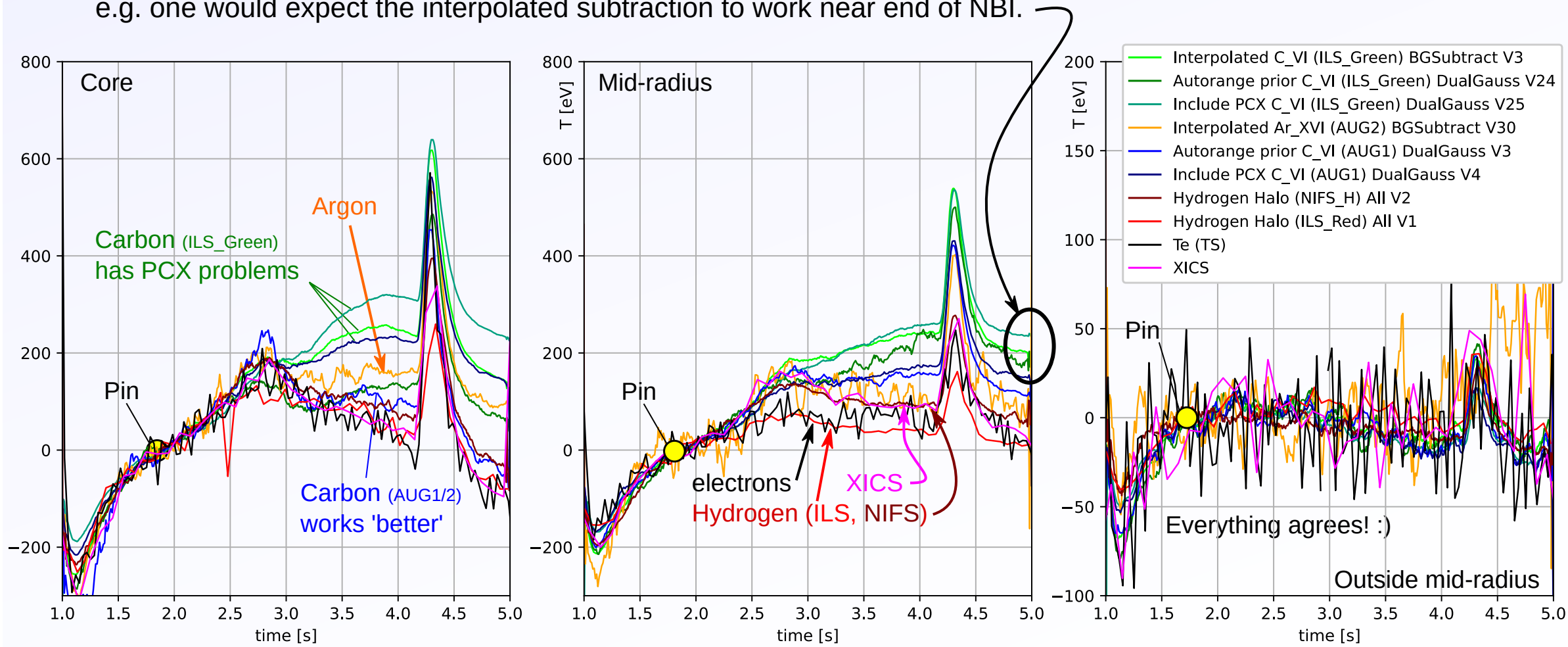




# Te, Ti, Tz profiles

Generally visible:

- Behaviour of  $T_e$  and  $T_i$  (hydrogen) mostly agree --> Expected as  $(T_i - T_e) > \pm 50\text{eV}$  will lead to  $P_{ei} \gg$  available power.
- All temperatures agree outside gradient region, and in the one case where peaking does not occur (#20181009.018)
- XICS Tz in very core seems to agree with Te, Ti, but maybe shows similar higher Ti in steepest gradient region near mid-radius.
- Passive CX ... is a big complication, but the doesn't quite seem to fit.  
e.g. one would expect the interpolated subtraction to work near end of NBI.

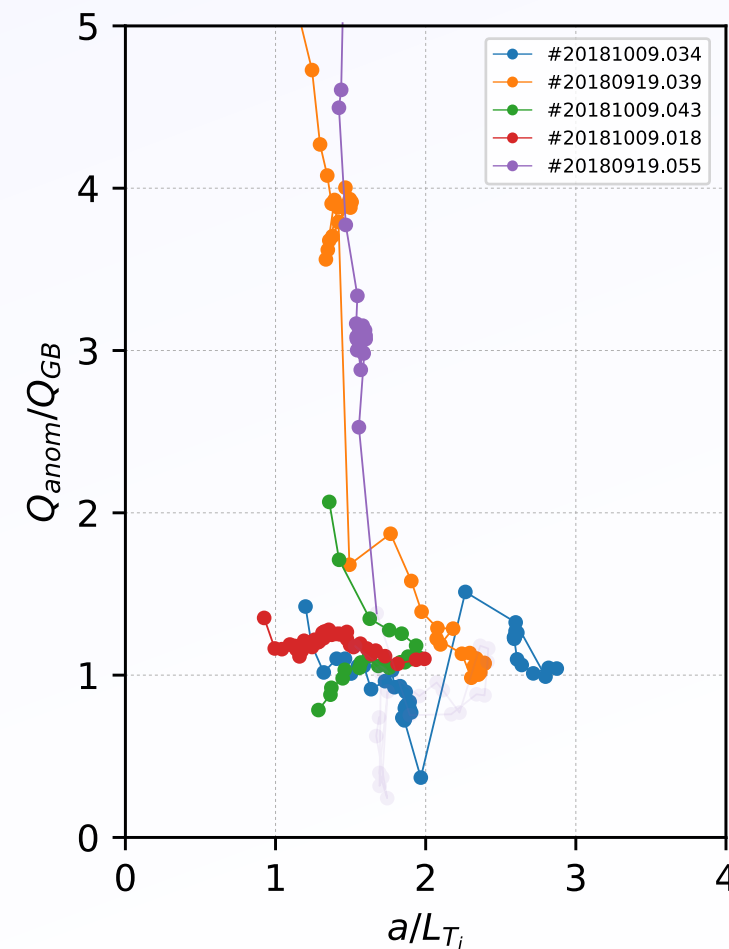
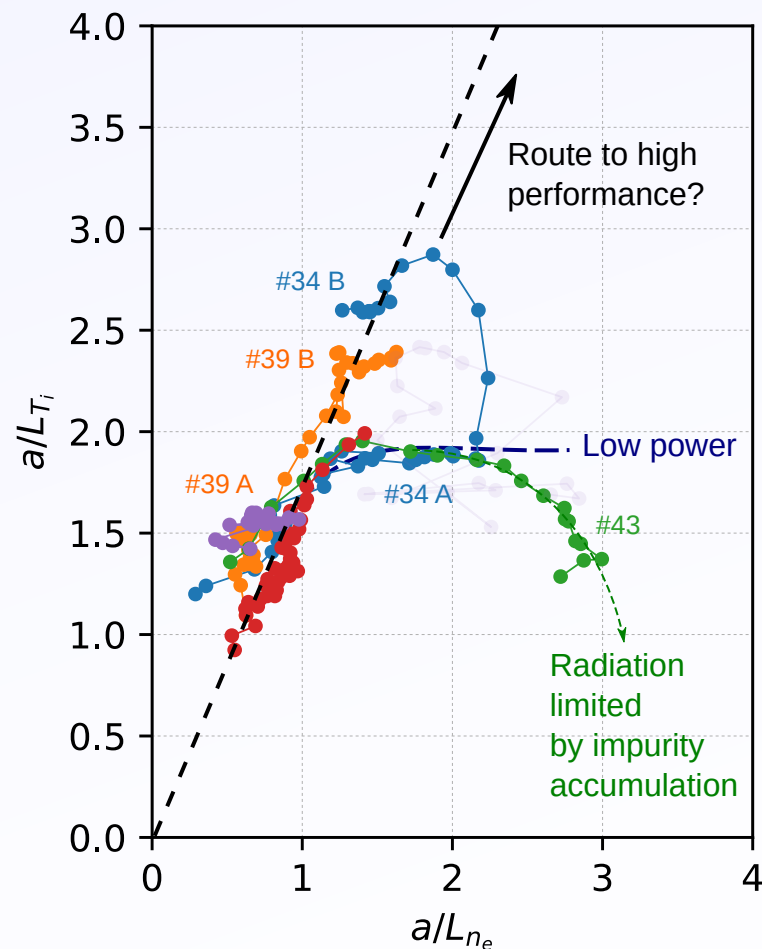
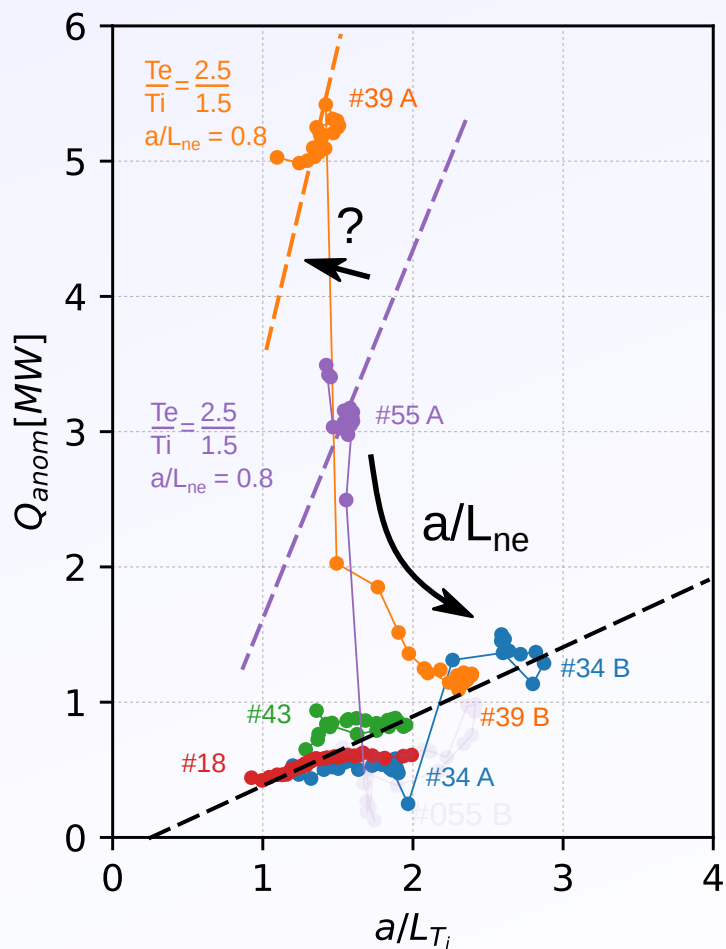




## Total power flux

Over multiple shots, a pattern emerges:

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- #55: Initially high ECRH like #39 but for some reason same  $T_e, T_i, n_e$  at less power...?
- #18: Pure NBI. Density gradient builds up late and power is not added.







...

...



# Total power balance

We can ignore  $P_{ei}$  and examine the balance of  $Q_{total}$ :

- Dominated by anomolous transport at all time points for  $\rho \geq 0.3$ .
- Both radiation and (neo)classical transport are small but significant (classical  $\sim 20\%$  x neoclassical).
- Significant increase in anomolous transport with ECRH.
- Possibly a very weak sign of very small temporary decrease in heat transport coincident with particles.

