



Motional Stark Effect Coherence Imaging for ASDEX Upgrade.

Design and evaluation progress.

O. P. Ford,¹ J. Howard,² R. König,¹ J. Svensson,¹ R. Wolf¹

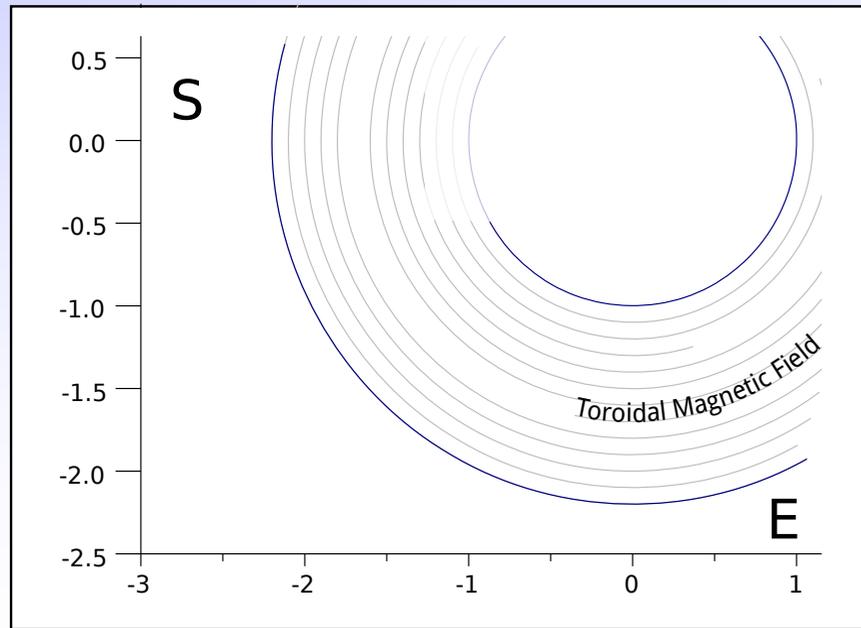
1: Max-Planck Institut für Plasmaphysik, Greifswald, Germany

2: Plasma Research Laboratory, Australian National University, Canberra

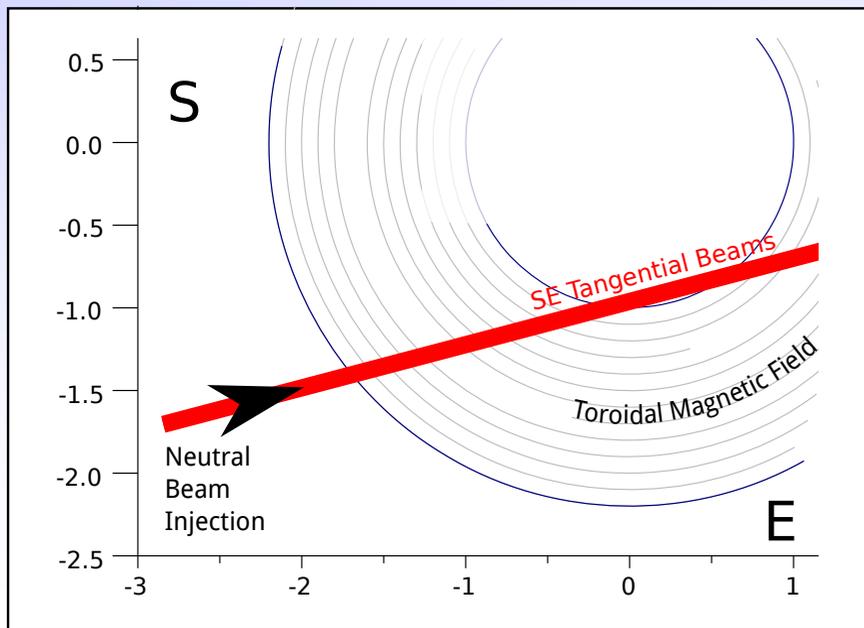
- Brief (re)introduction.
- Line integration and resolution.
- What do we gain with 2D measurements?
 - in Theory,
 - in Practice
- Model improvements.
- Outlook



Very Brief Introduction

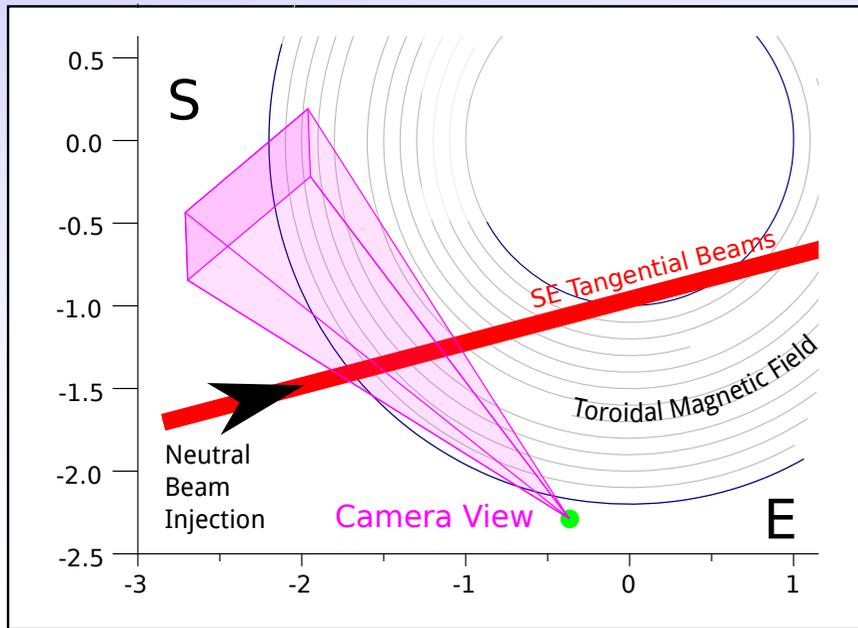


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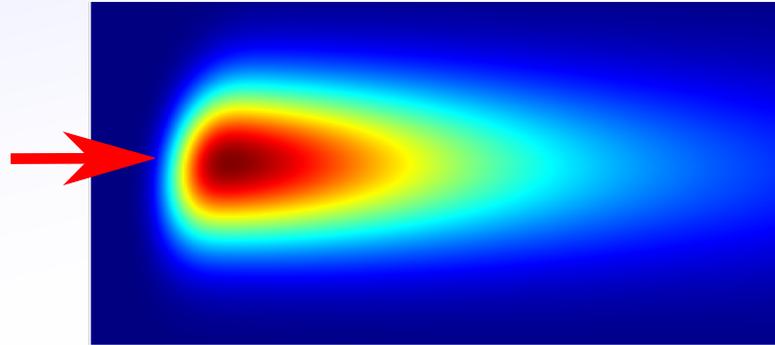


Neutral beam atoms injected into plasma.
Excited by plasma, then emit $H\alpha/D\alpha$ radiation.

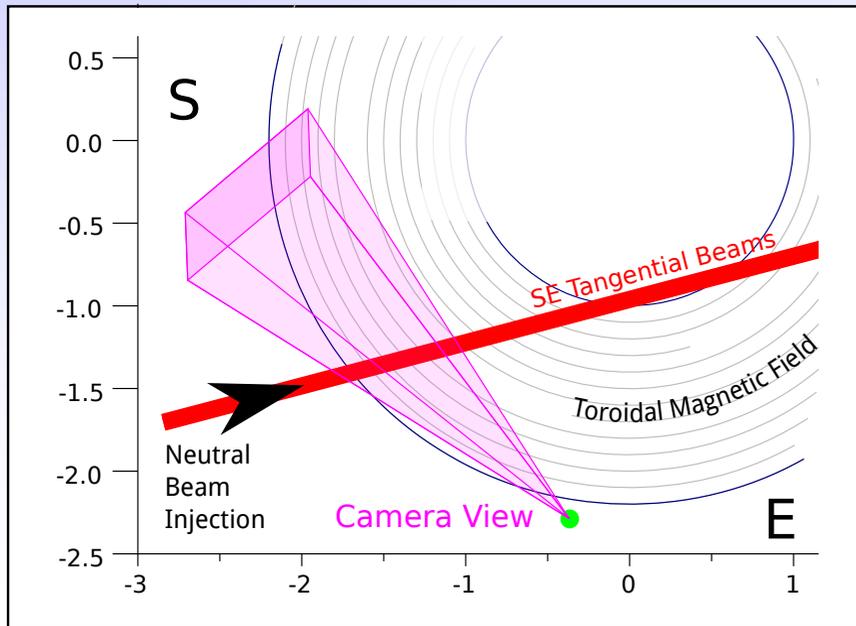
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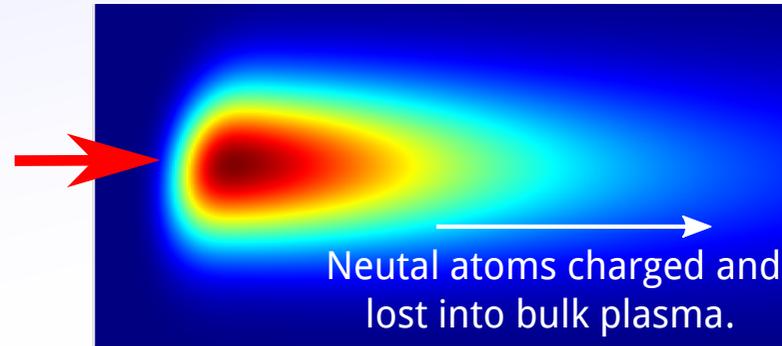
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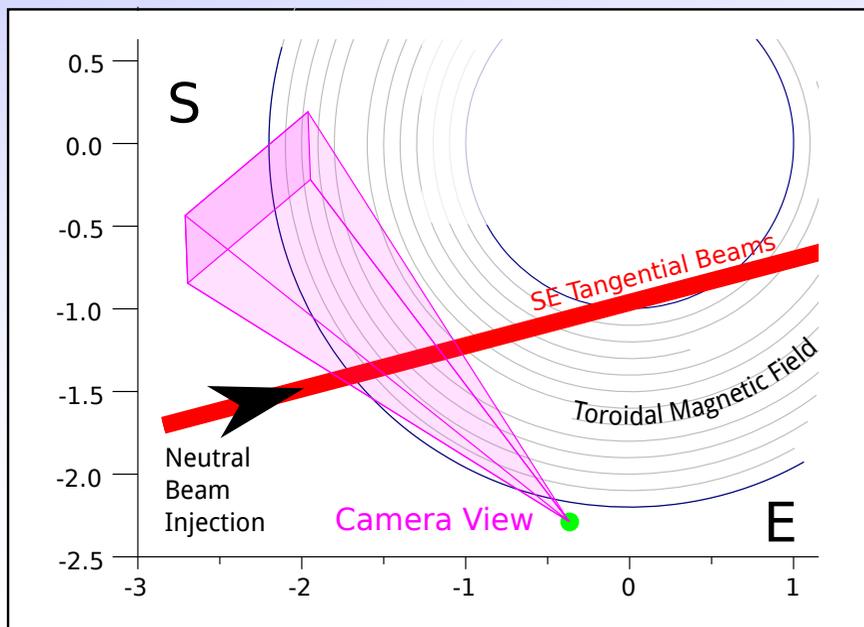
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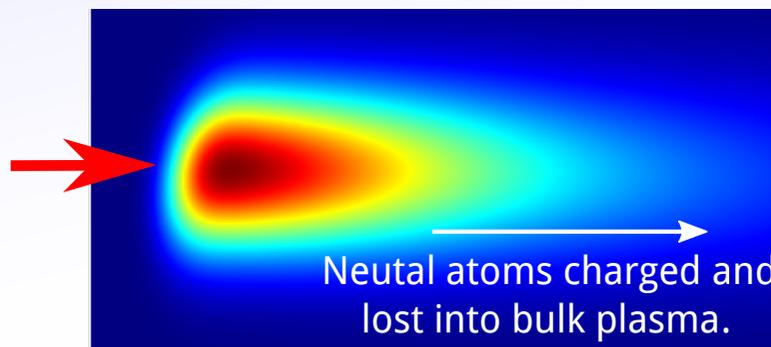
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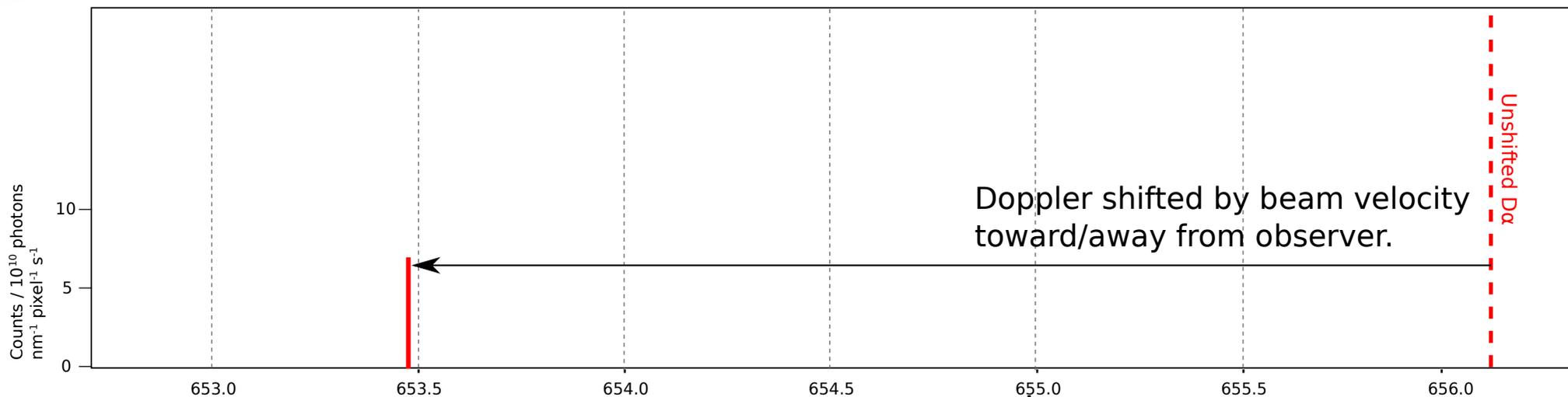
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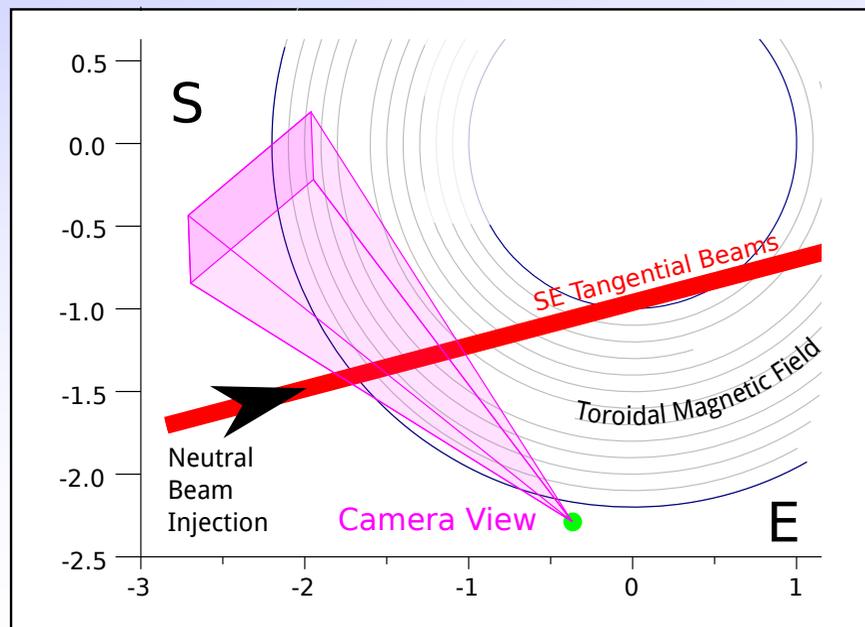
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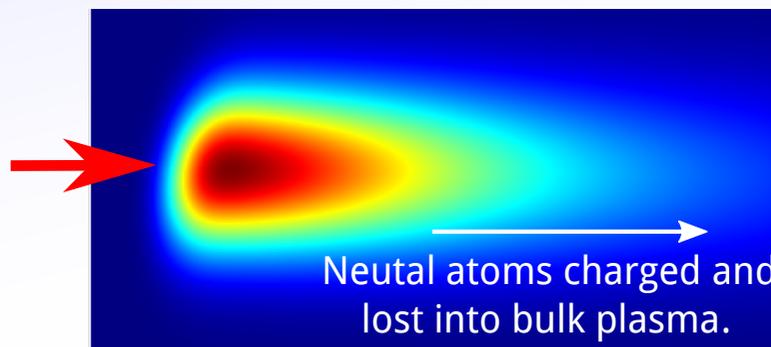
Spectrum from a single pixel:



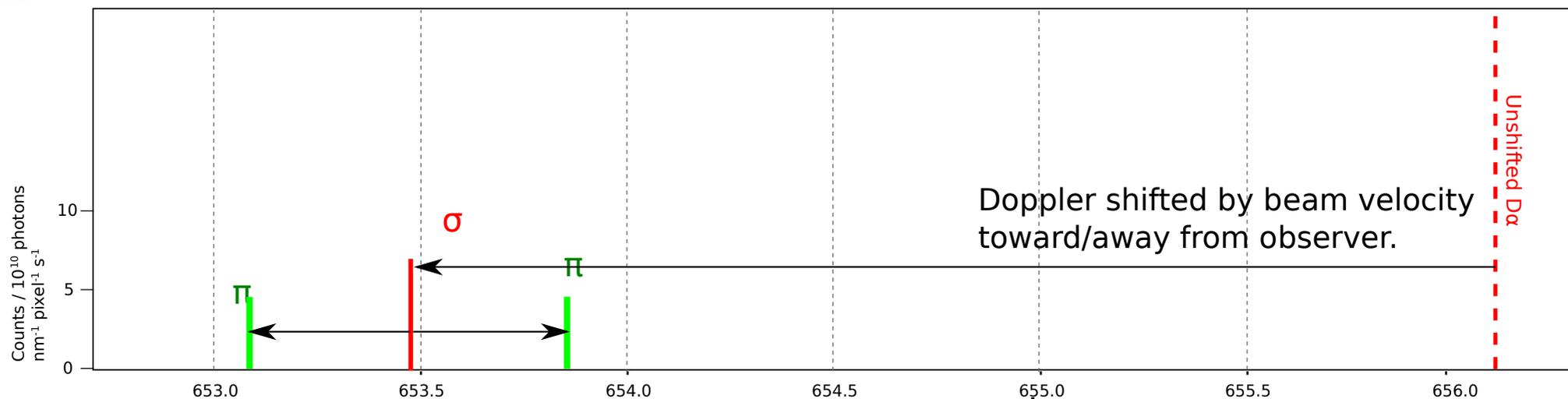
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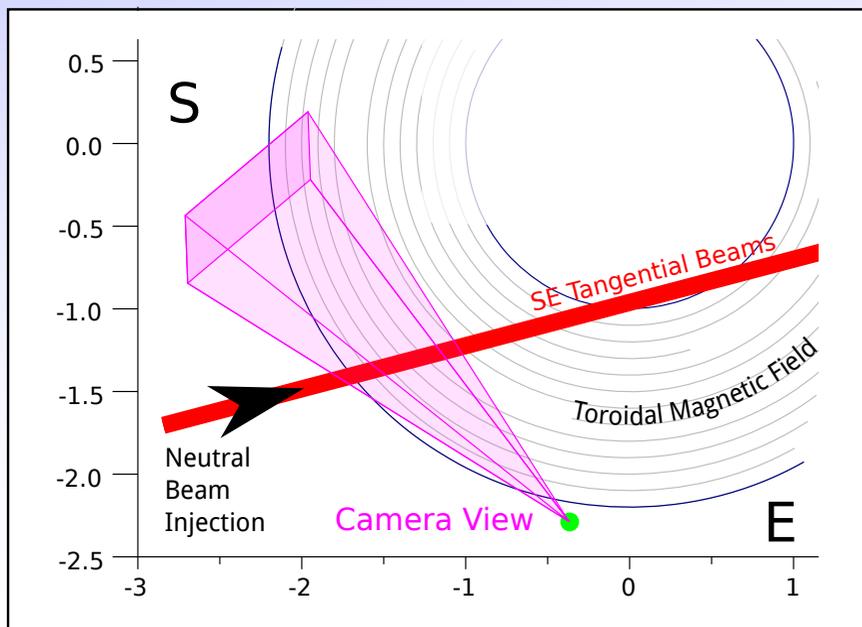
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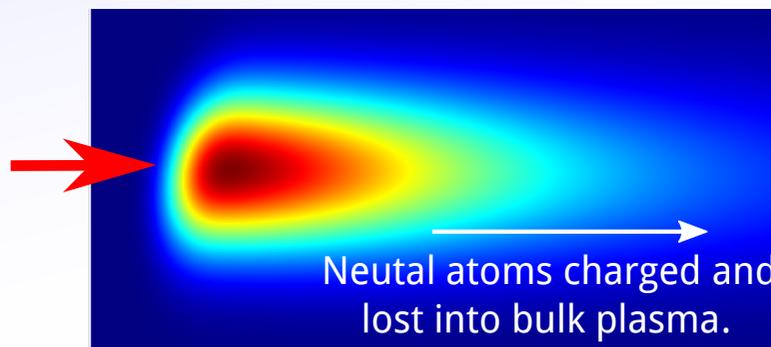
Stark split by electric field in rest frame of atom:

$$\mathbf{E} = \mathbf{v} \times \mathbf{B}$$

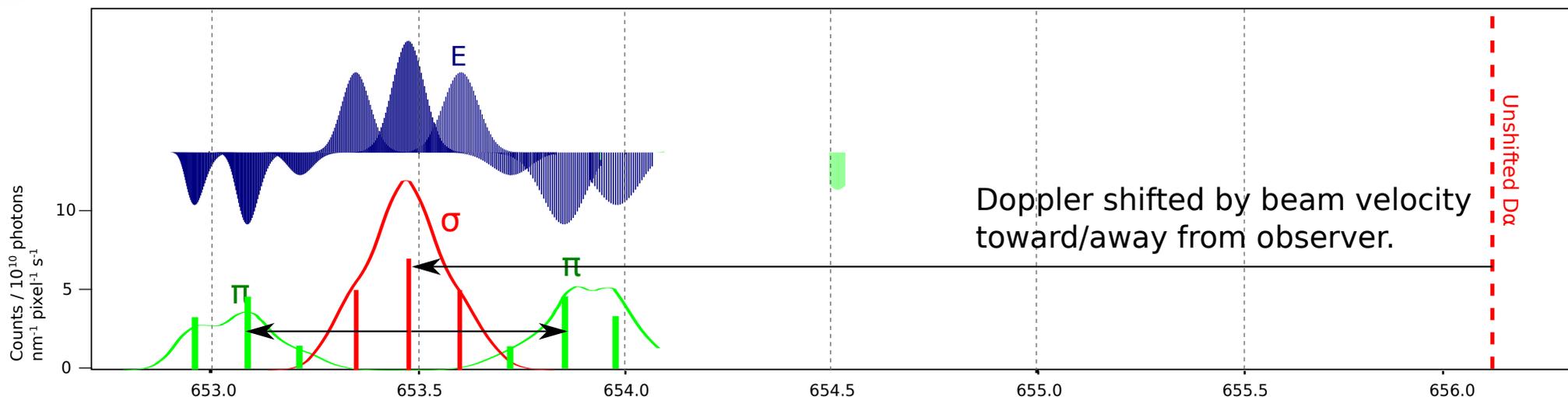
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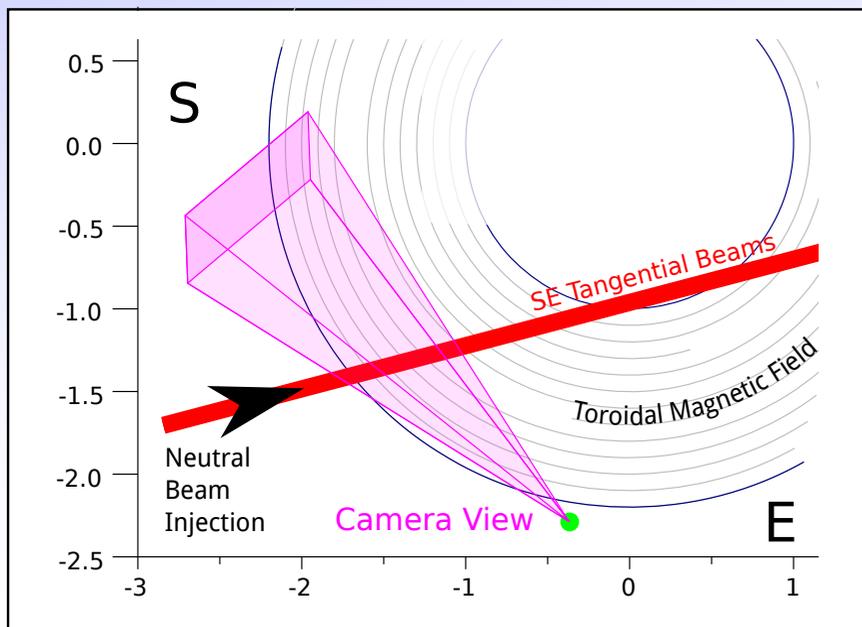
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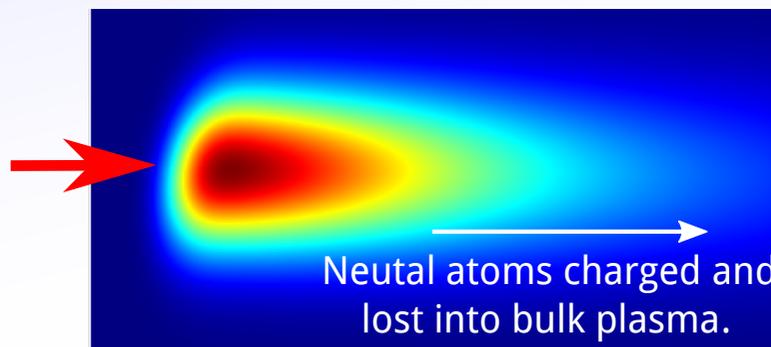
Roughly: π polarised parallel to E.

σ polarised perp' to E.

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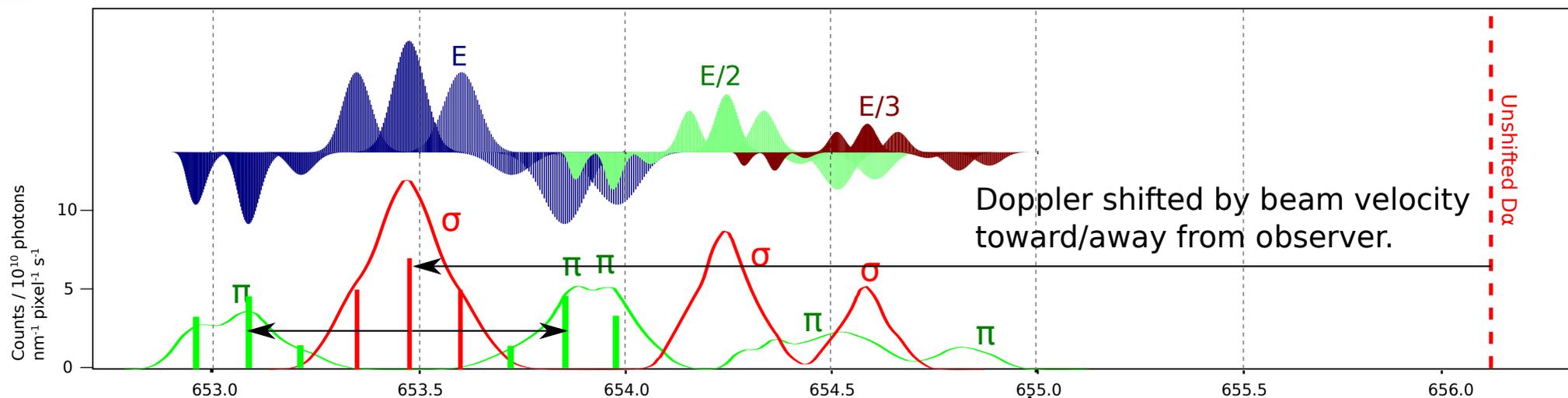


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Complications:
Atoms with different injection energy: different Doppler shift.

Spectrum from a single pixel:



Doppler shifted by beam velocity toward/away from observer.

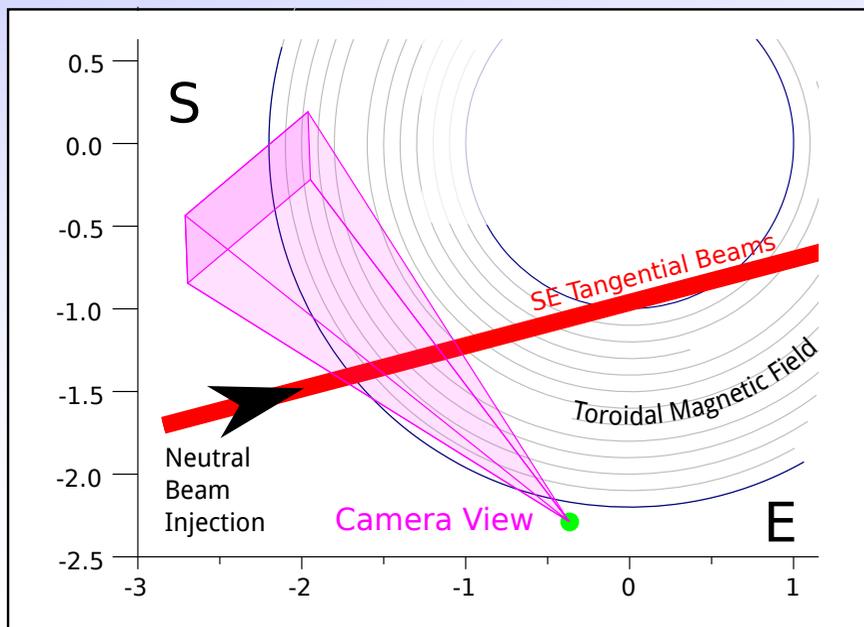
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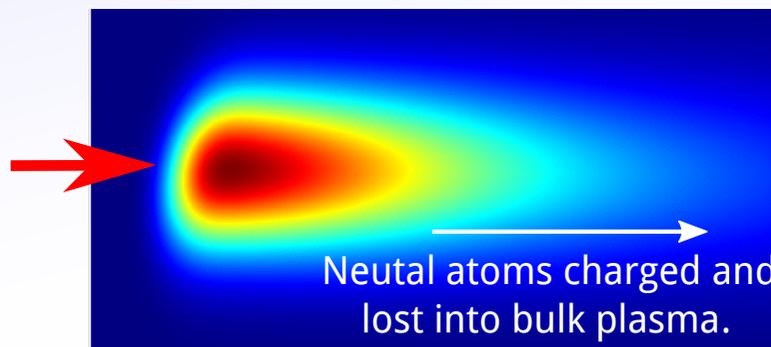
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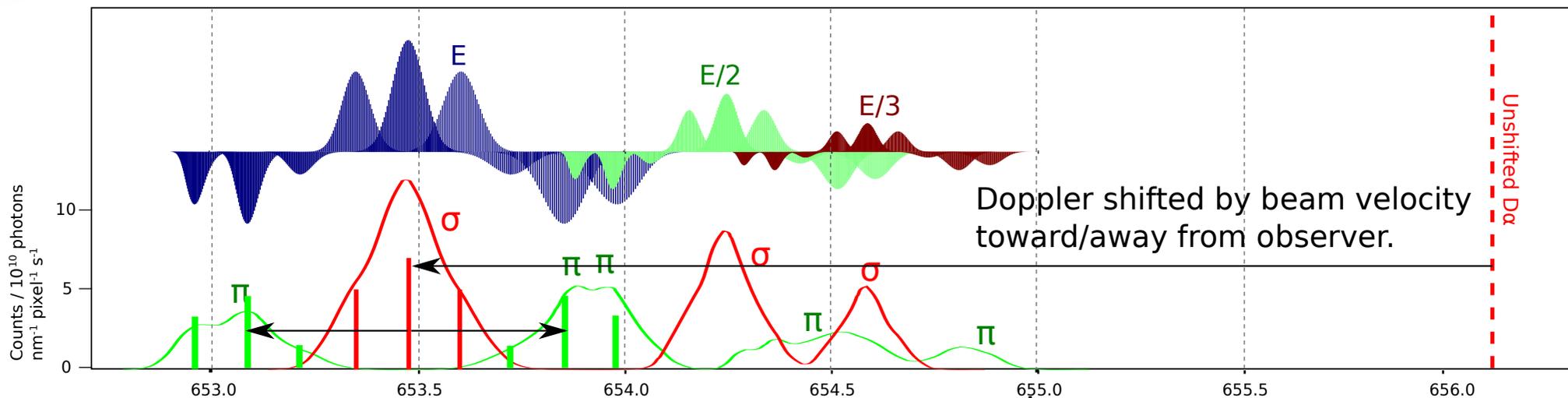
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Doppler broadening: Beam divergence, line integration etc.
Background D α (not shown).

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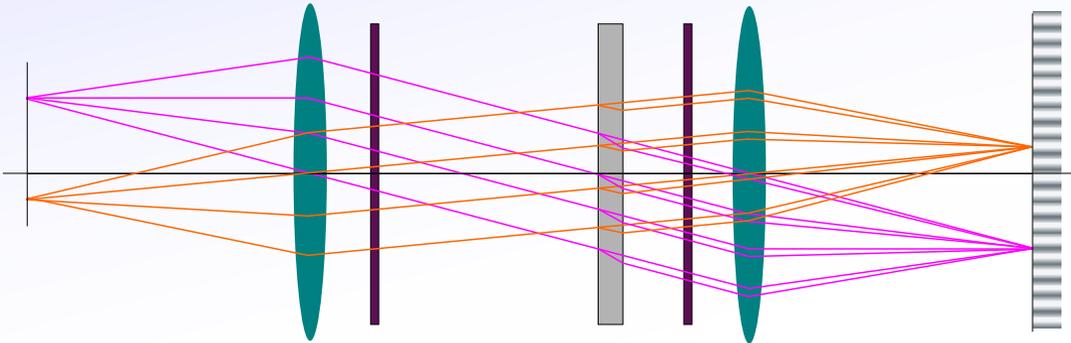
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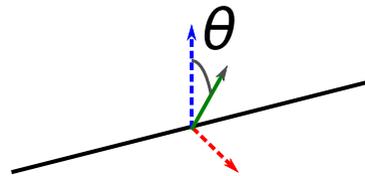
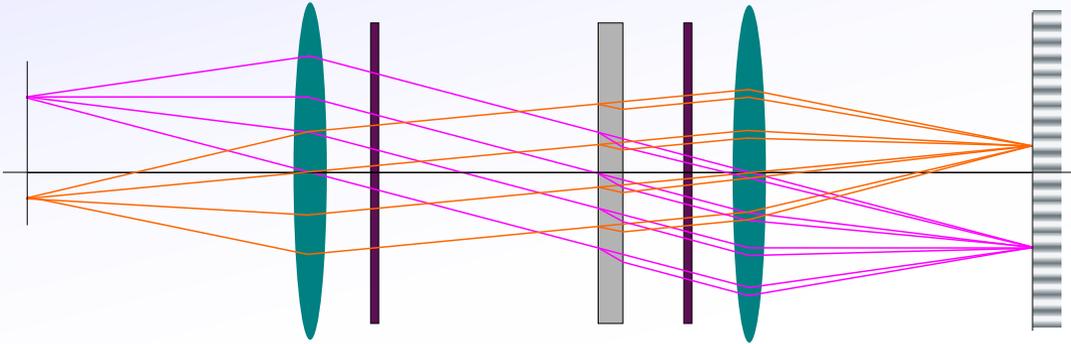
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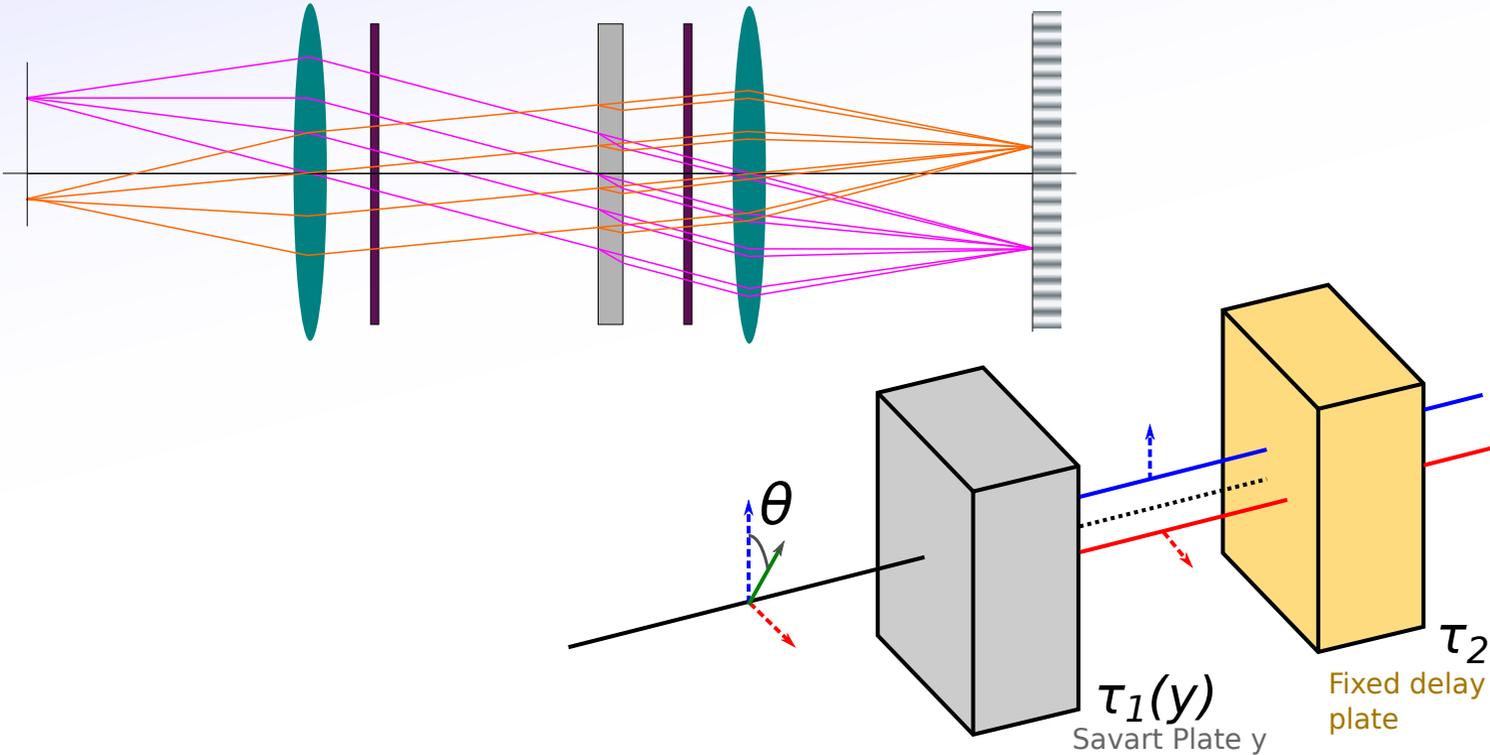
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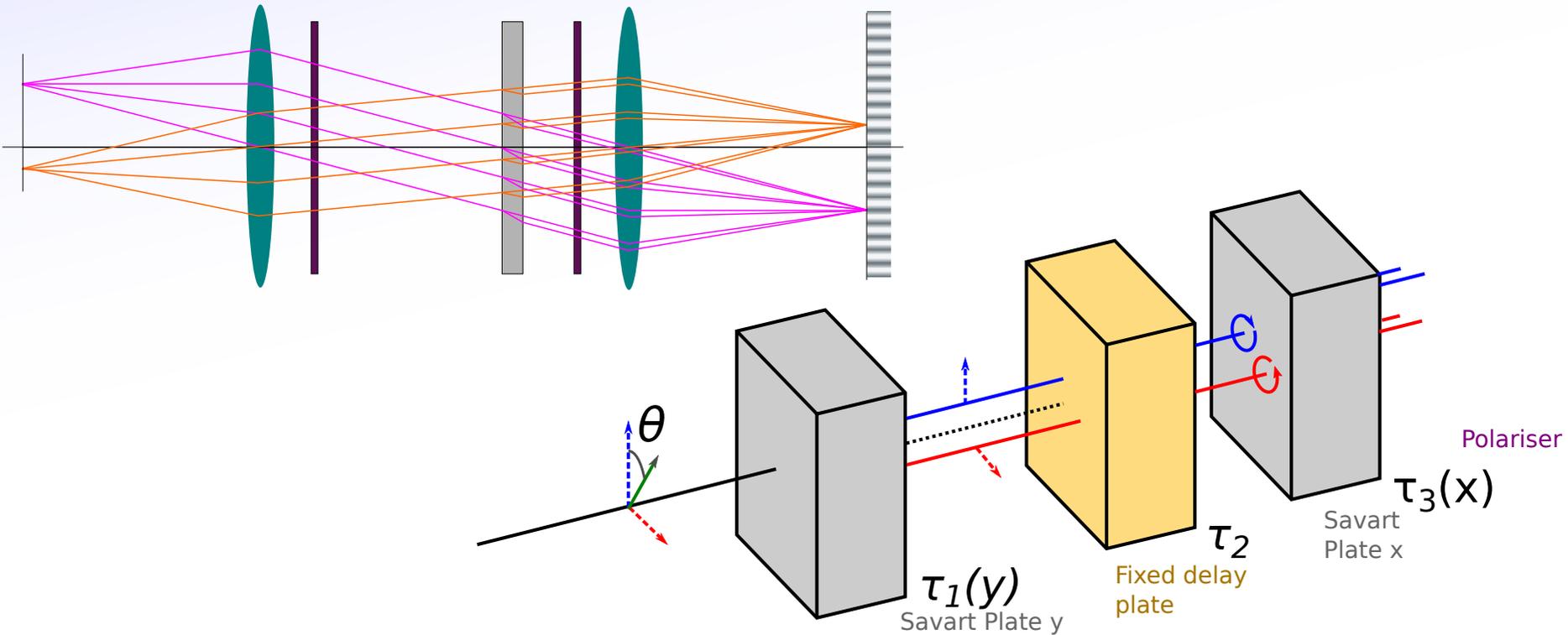
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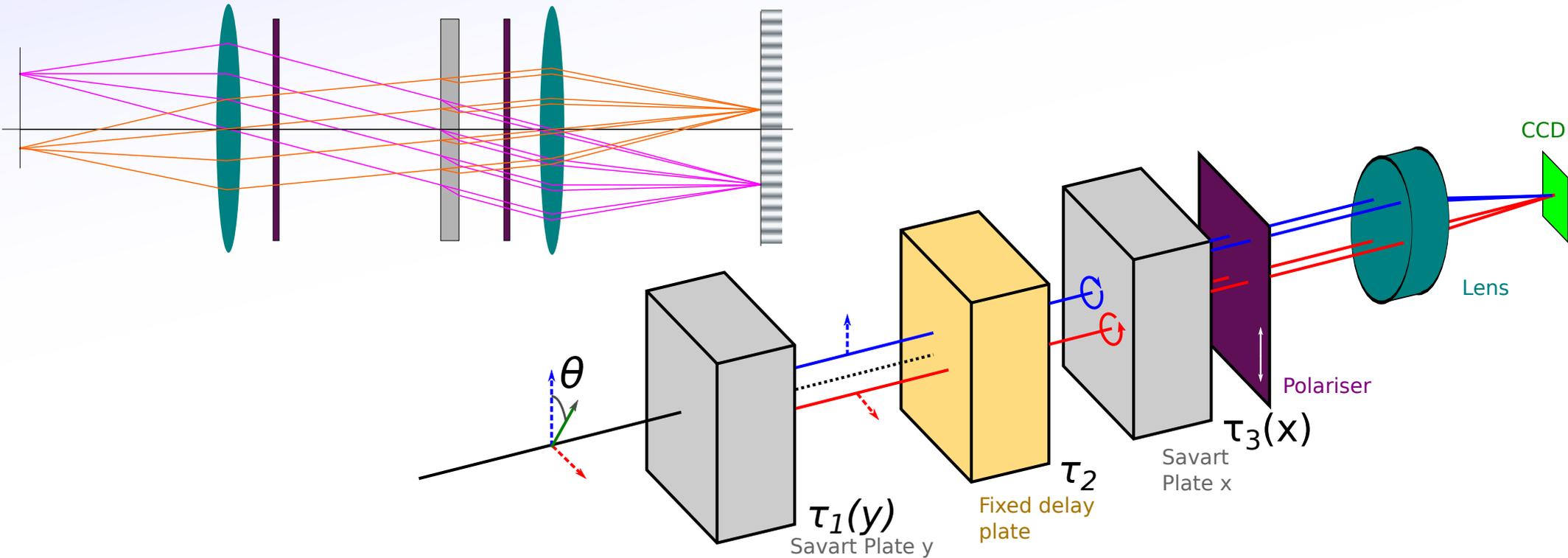
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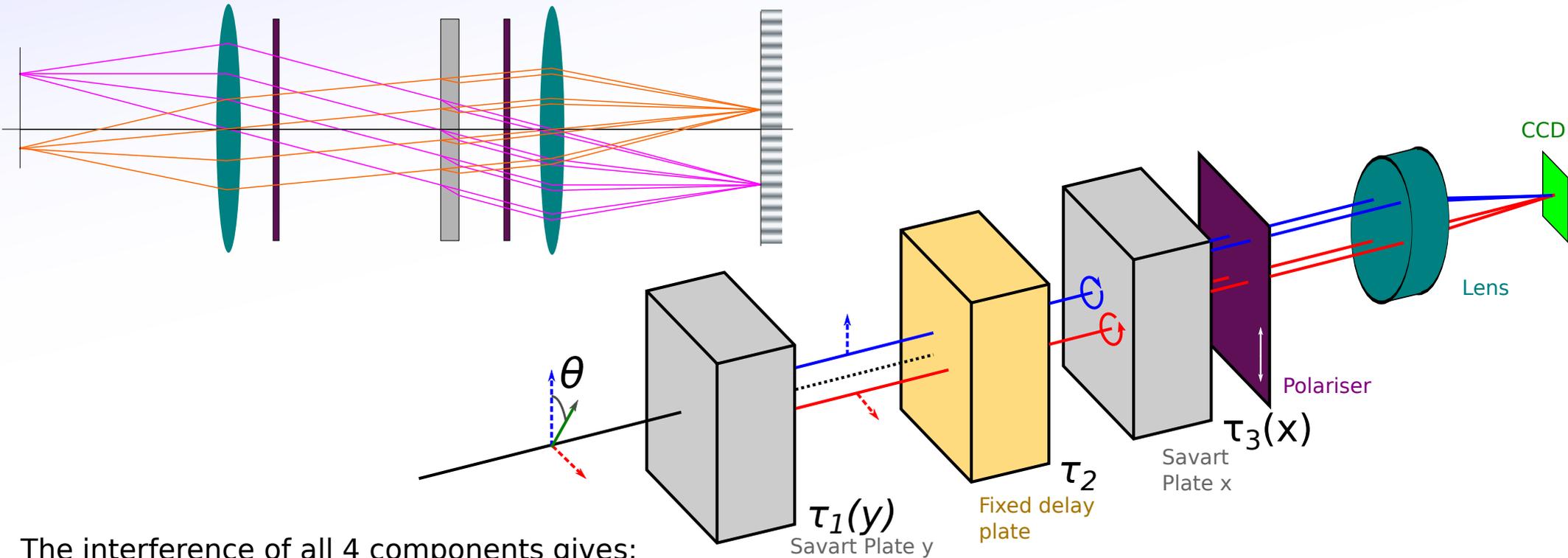
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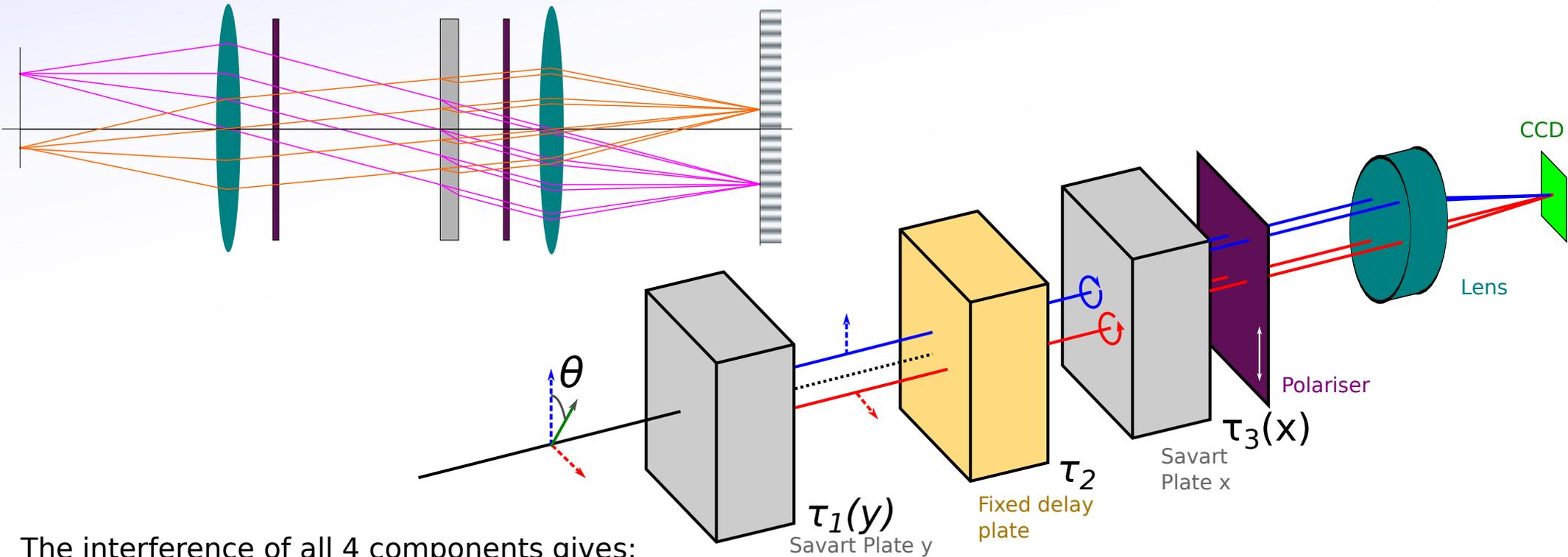
The interference of all 4 components gives:

$$I = \frac{I_0}{2} [1 + \zeta (\cos 2\theta \cos(x) + \sin 2\theta \sin(x) \sin(y))]$$

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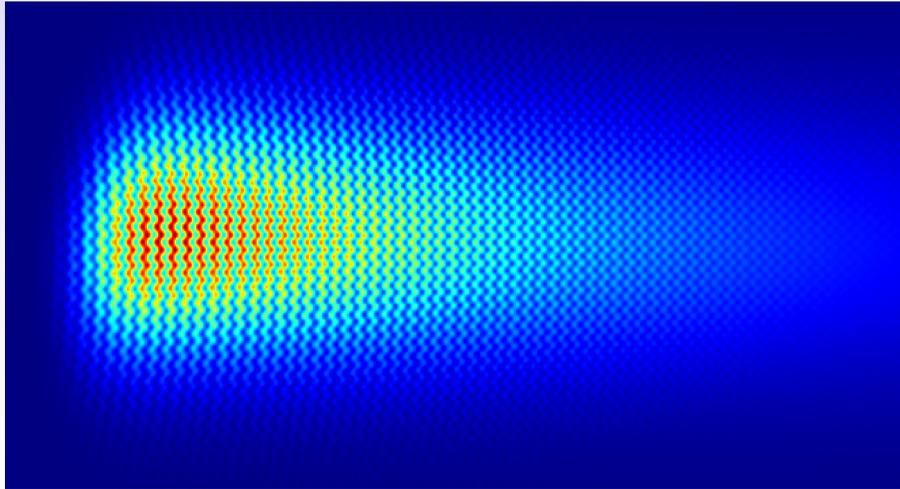
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By demodulating the image in x and y, we can find θ , I_0 and ζ .

(For the record: This is the 'Amplitude Modulated Double Spatial Hetrodyne' system).

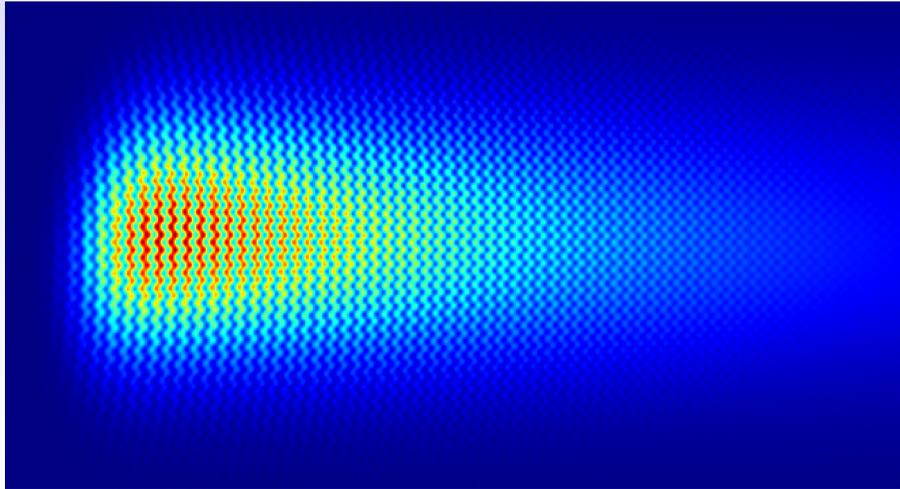
Simple Demodulation

The image will look something like this:

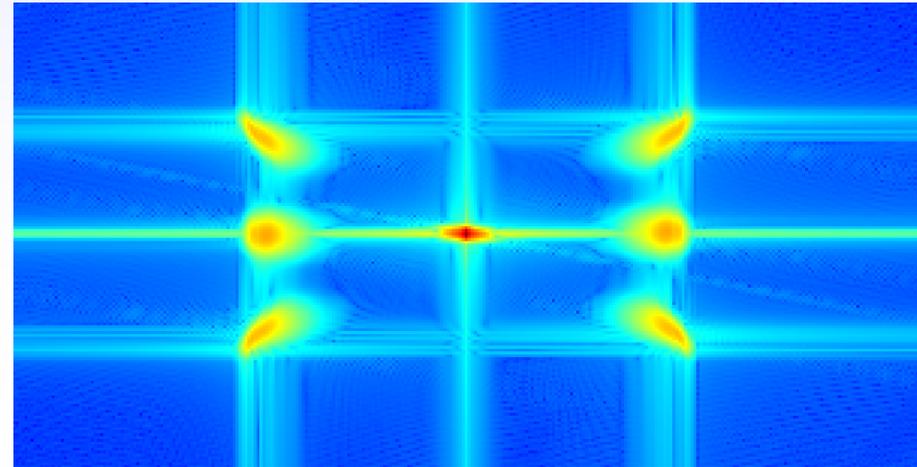


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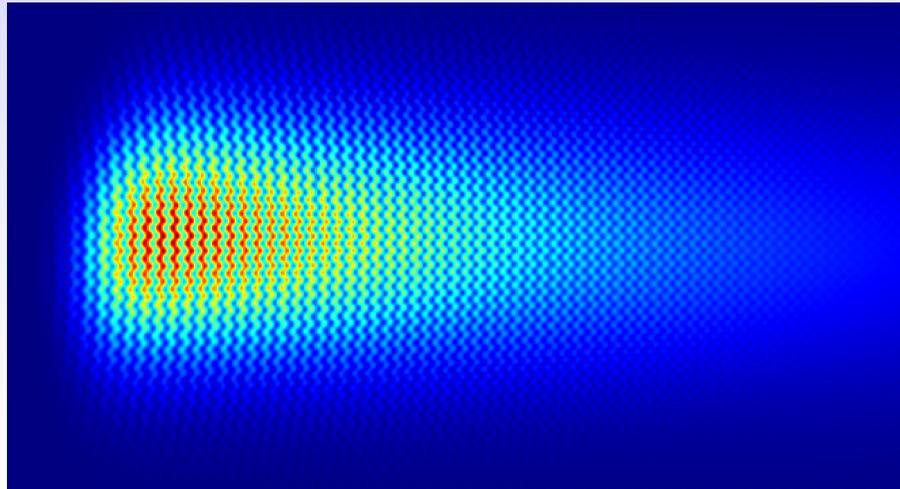


FT
→

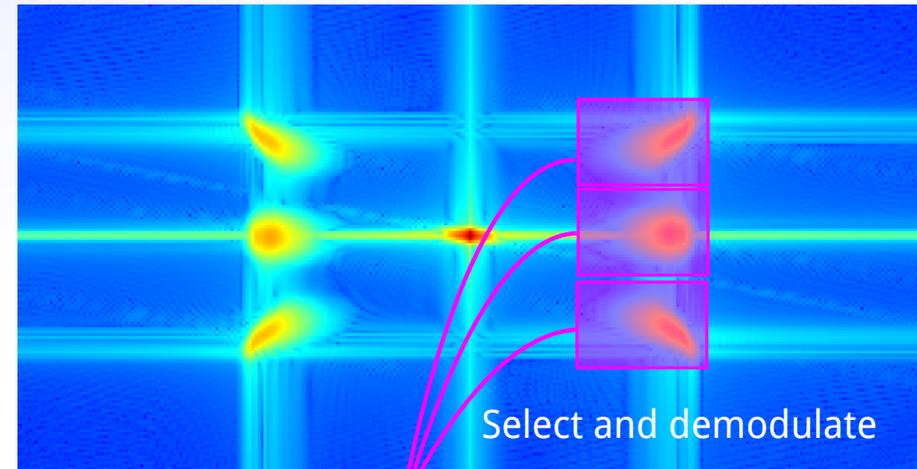


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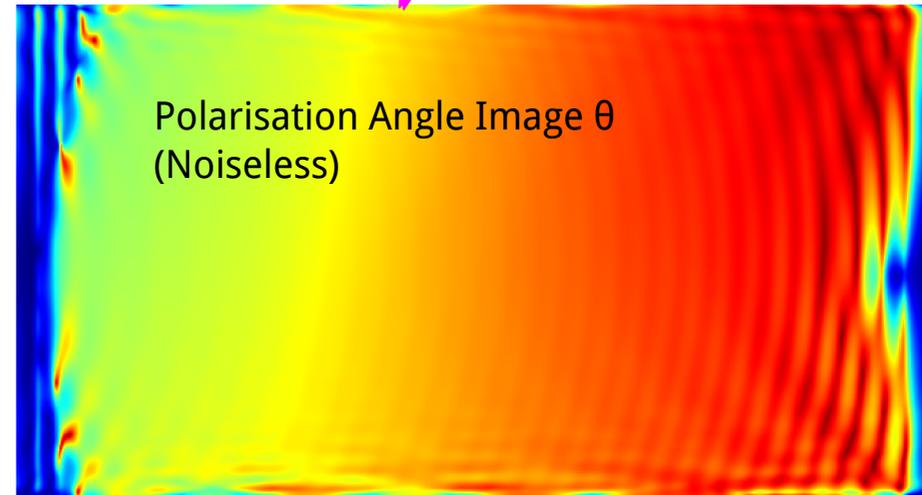
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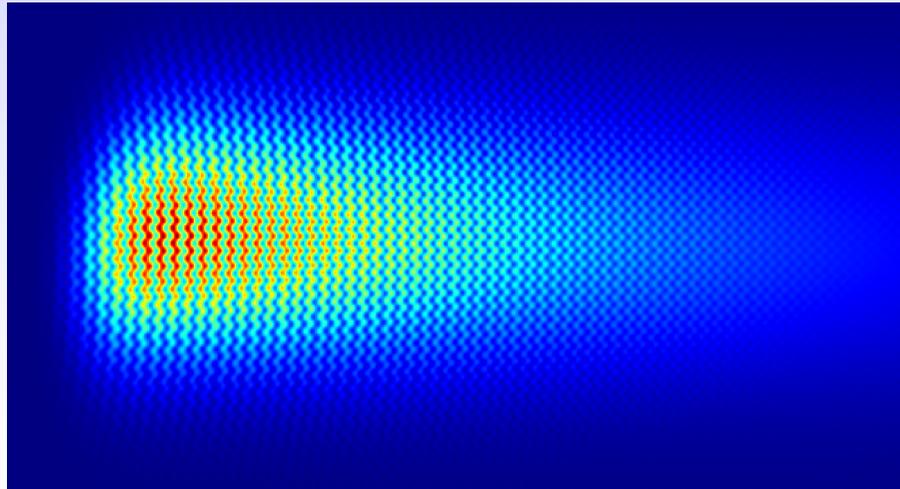
Select and demodulate



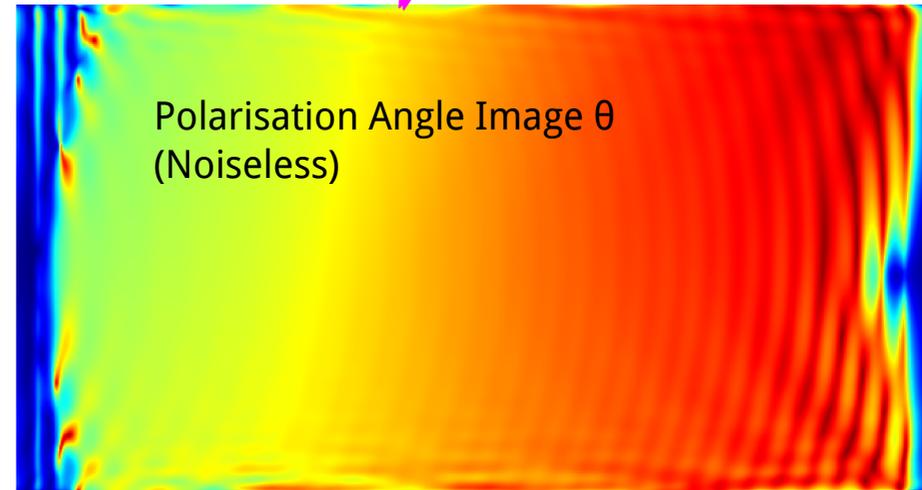
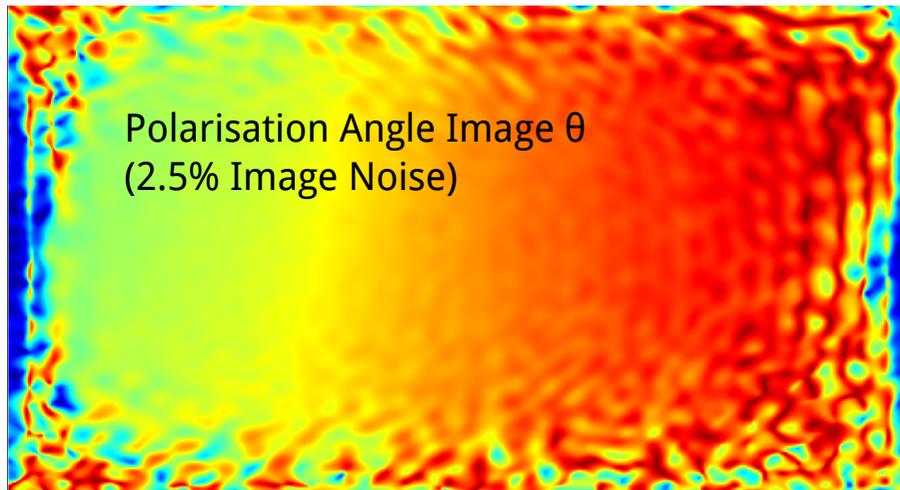
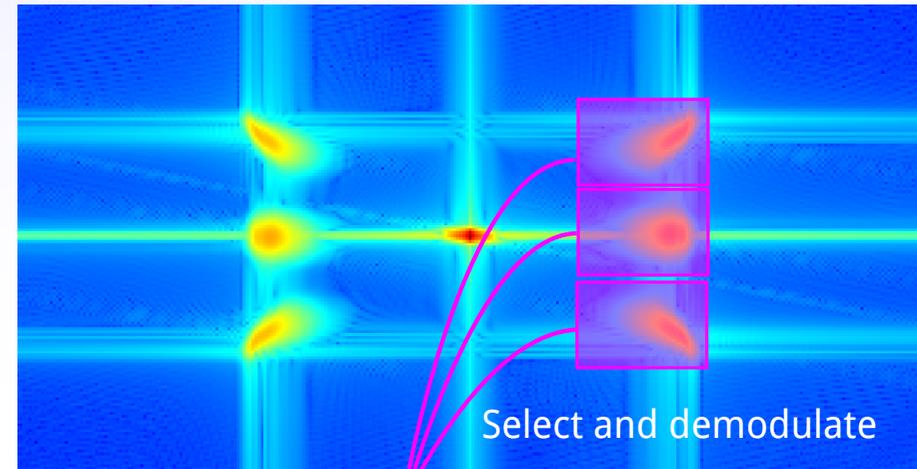
Polarisation Angle Image θ
(Noiseless)

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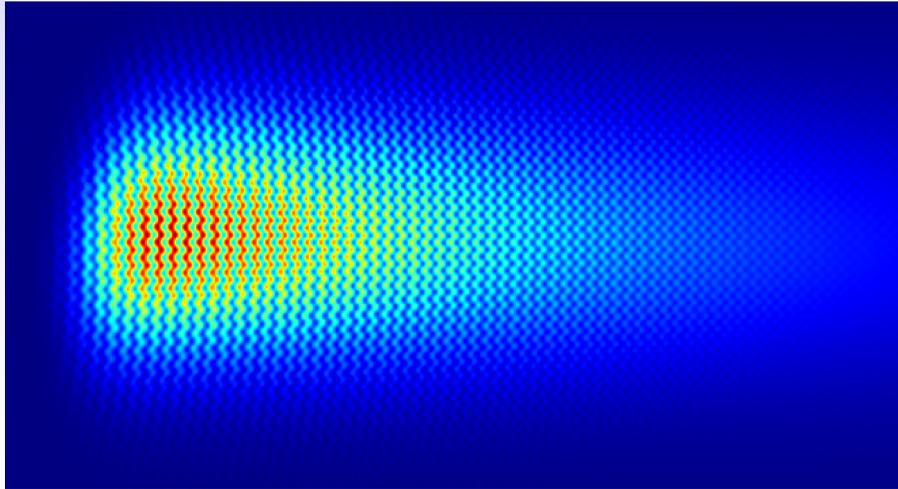
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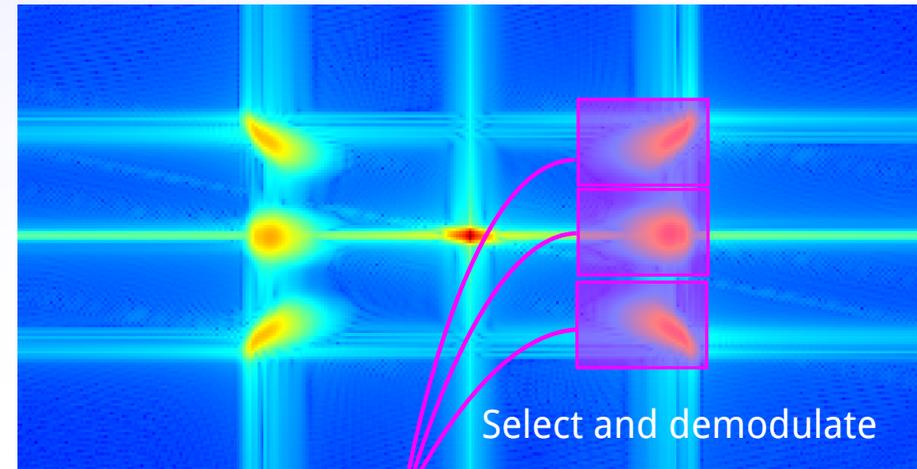
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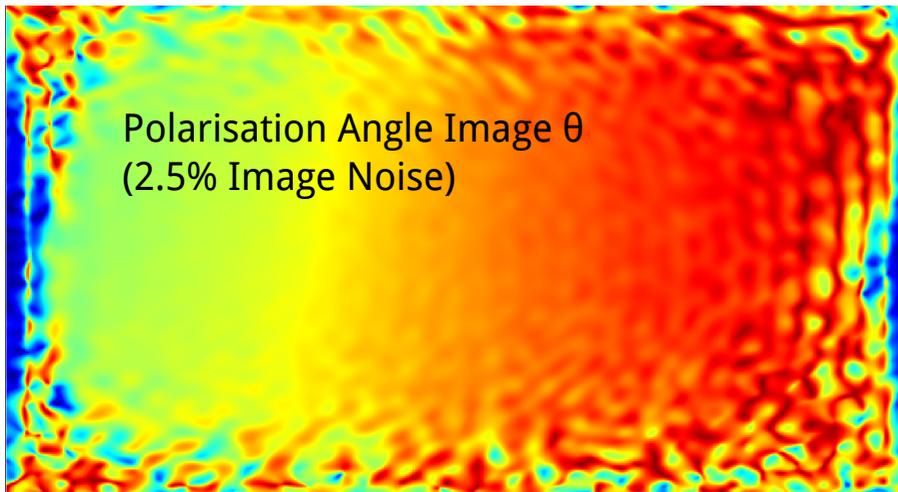
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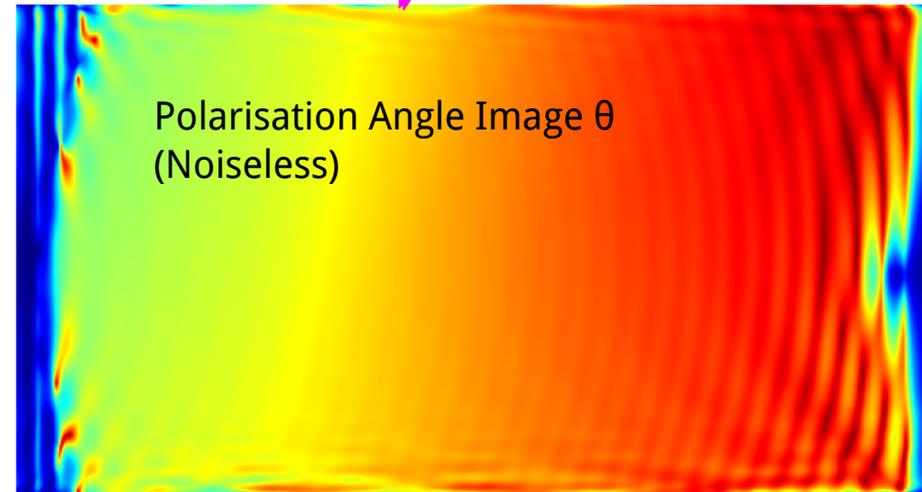
FT
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Select and demodulate



Polarisation Angle Image θ
(2.5% Image Noise)



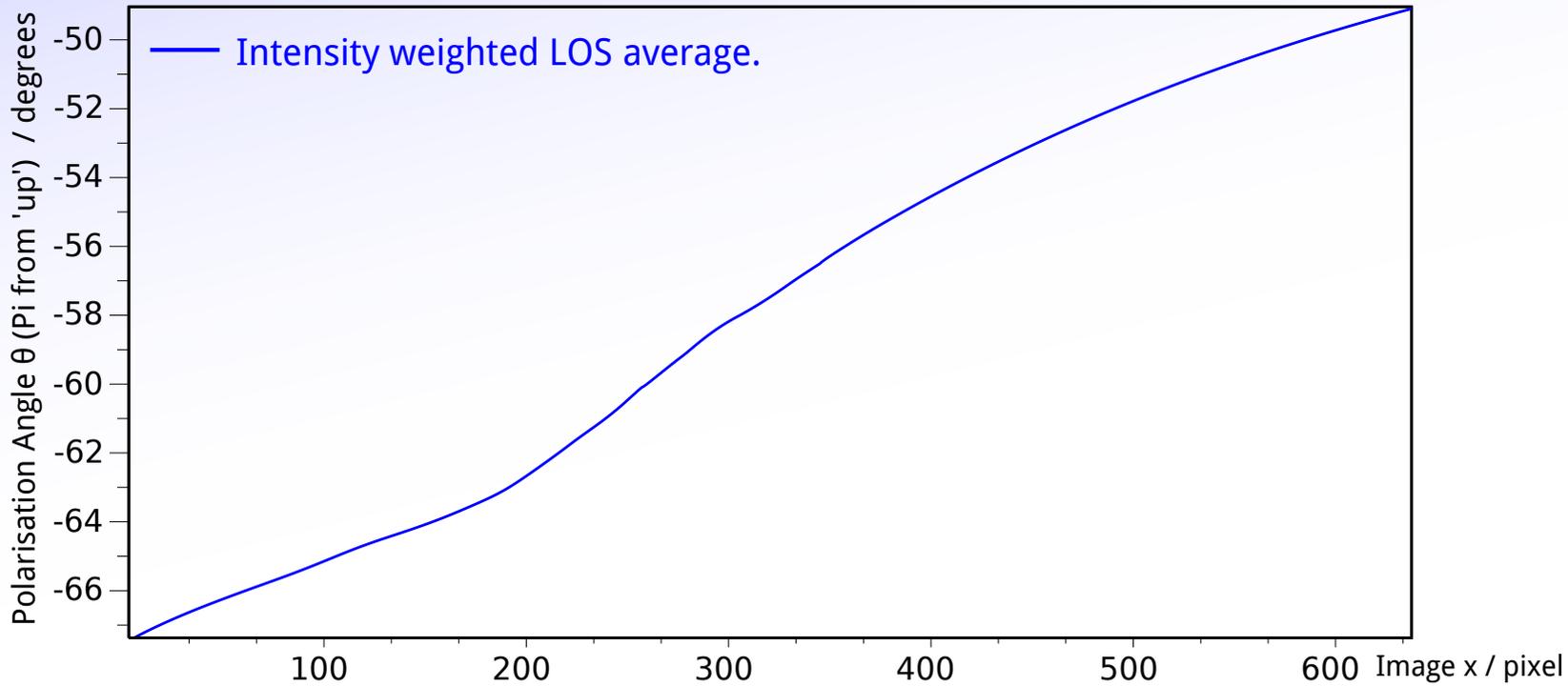
Polarisation Angle Image θ
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$$I = \int \frac{I_0}{2} [1 + \zeta (\cos 2\theta \cos(x) + \sin 2\theta \sin(x) \sin(y))] dl$$
$$\neq \frac{\langle I_0 \rangle}{2} [1 + \langle \zeta \rangle (\cos \langle 2\theta \rangle \cos(x) + \sin \langle 2\theta \rangle \sin(x) \sin(y))]$$

Actually, this is wrong. The image is really the integral of this over the LOS.
However - it seems that if we assume it is, the recovered θ is the same as the LOS average for each pixel.
The other terms are not equal to their LOS averages and introduce extra phases.

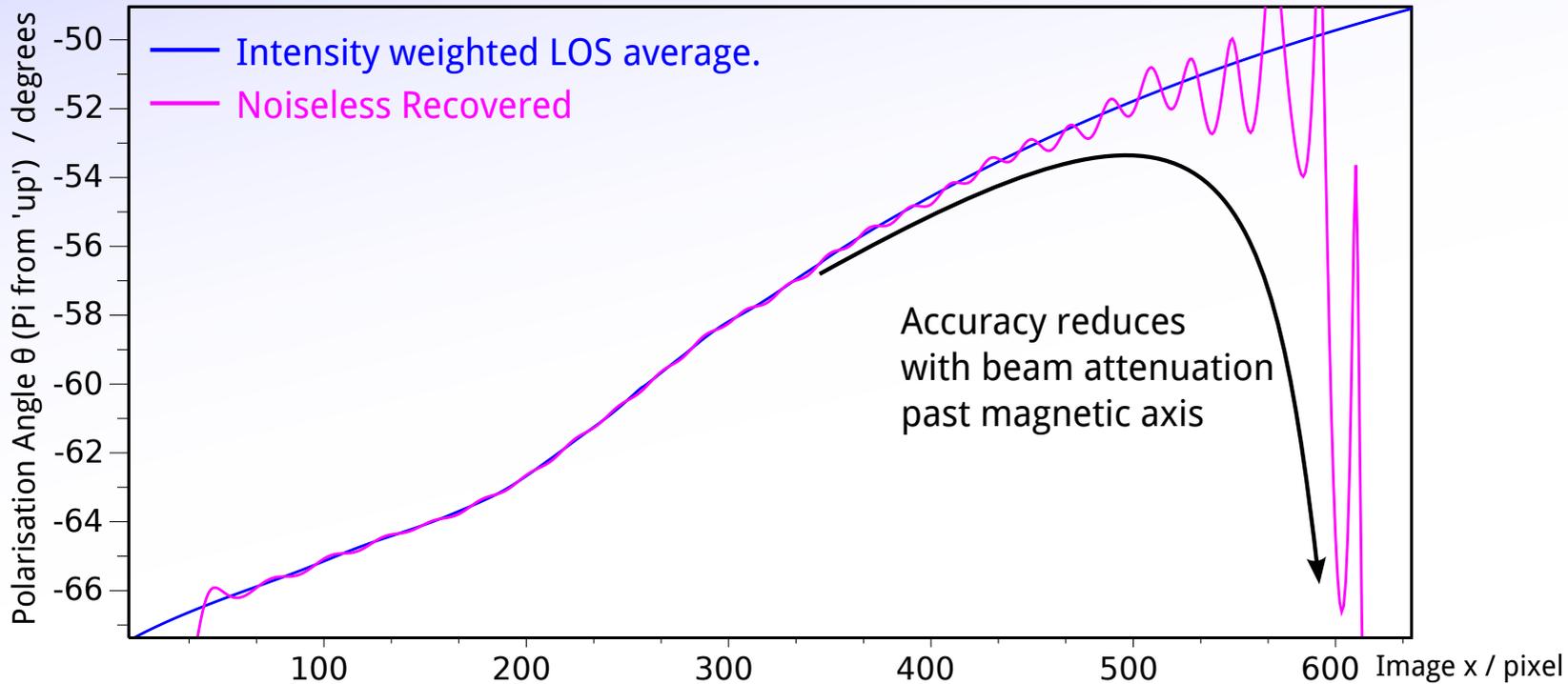
Demodulation: Accuracy of θ recovery.

Recovery of θ from a noiseless and noisy images gives us our probably accuracy:



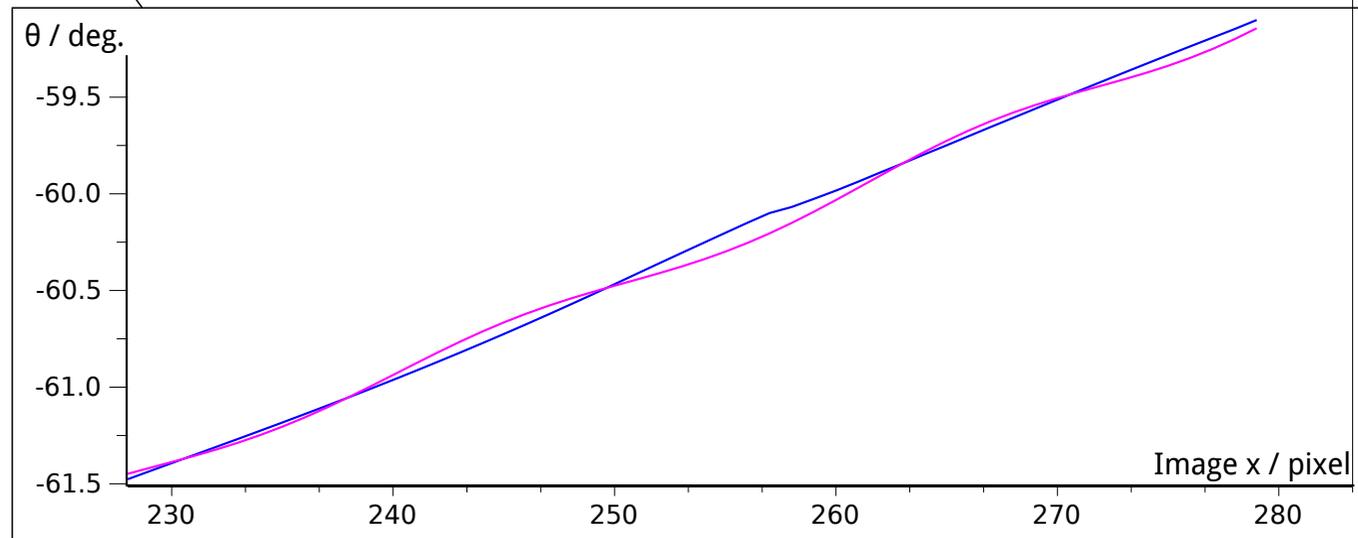
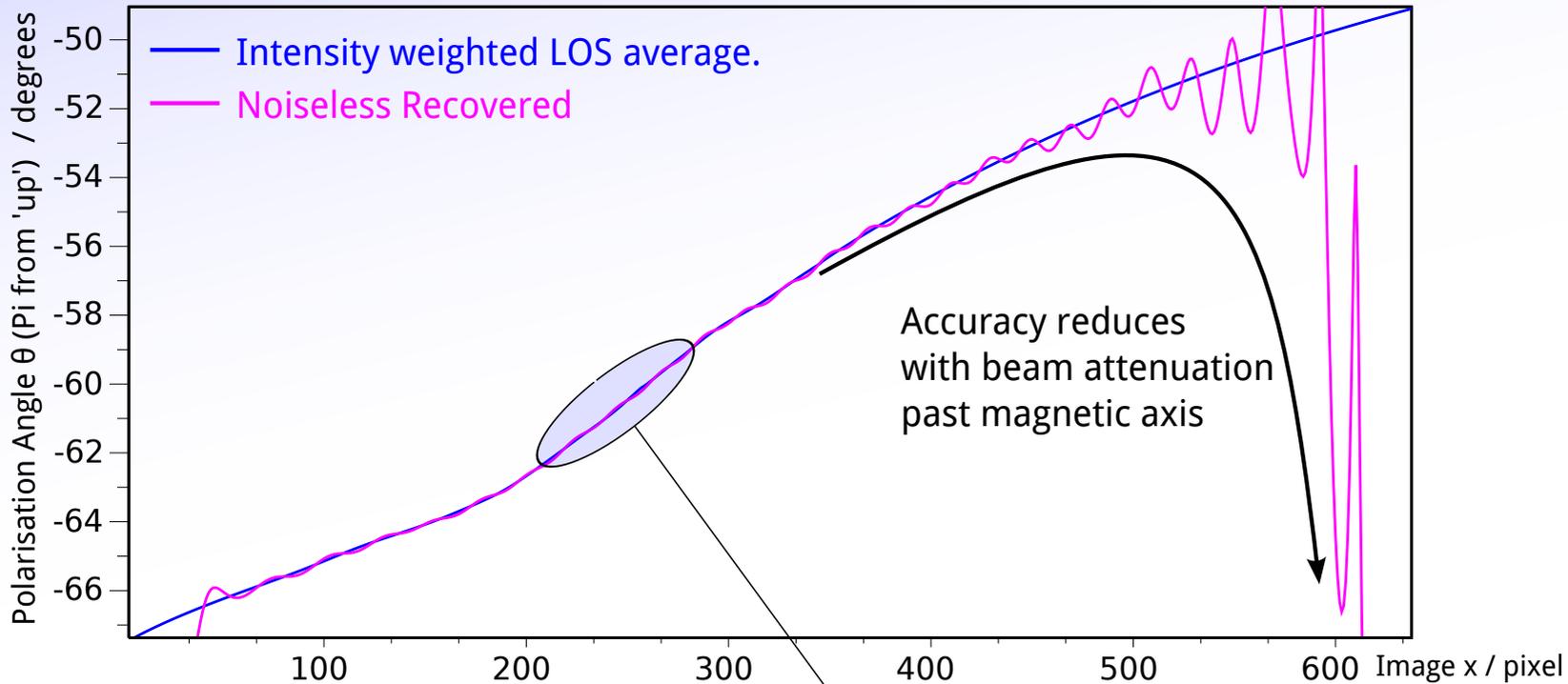
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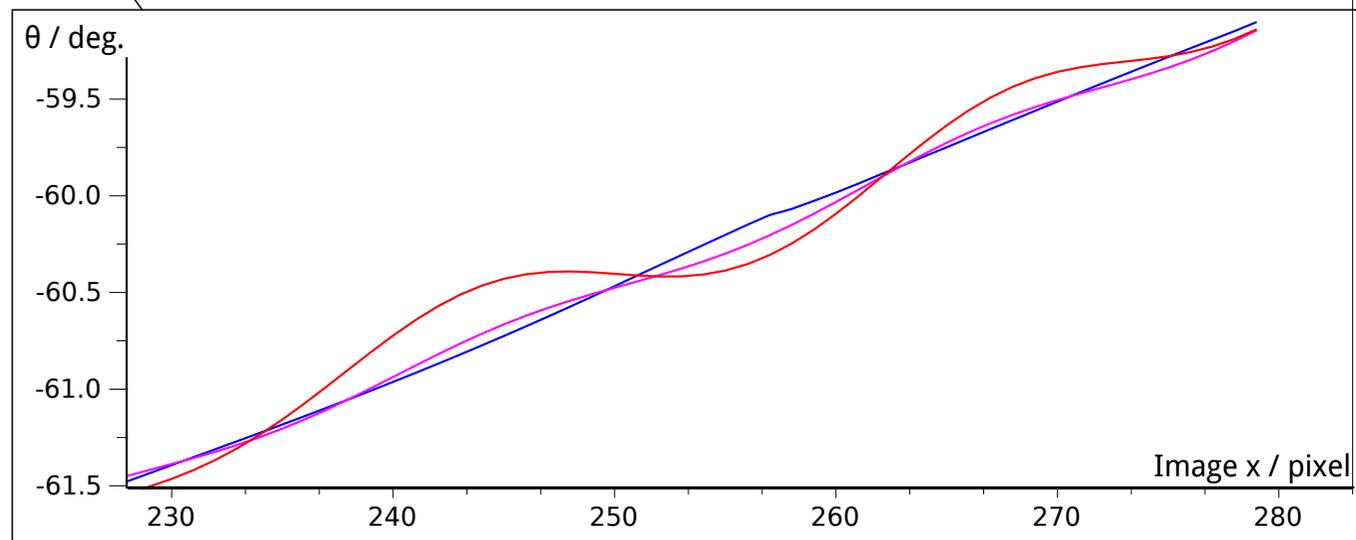
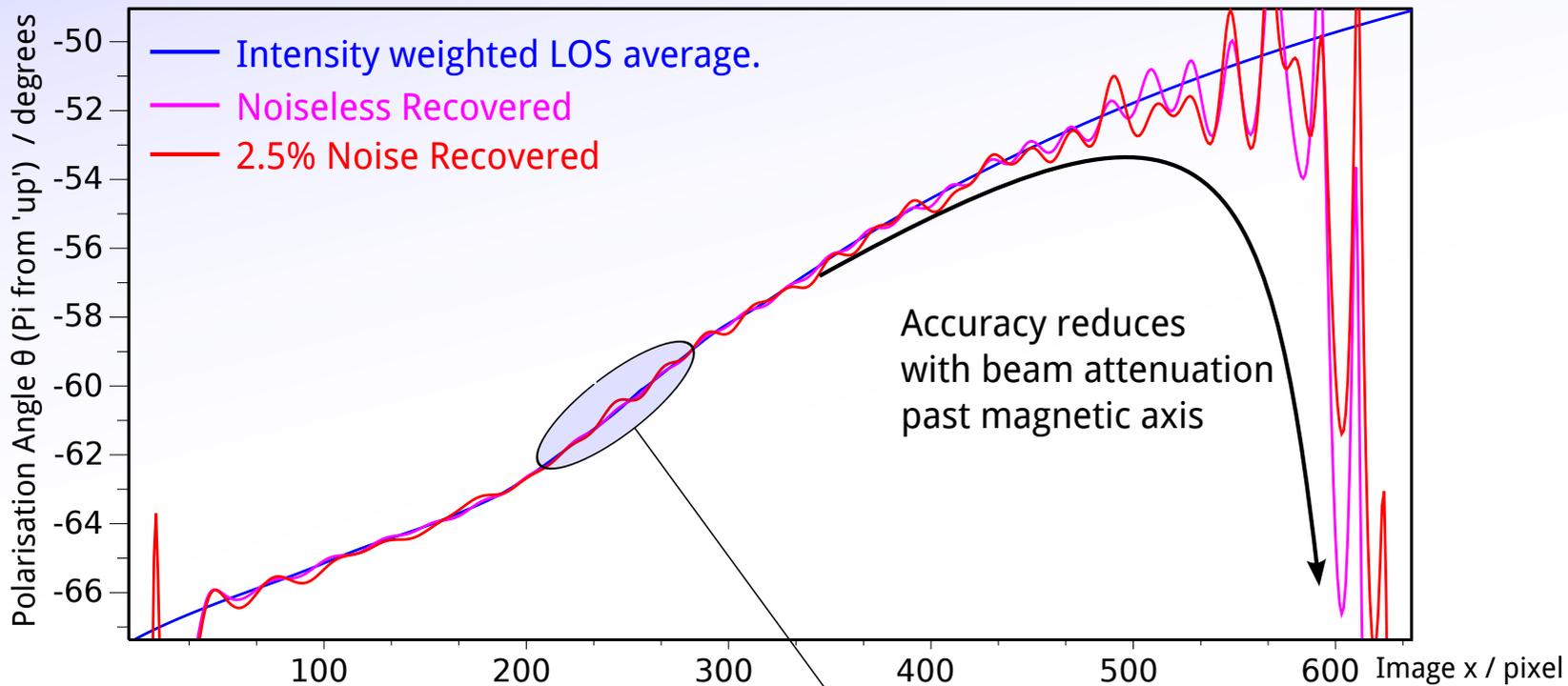
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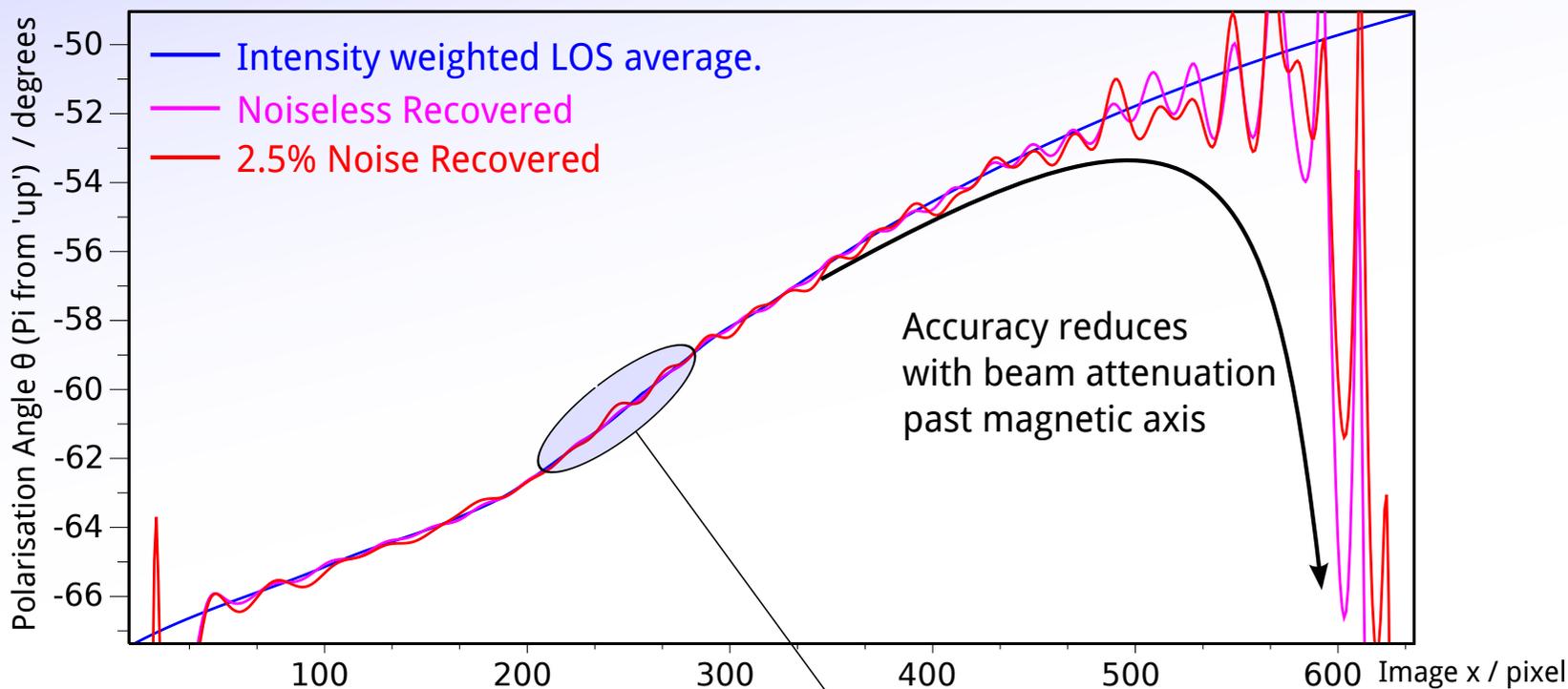
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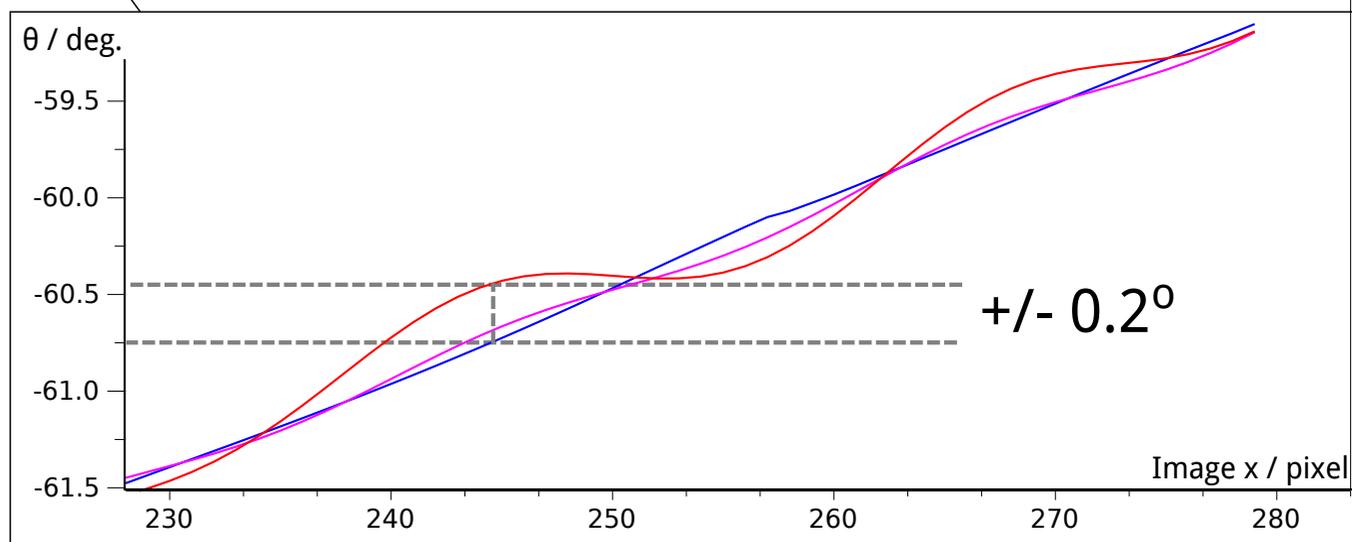


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It might be better than this, since these are partly artifacts of the demodulation. However, for now we can take $\theta \pm 0.2^\circ$ as the accuracy.





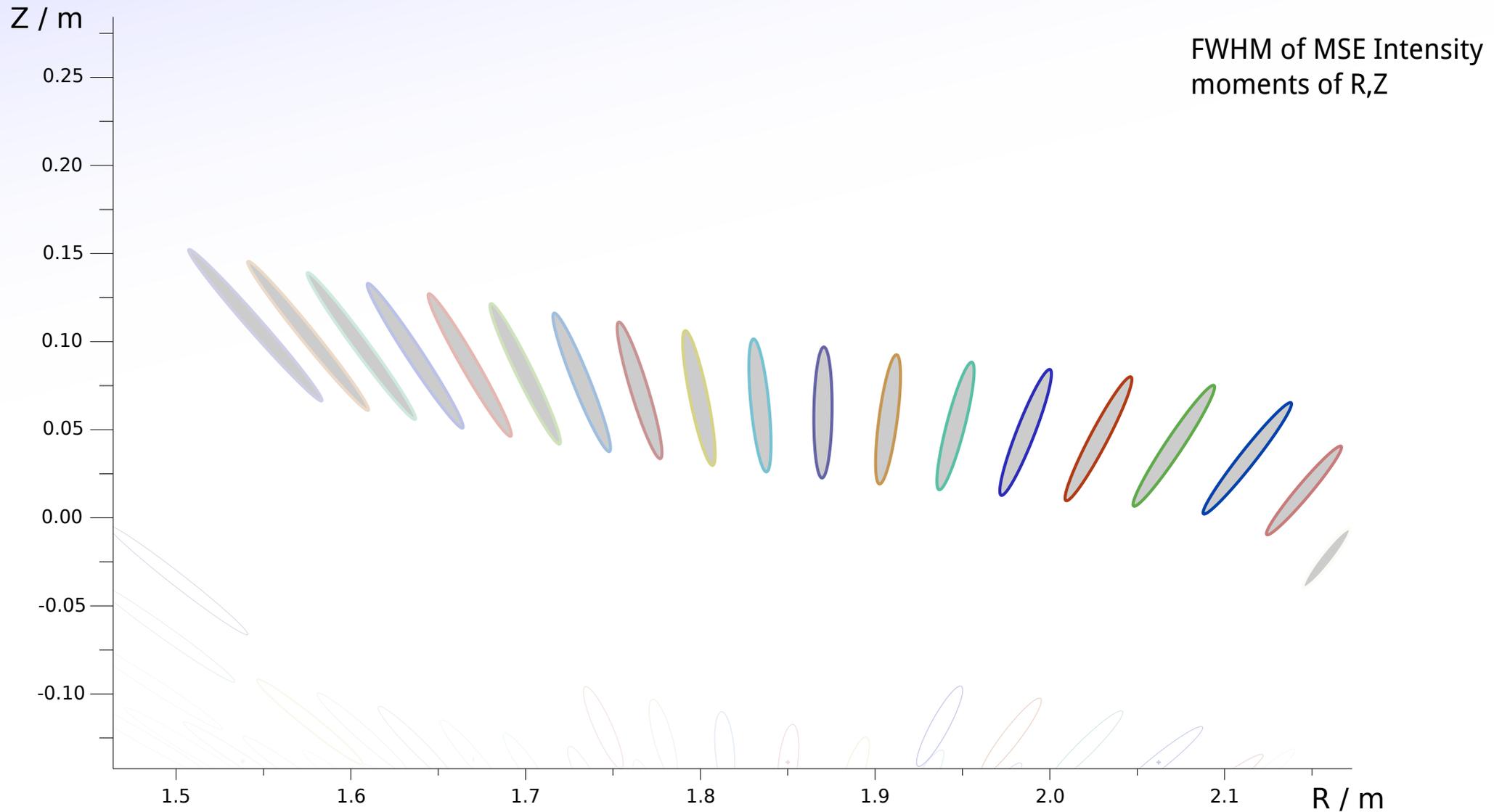
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The recovered θ are really $\langle \theta \rangle$ over the LOS. Spatial resolution is a combination of pixel-pixel averaging due to modulation (1cm) and the LOS averaging.

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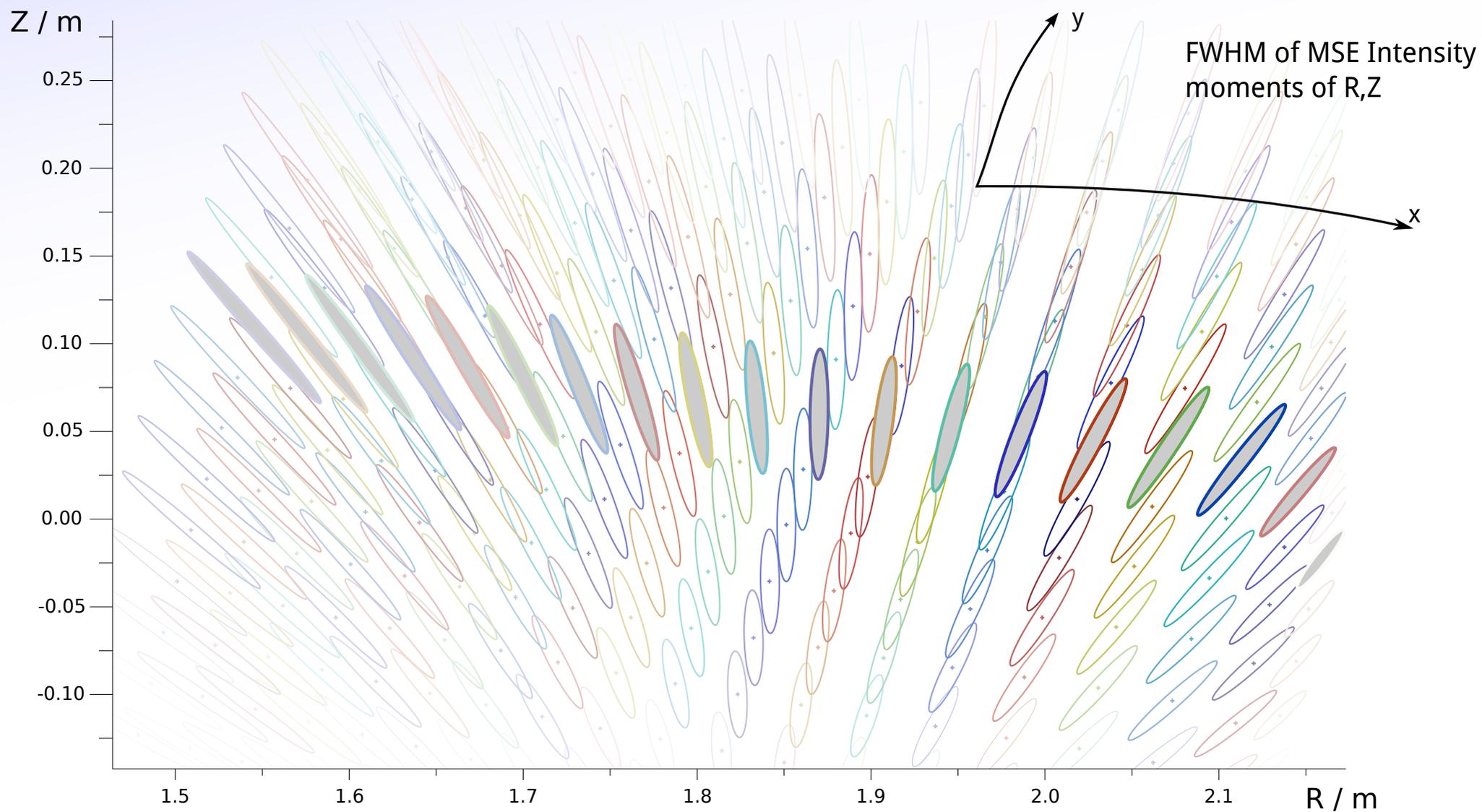
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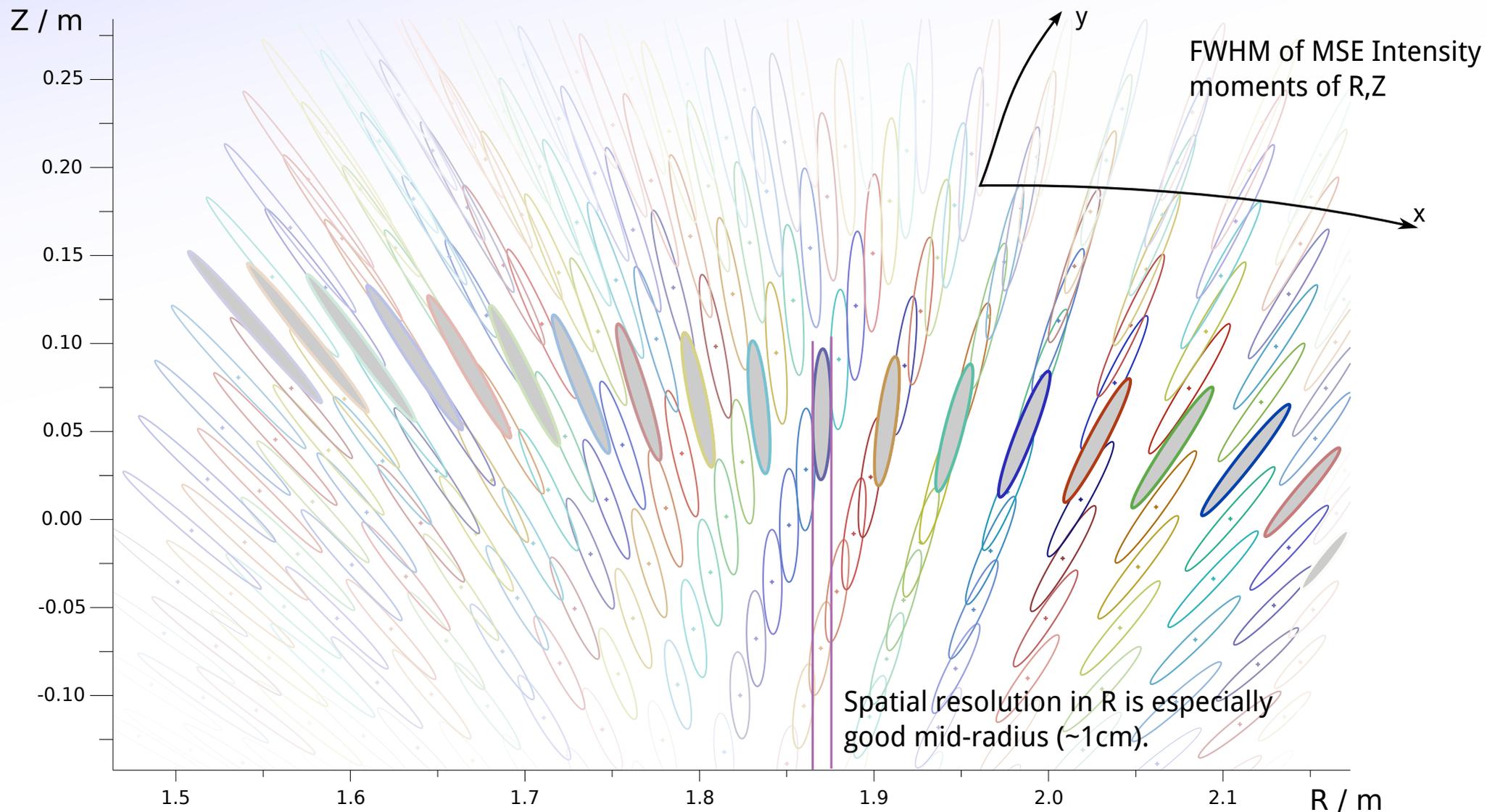
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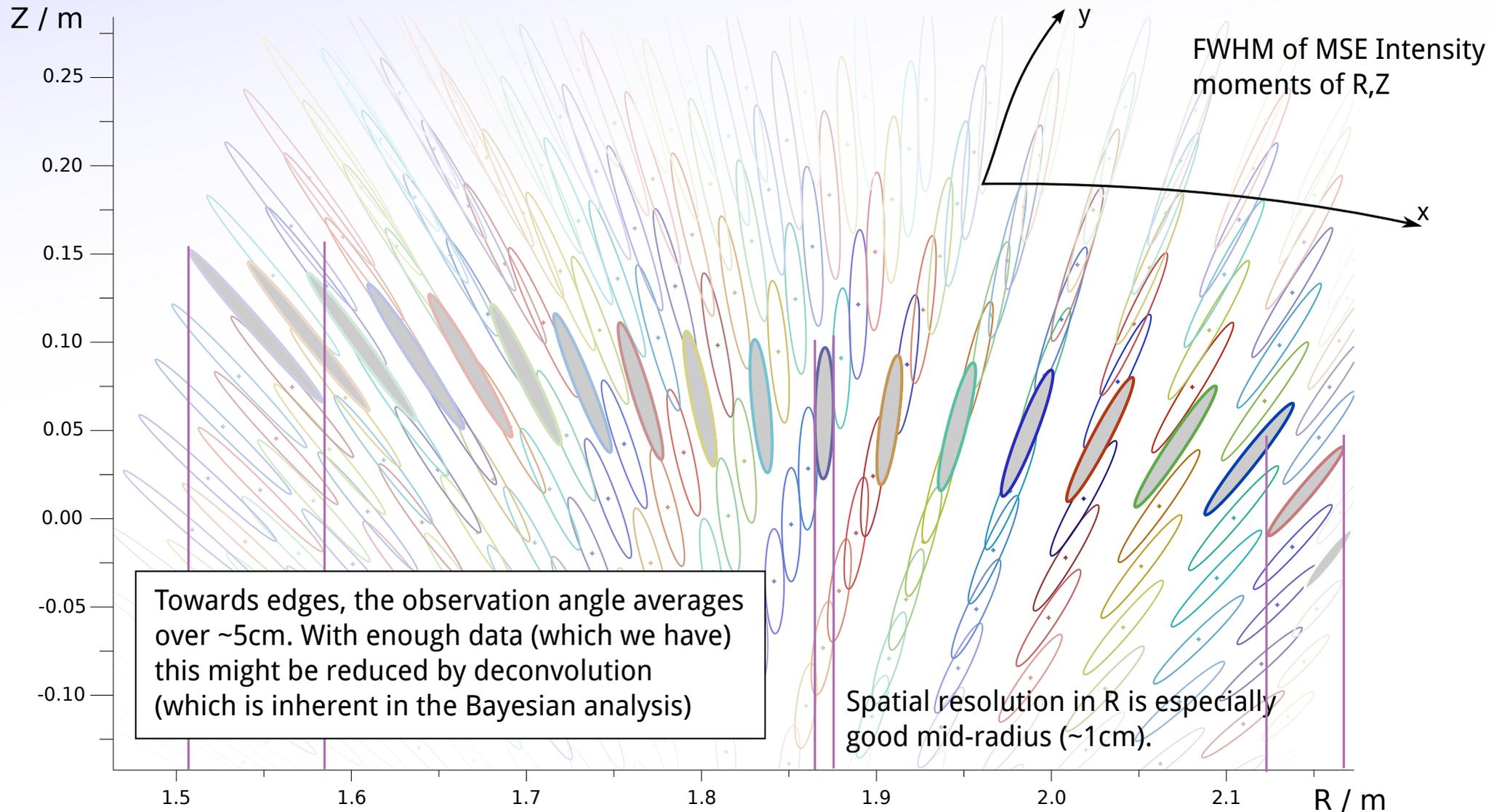
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Recovery of plasma current.

To final objective is to measure plasma current j .

For normal 1D measurements: not possible so θ used as a constraint for equilibrium.

Does having 2D measurements make it possible to calculate j without equilibrium?



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Assuming toroidal symmetry, the current is:

$$-\mu_0 j_\phi = \frac{\partial B_z}{\partial R} + \frac{1}{R} \frac{\partial^2}{\partial Z^2} \int_0^R R' B_z(R', Z) dR'$$

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The new term gives localisation of current in Z (\sim via curvature of field).

Recovery of plasma current.

To final objective is to measure plasma current j .

For normal 1D measurements: not possible so θ used as a constraint for equilibrium.

Does having 2D measurements make it possible to calculate j without equilibrium?

Assuming toroidal symmetry, the current is:

$$-\mu_0 j_\phi = \frac{\partial B_z}{\partial R} + \frac{1}{R} \frac{\partial^2}{\partial Z^2} \int_0^R R' B_z(R', Z) dR'$$

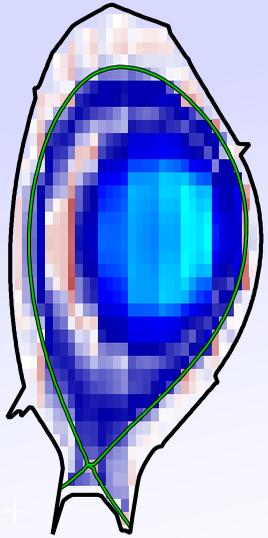
Assume we know B_ϕ as the vacuum field, then we can calculate B_z from θ .

However, we only see where the MSE emission is, so can only integrate from some $R = R_0$:

$$-\mu_0 j_\phi = \underbrace{\frac{\partial B_z}{\partial R}}_{\text{This we have with 1D MSE.}} + \underbrace{\frac{1}{R} \frac{\partial^2 \psi(R_0, Z)}{\partial Z^2}}_{\text{Function of Z that we cannot know.}} + \underbrace{\frac{1}{R} \frac{\partial^2}{\partial Z^2} \int_{R_0}^R R' B_z(R', Z) dR'}_{\text{The new term gives localisation of current in Z (~via curvature of field).}}$$

A normal MSE system has only $B_z(R)$ so cannot calculate the 3rd term.

In theory, with 2D measurements, we can.



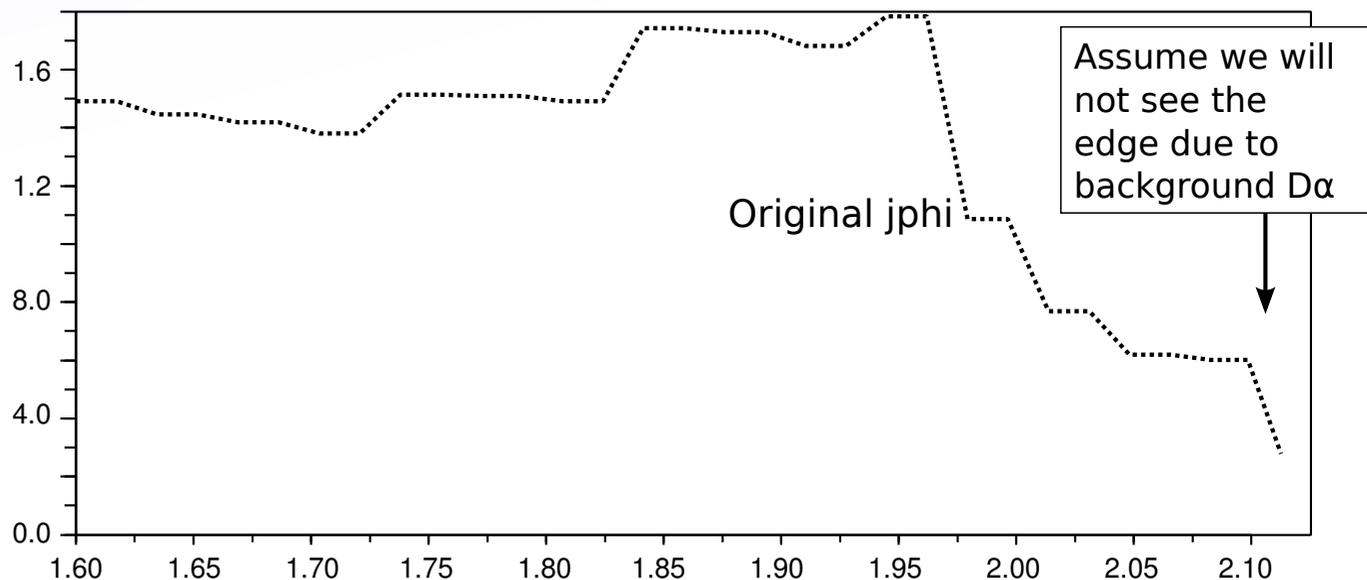
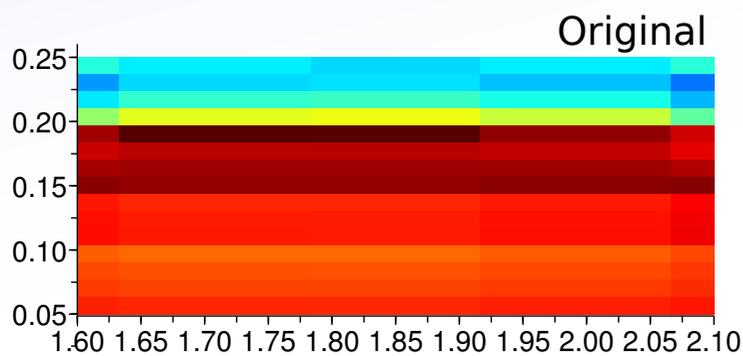
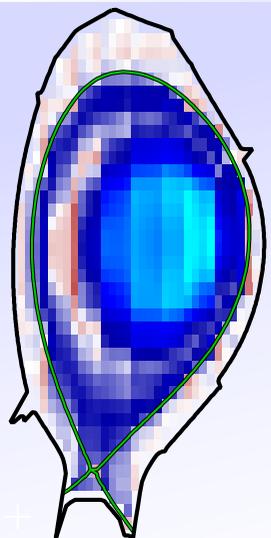
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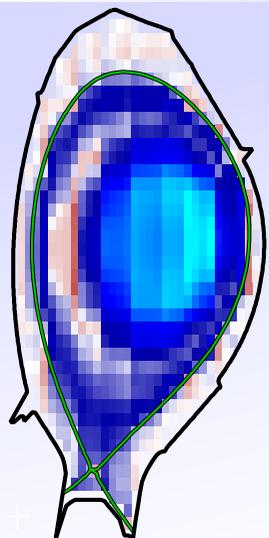
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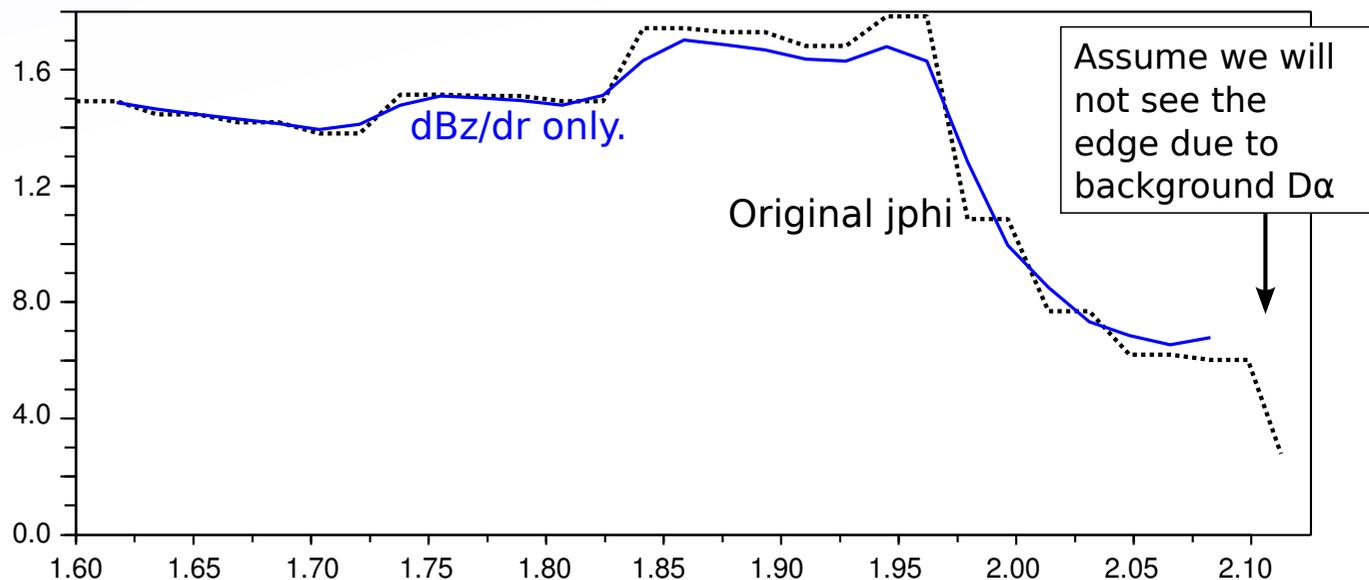
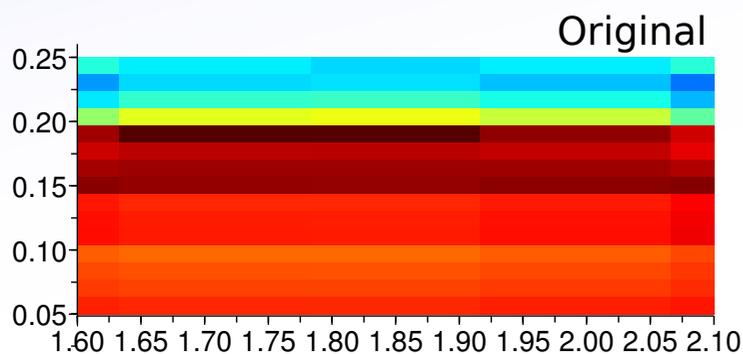


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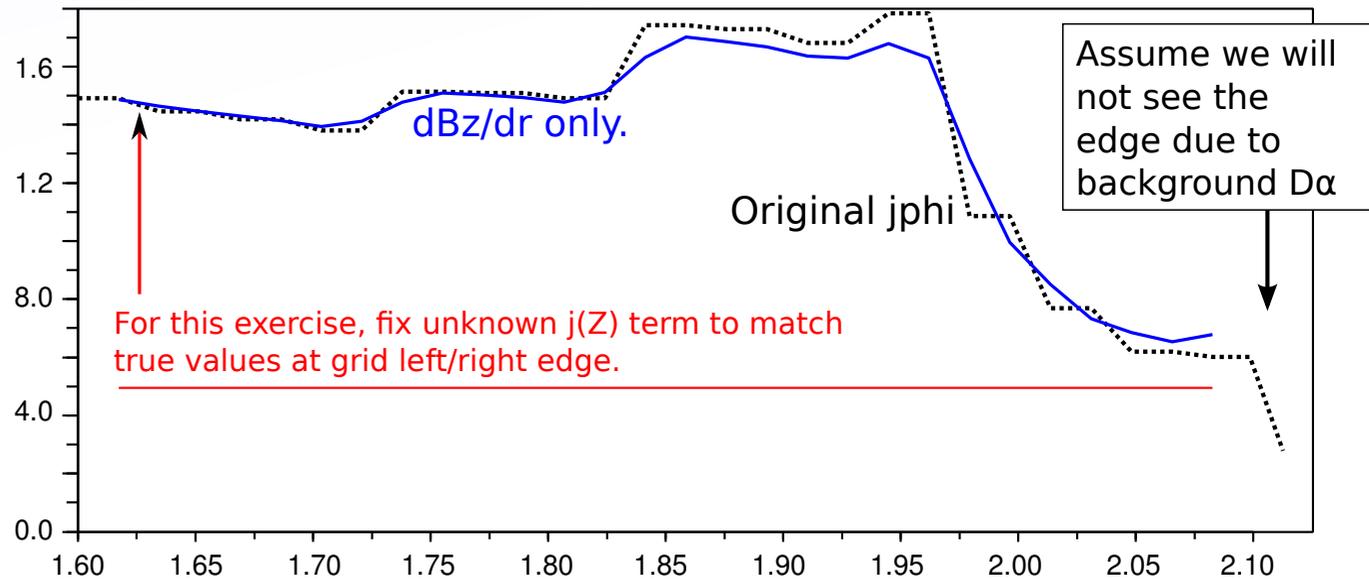
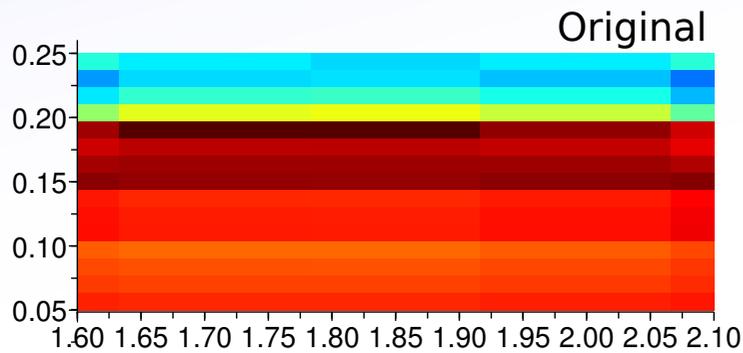
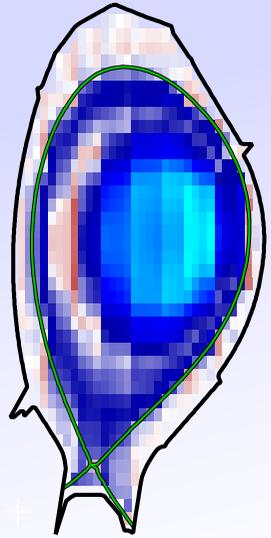
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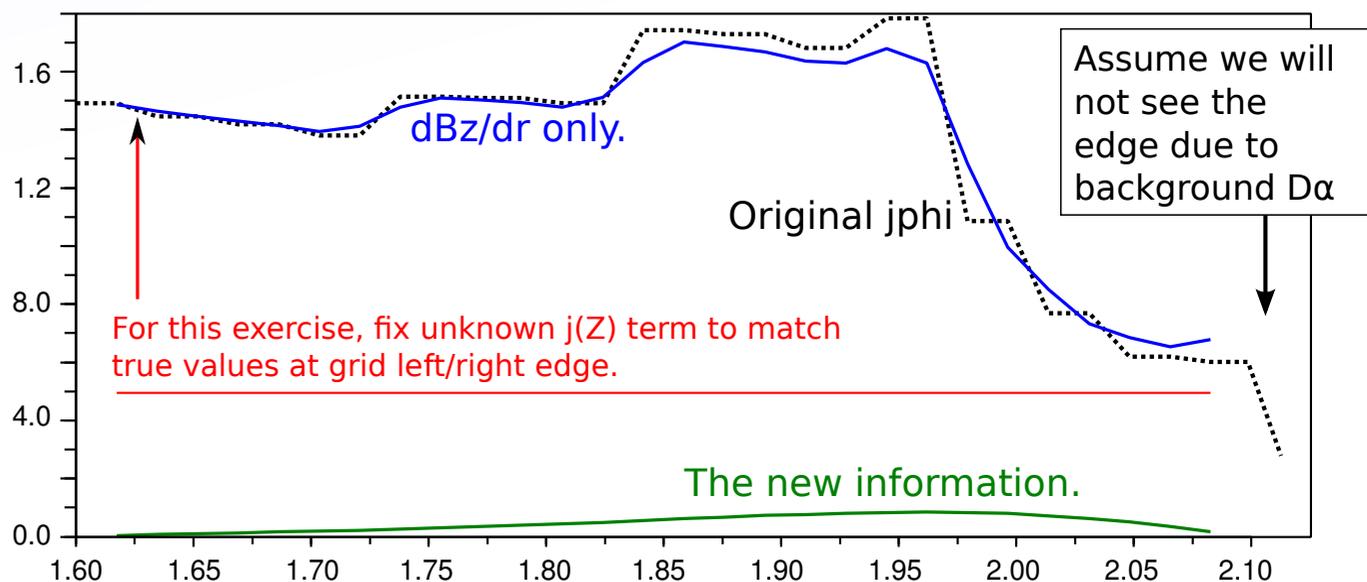
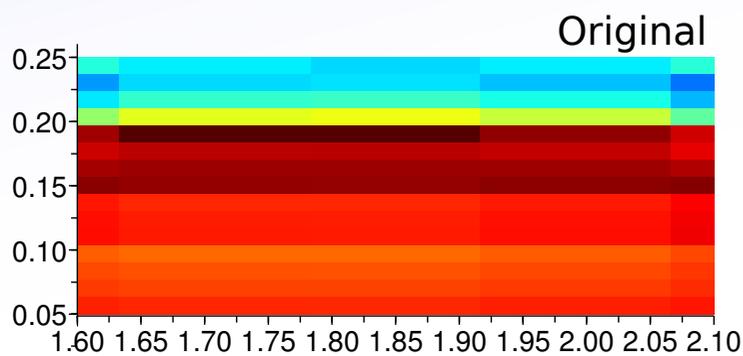
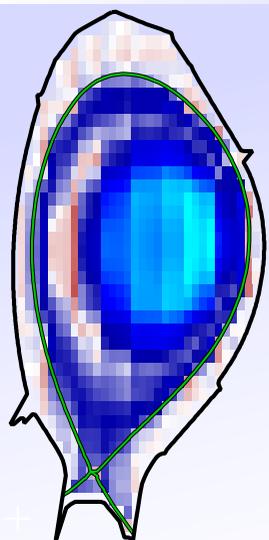
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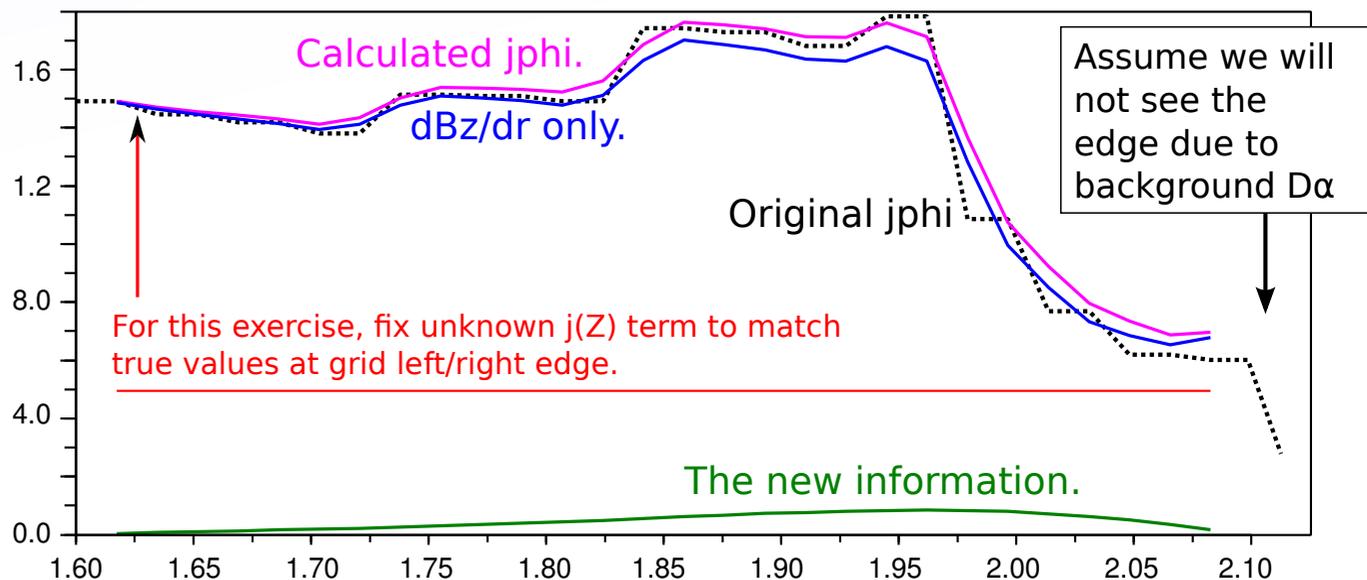
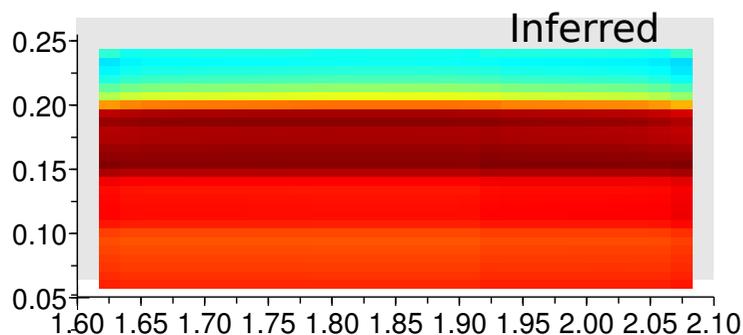
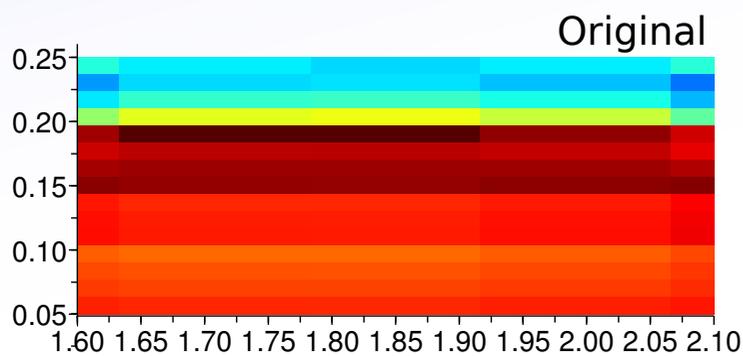
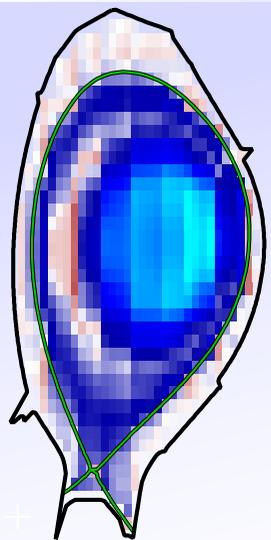
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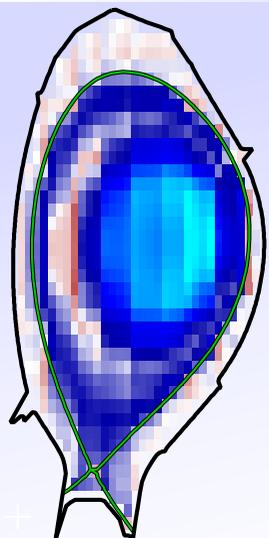
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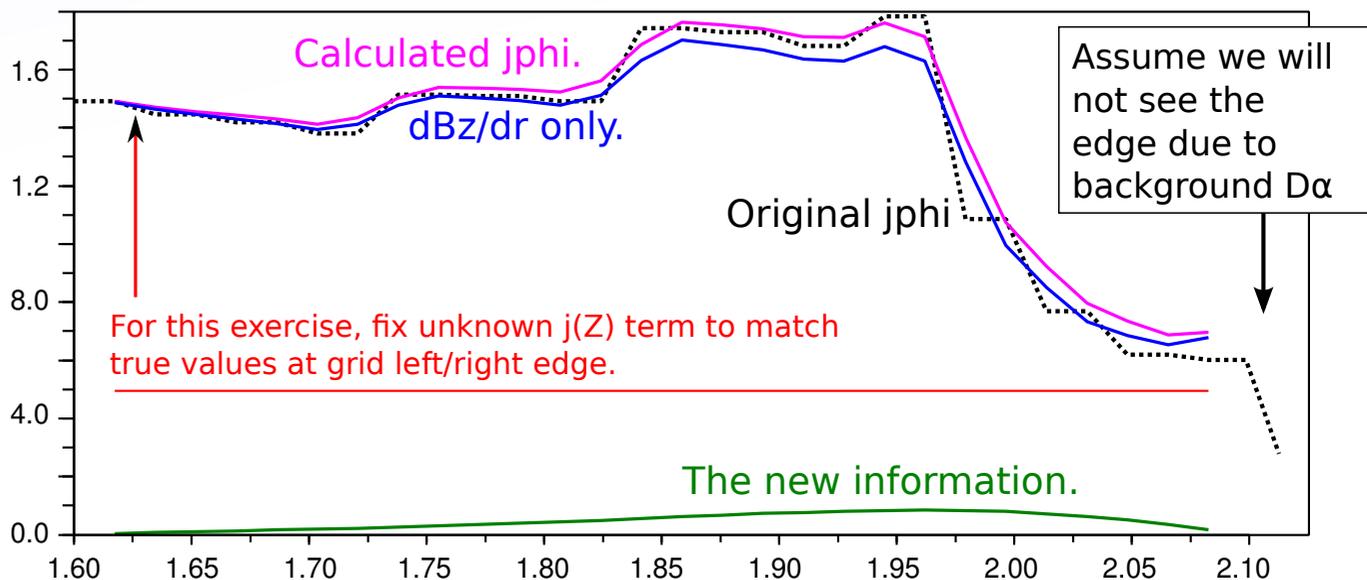
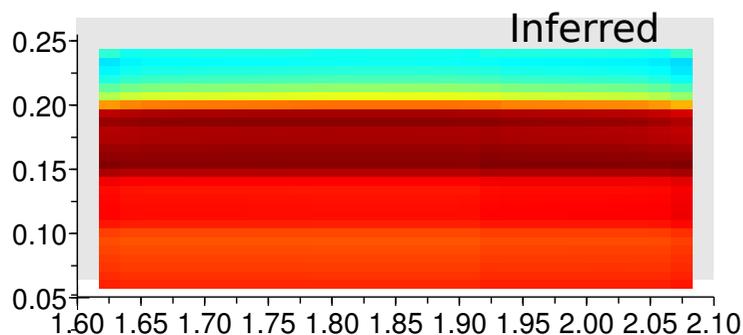
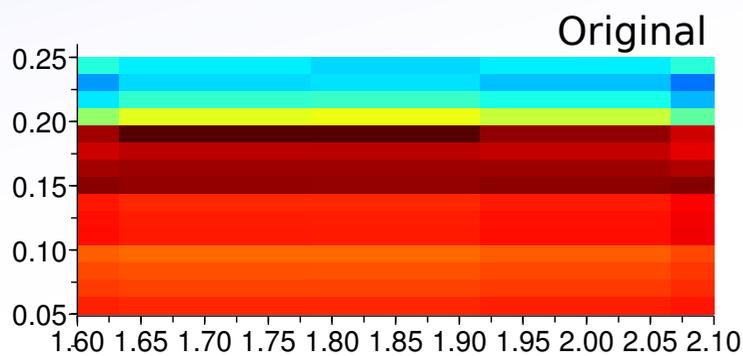


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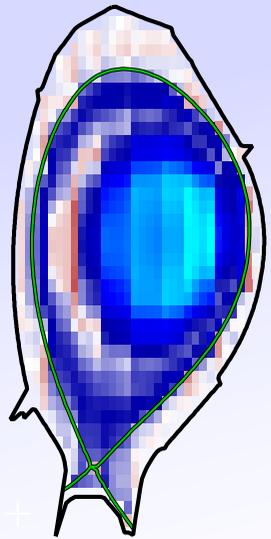
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The new part is < 10%. We do gain it mathematically but it will be entirely lost in the noise. Anyway, the $f(Z)$ term is still not known.

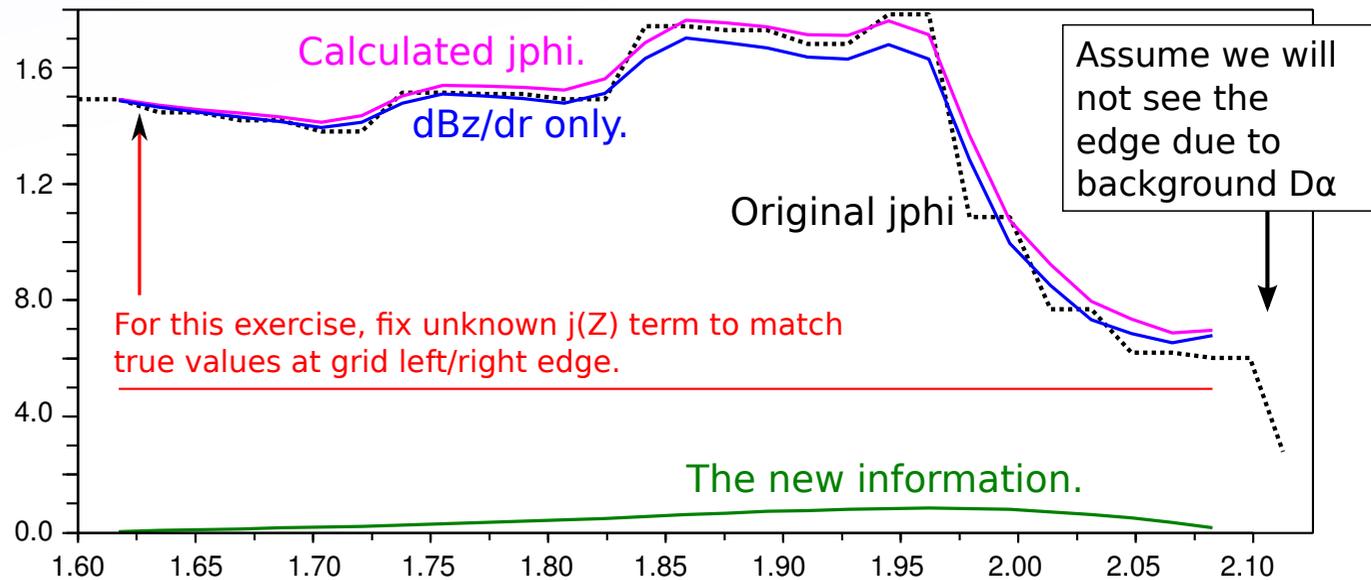
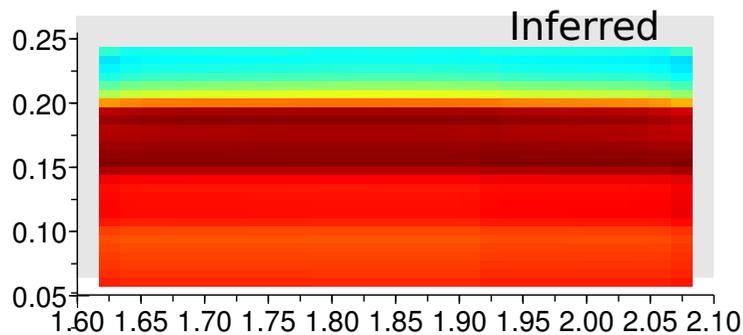
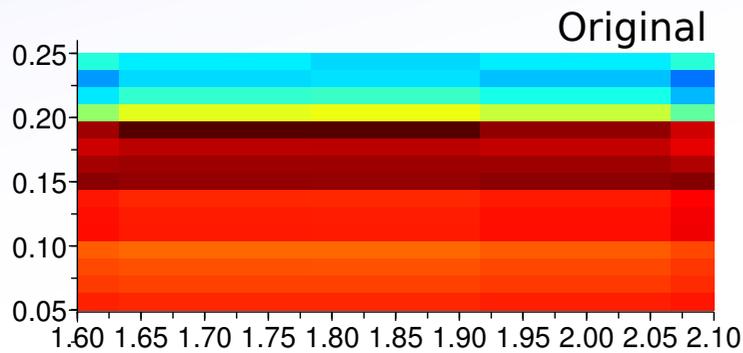
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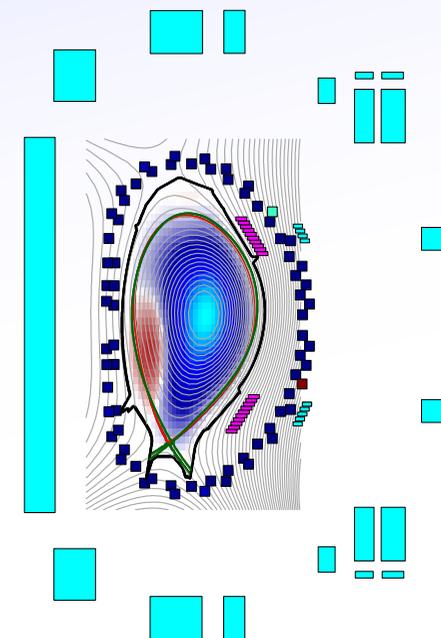
Conclusion: **No**. You still cannot exactly calculate $j\phi$ directly.

However, we still might not need to go as far as equilibrium as we also gain measurements of $\frac{dB_z}{dR}$ at different Z s. Together with normal coil measurements, it is now part of a complex tomography problem that we have done before.



By current tomography...

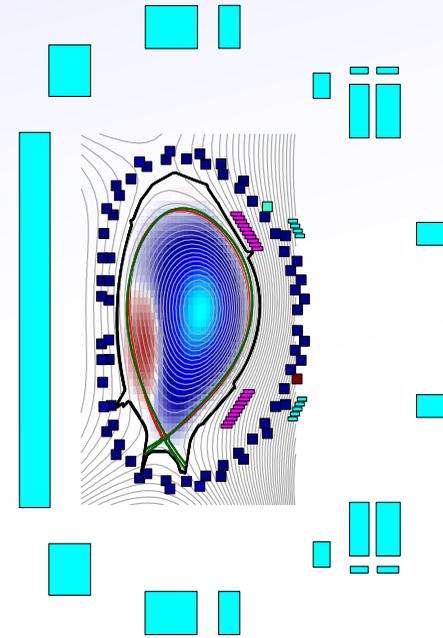
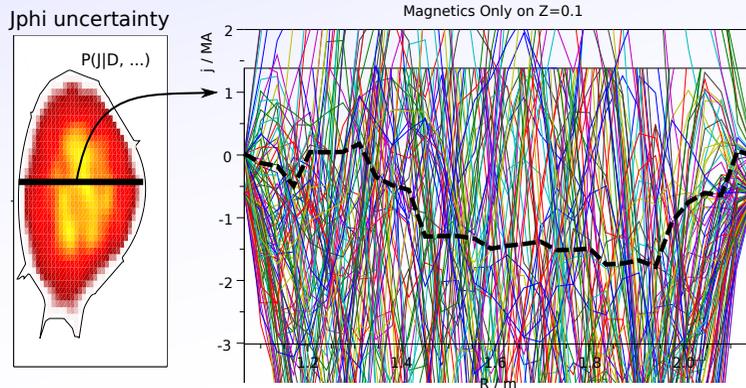
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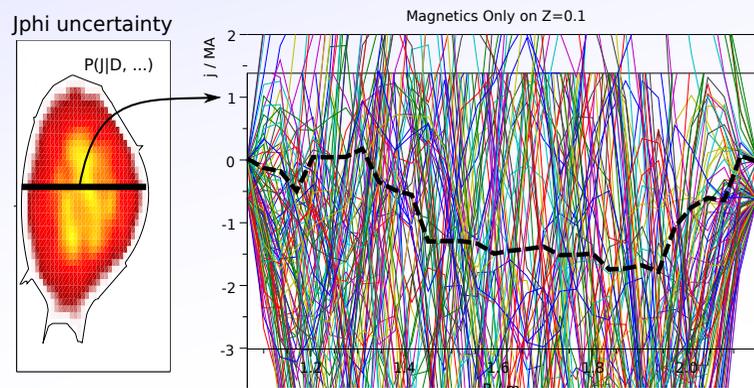
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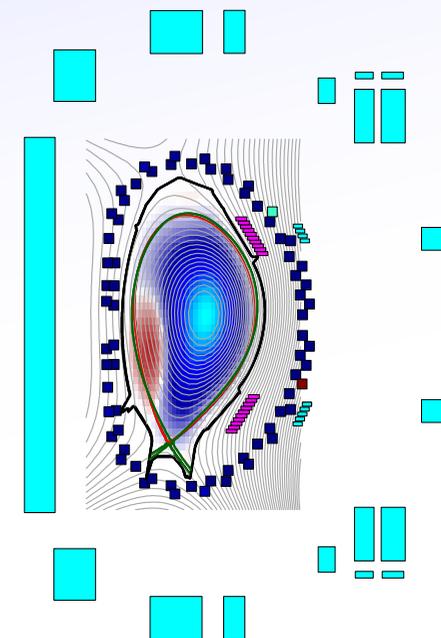
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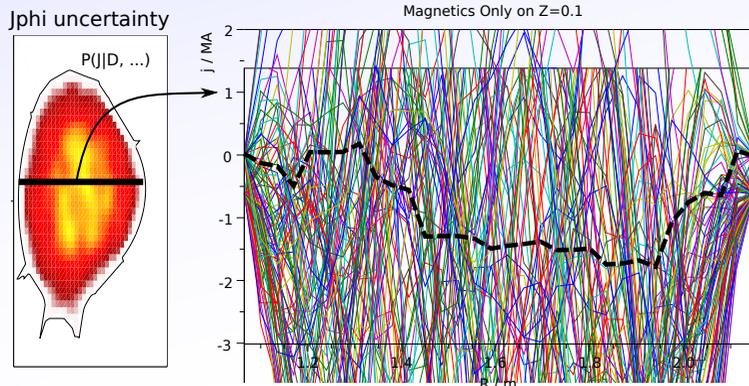
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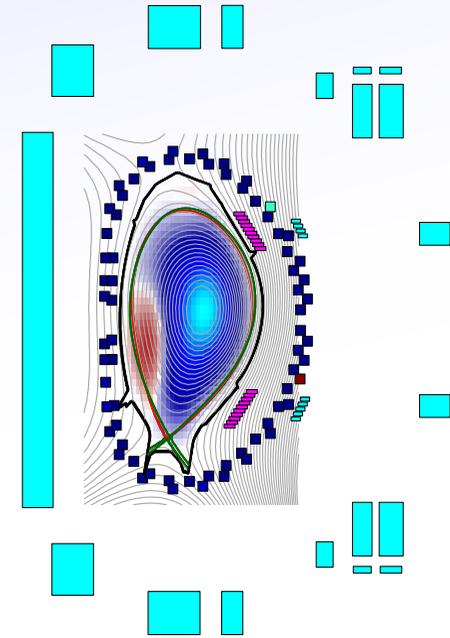
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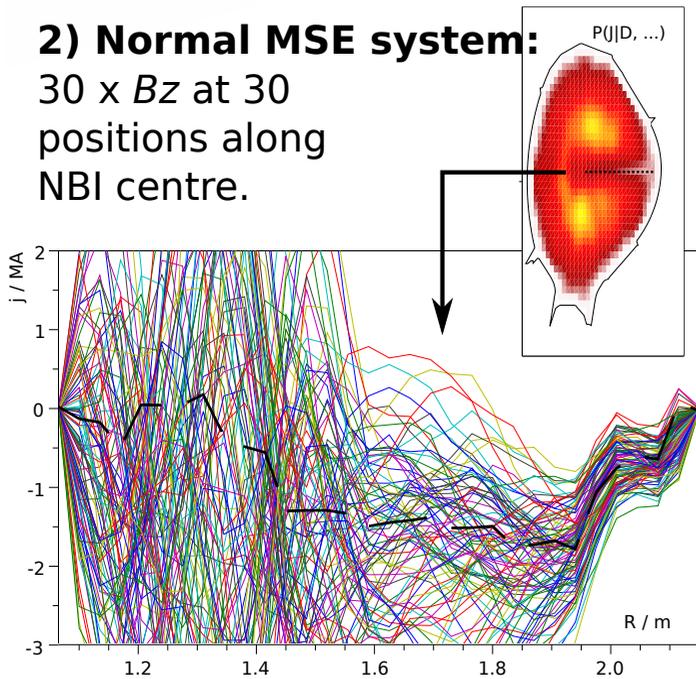


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2) Normal MSE system:

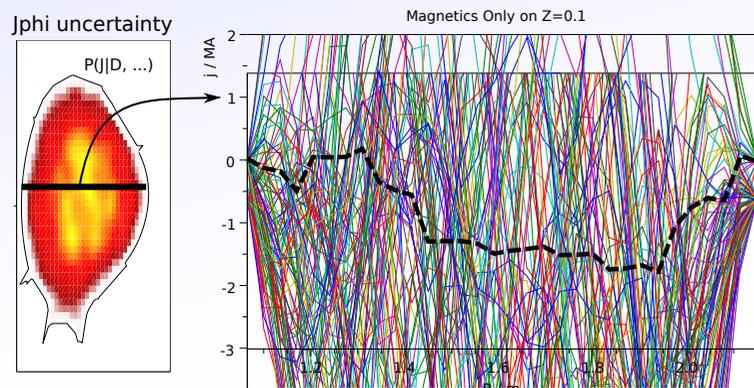
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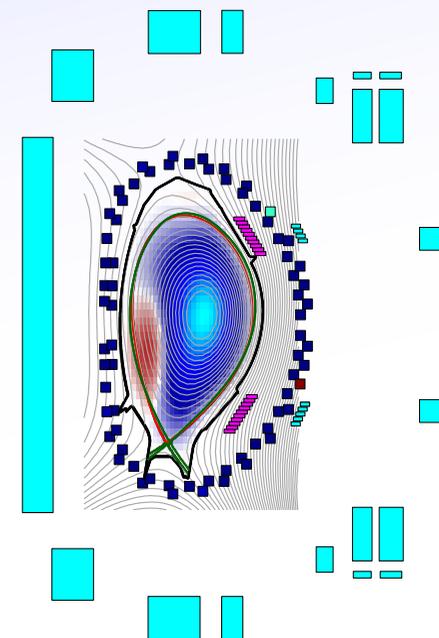
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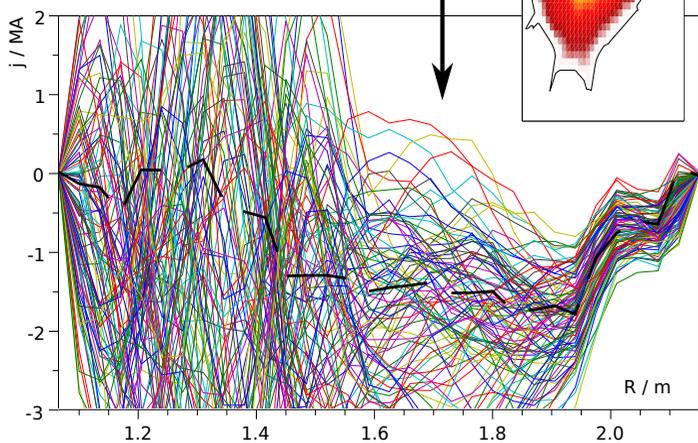
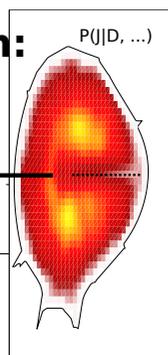


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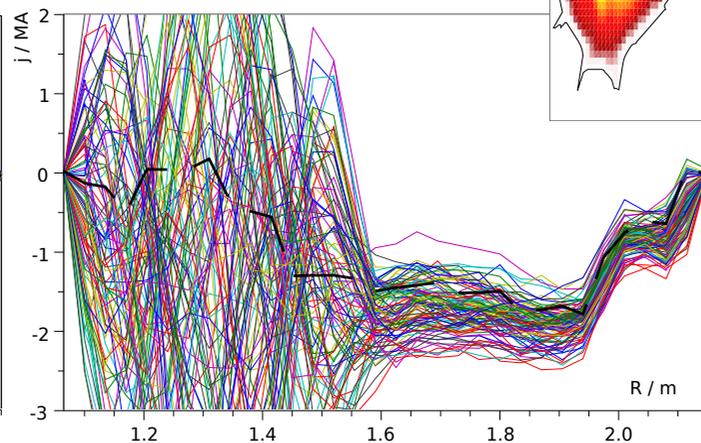
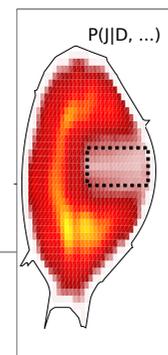
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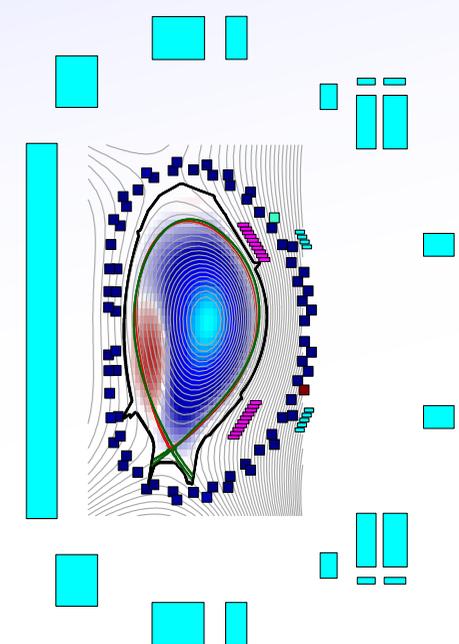
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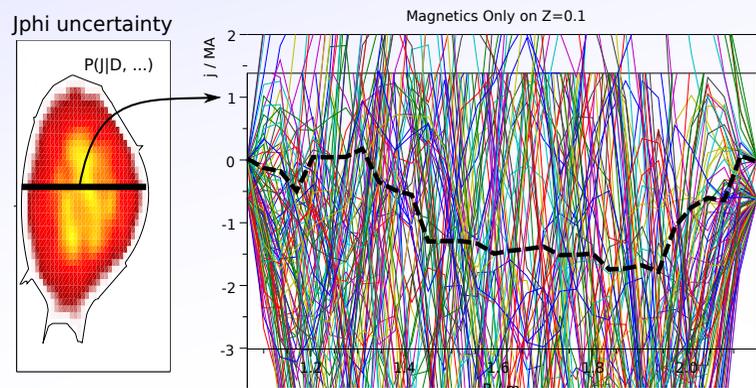


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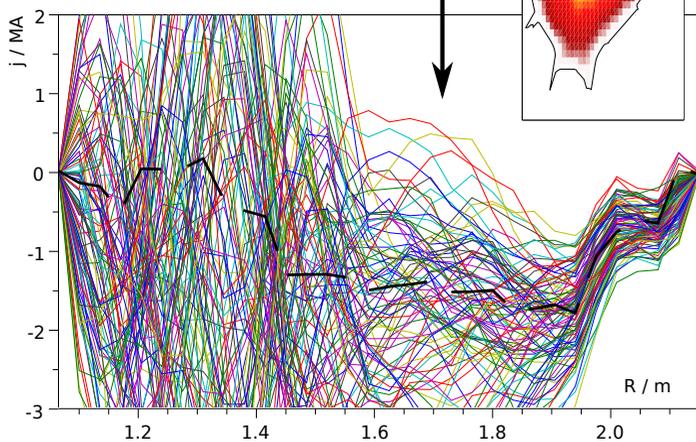
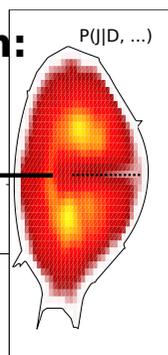
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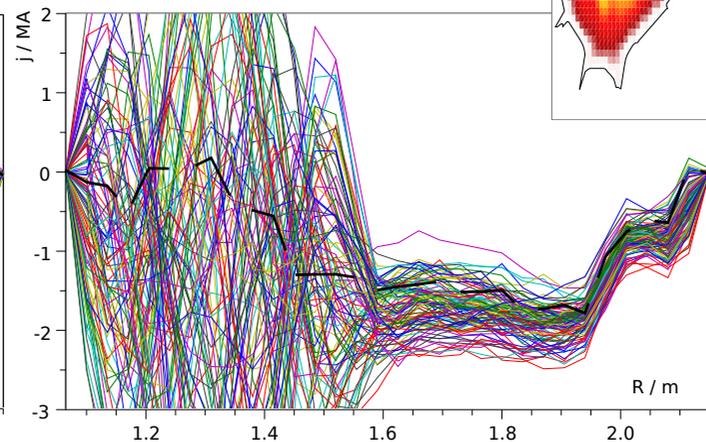
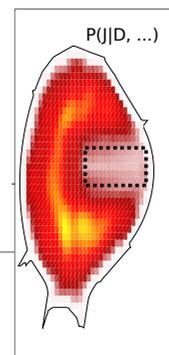
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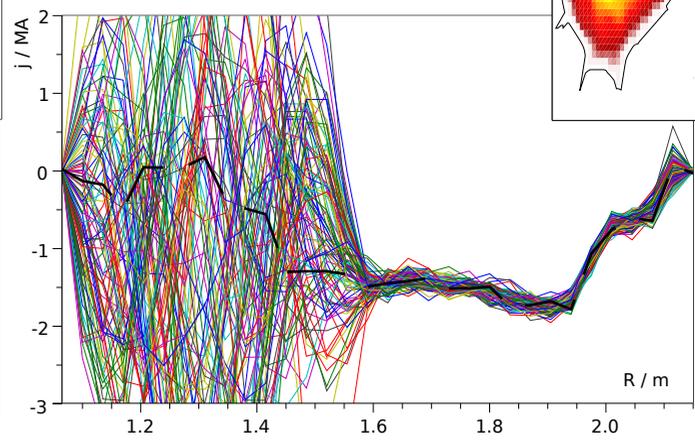
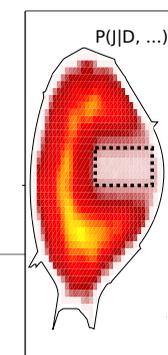
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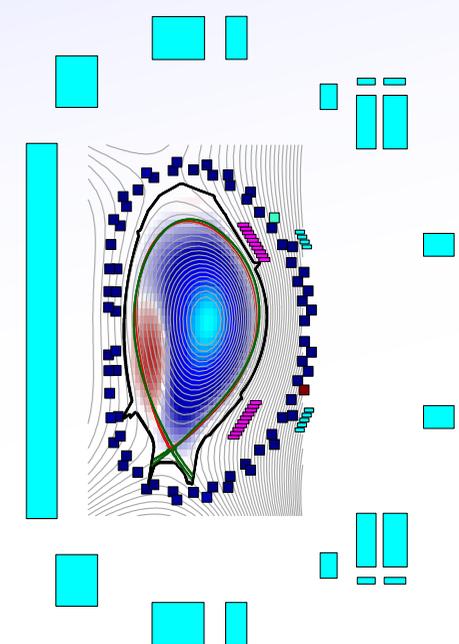
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30x15 grid of B_z
30x16 grid of Br .

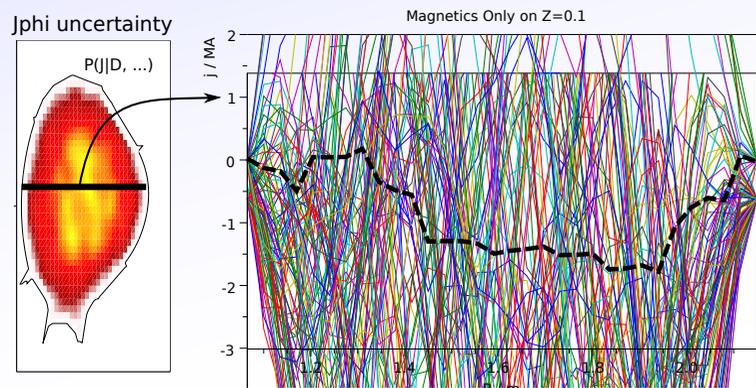


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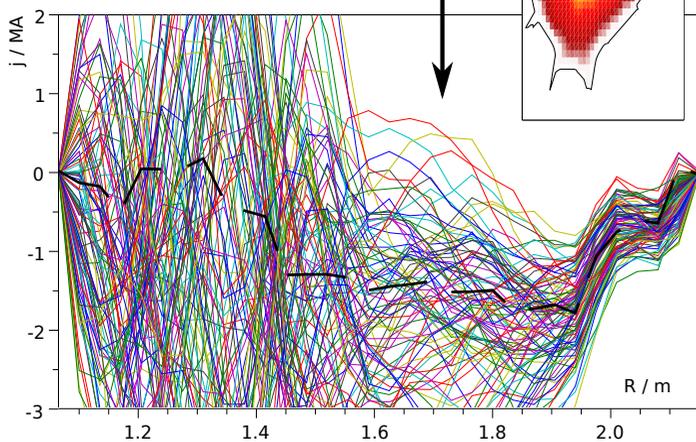
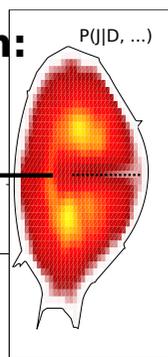


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Each case has 900 measurements at $\sigma = 10\text{mT}$.
So difference is only in the **type** of information.

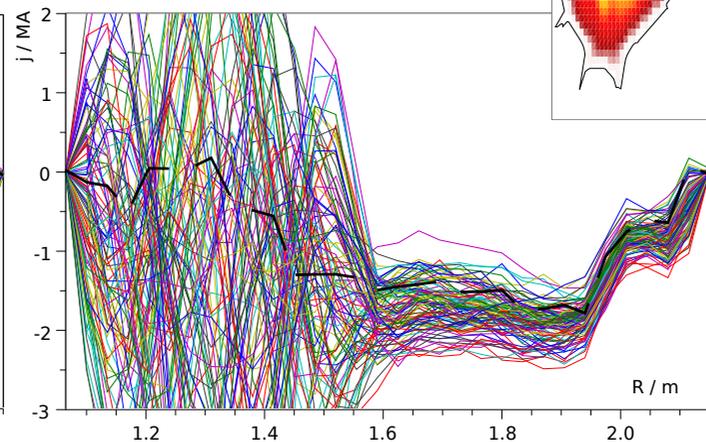
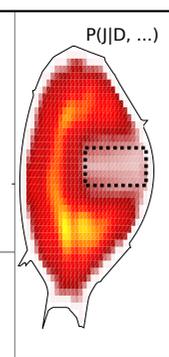
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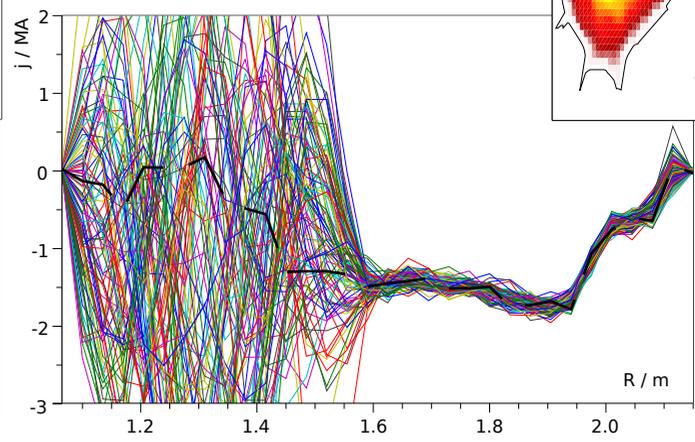
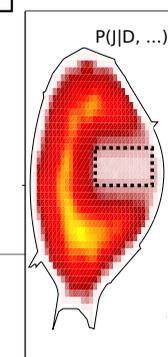
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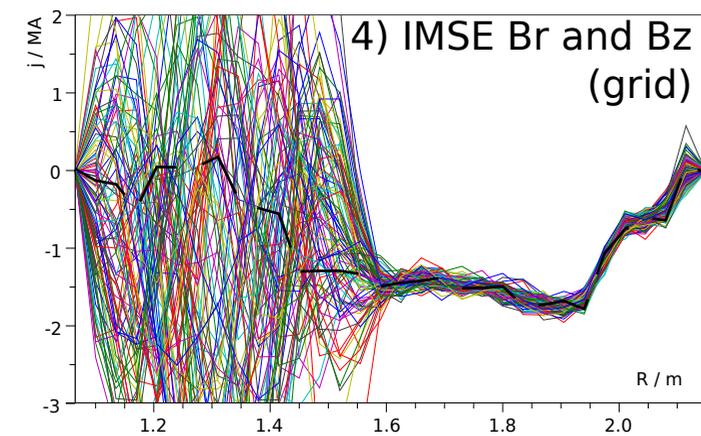
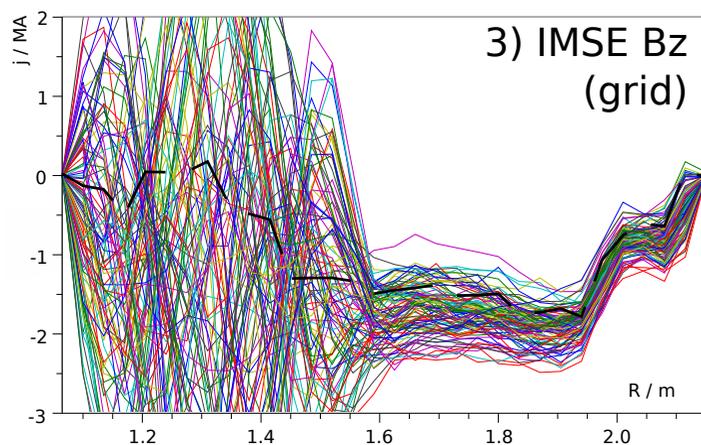
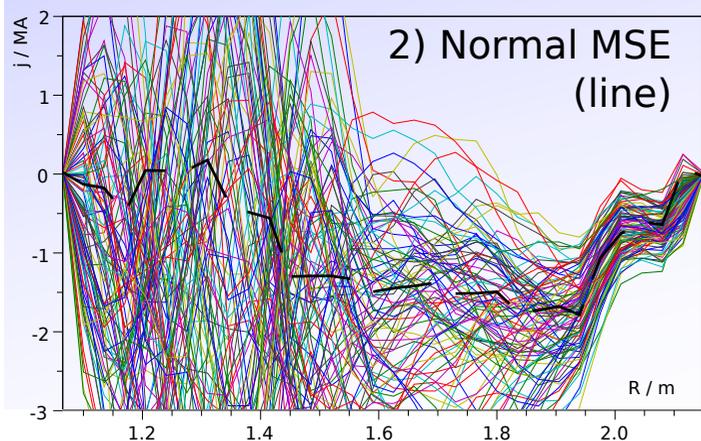




All $\sigma_{Br} = \sigma_{Bz} = 10\text{mT}$

By current tomography II

The IMSE still has a large uncertainty in j_ϕ offset. The unknown term it is not entirely pinned down by the magnetics.

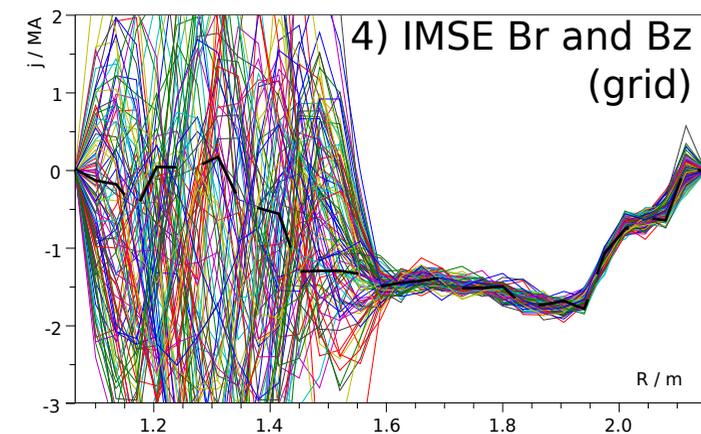
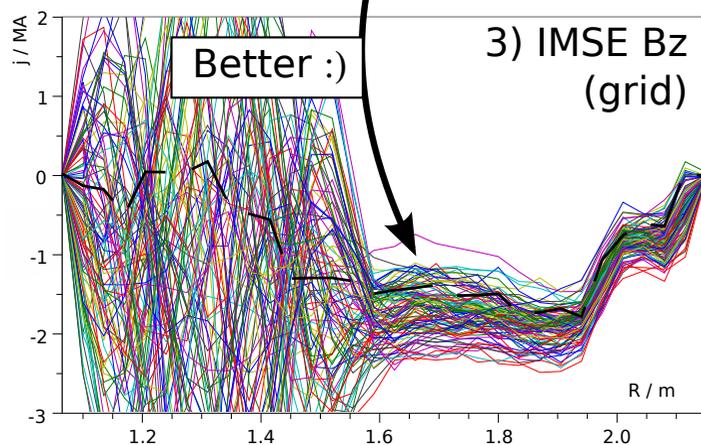
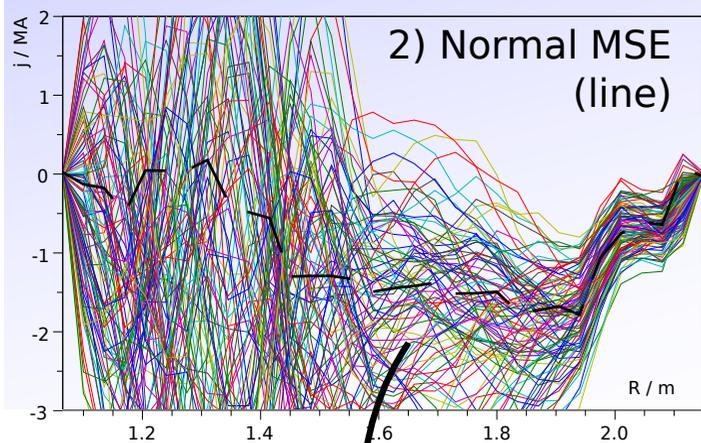


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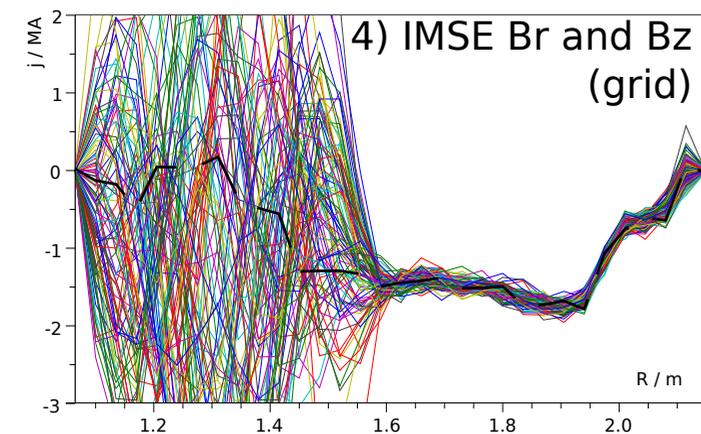
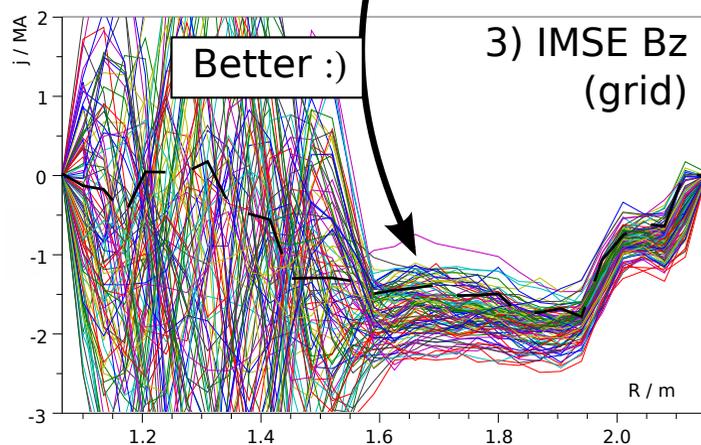
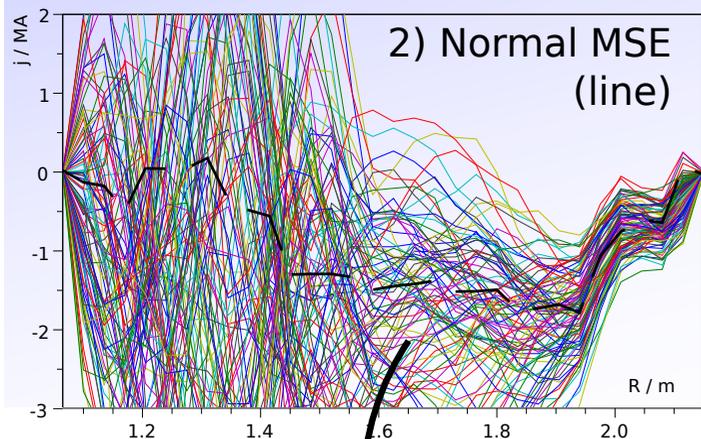
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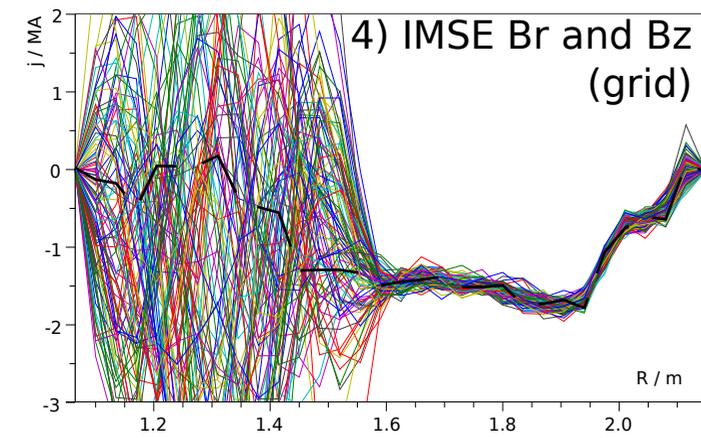
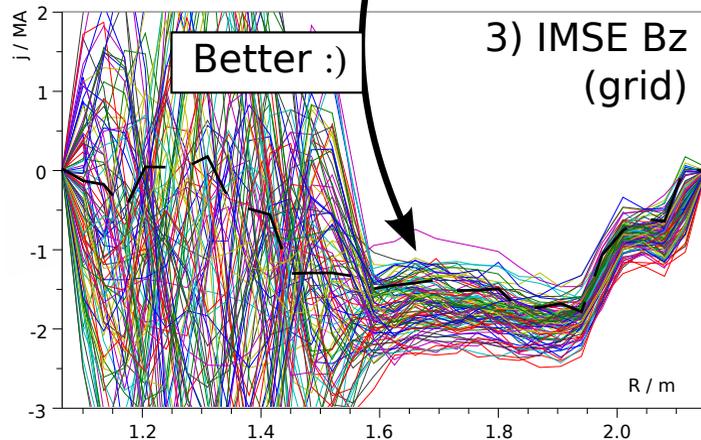
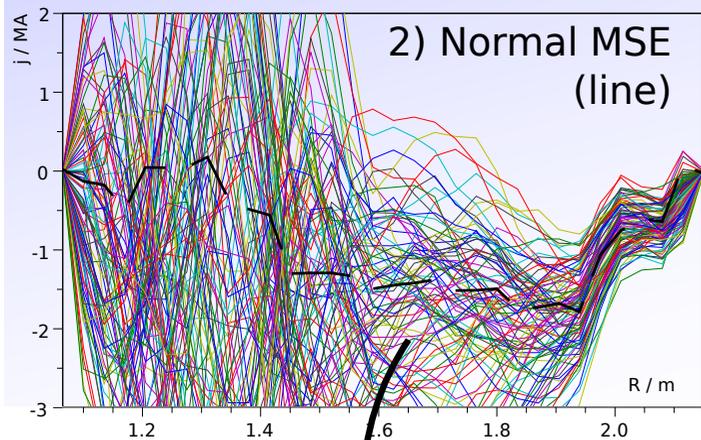
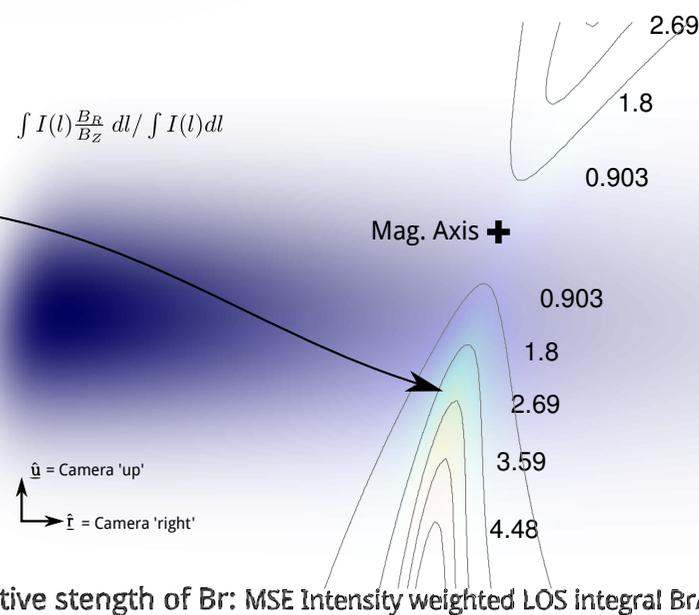
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All $\sigma_{Br} = \sigma_{Bz} = 10\text{mT}$

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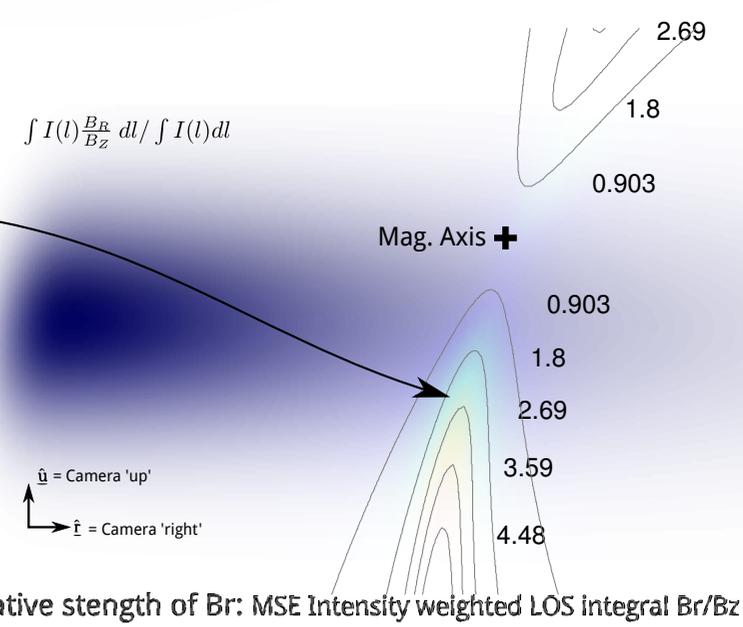
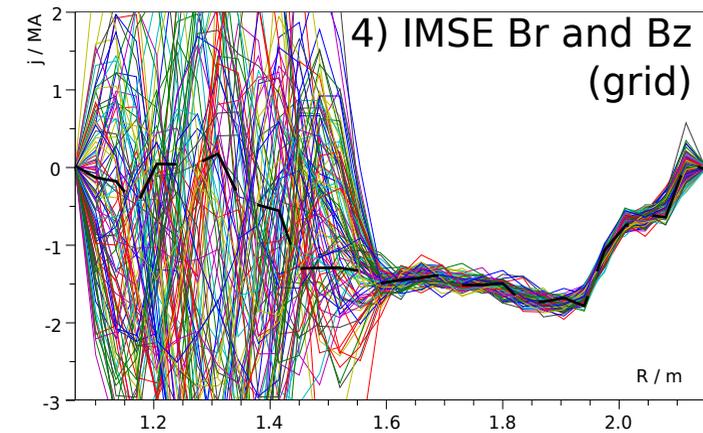
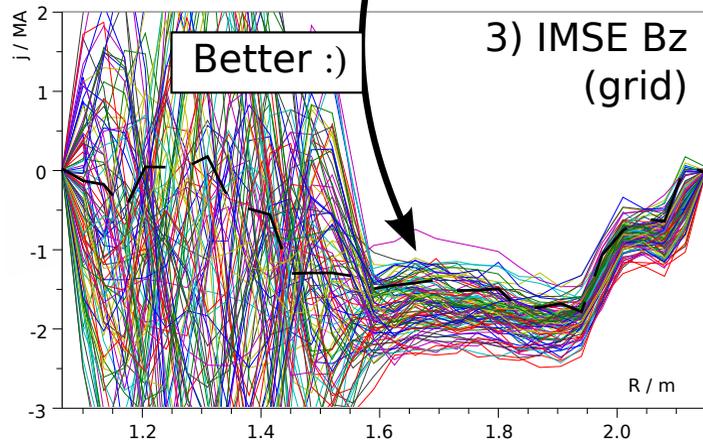
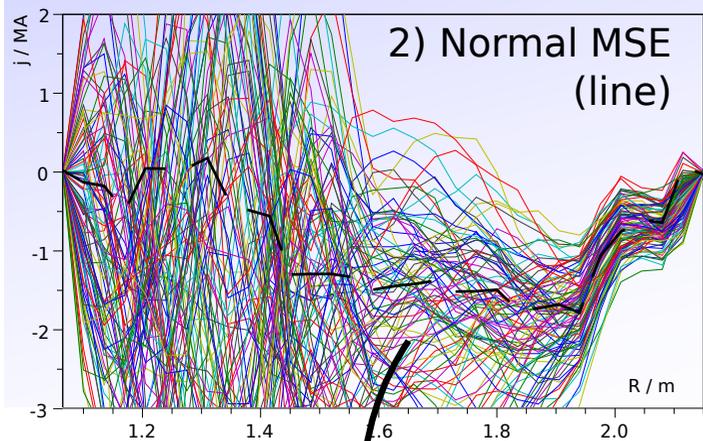
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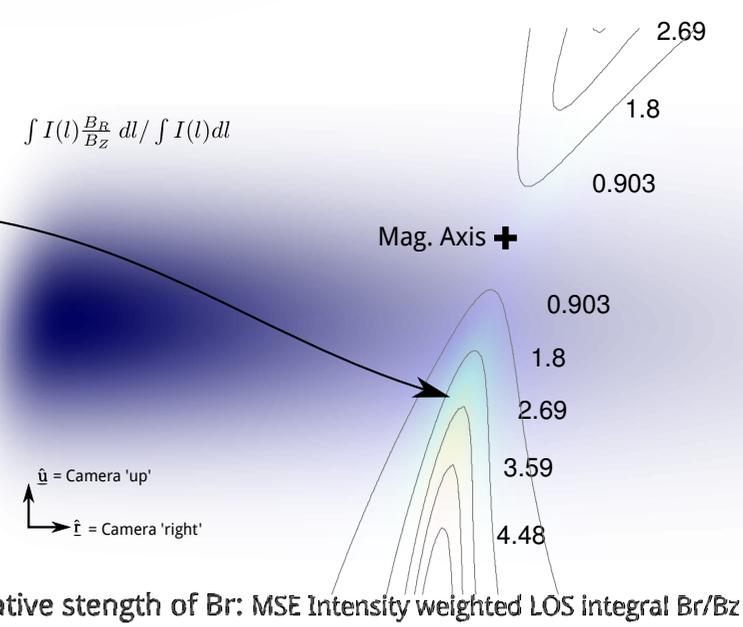
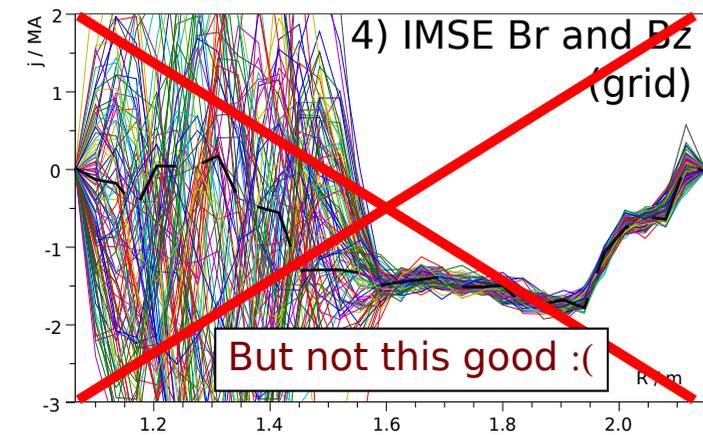
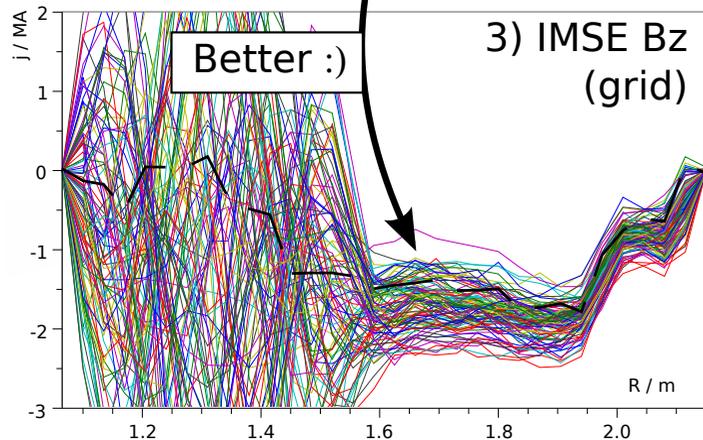
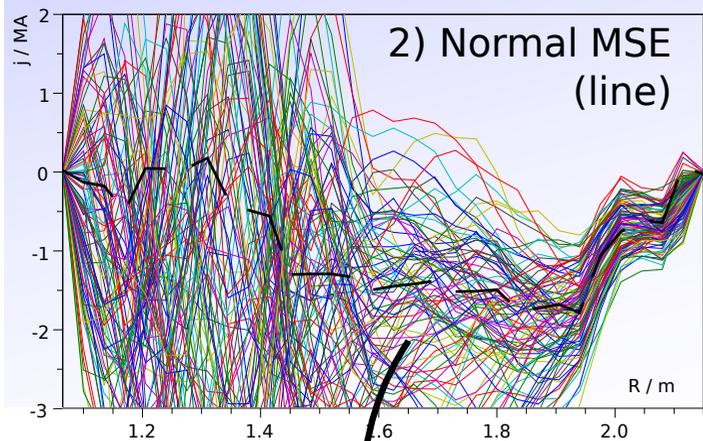
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At 5 - 10%, it will have an effect, but we do not expect to see the full current recovery from 2D tomography.





Full Inference

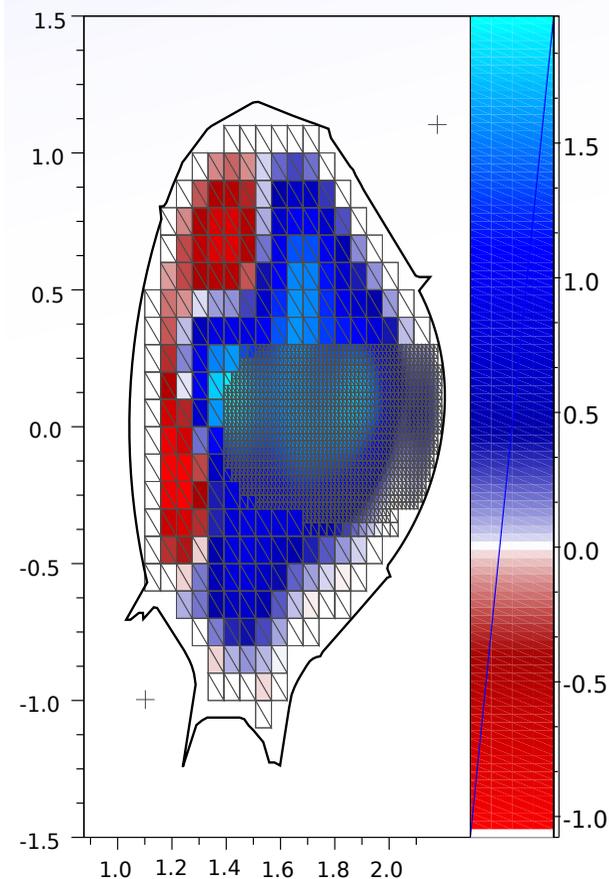
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Current beams

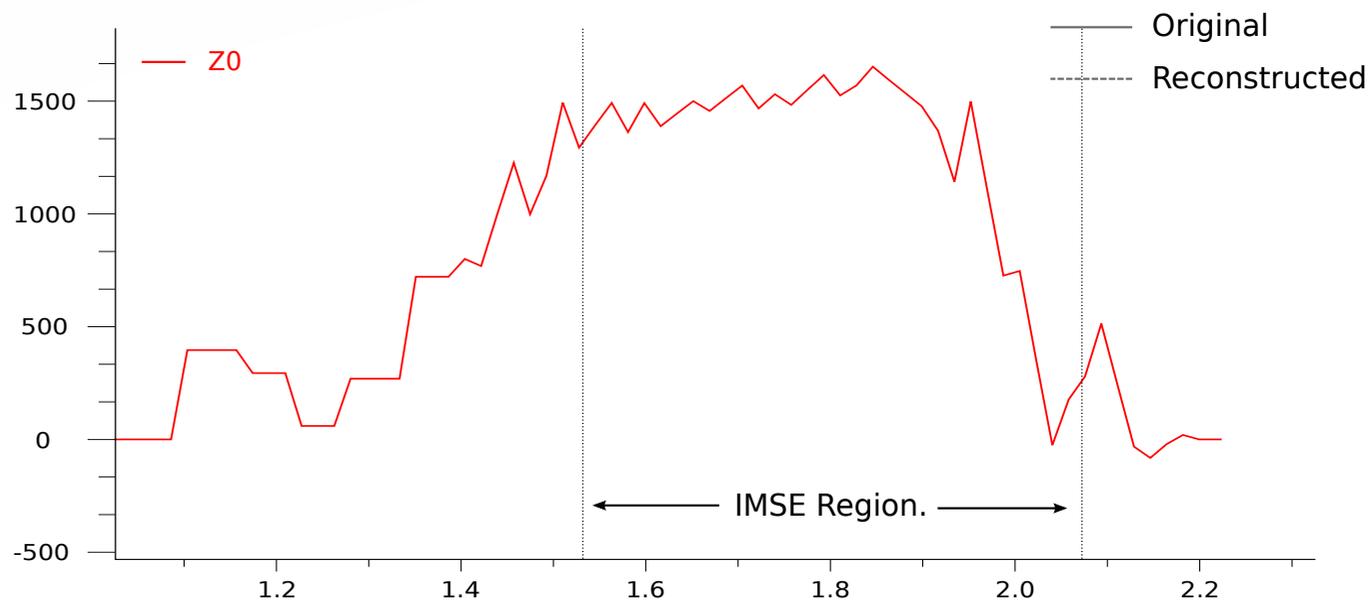
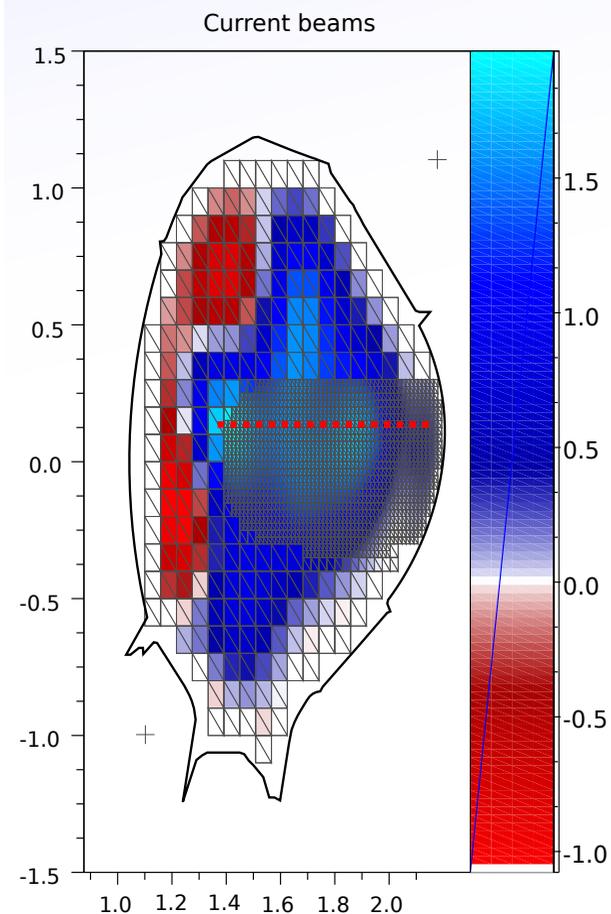


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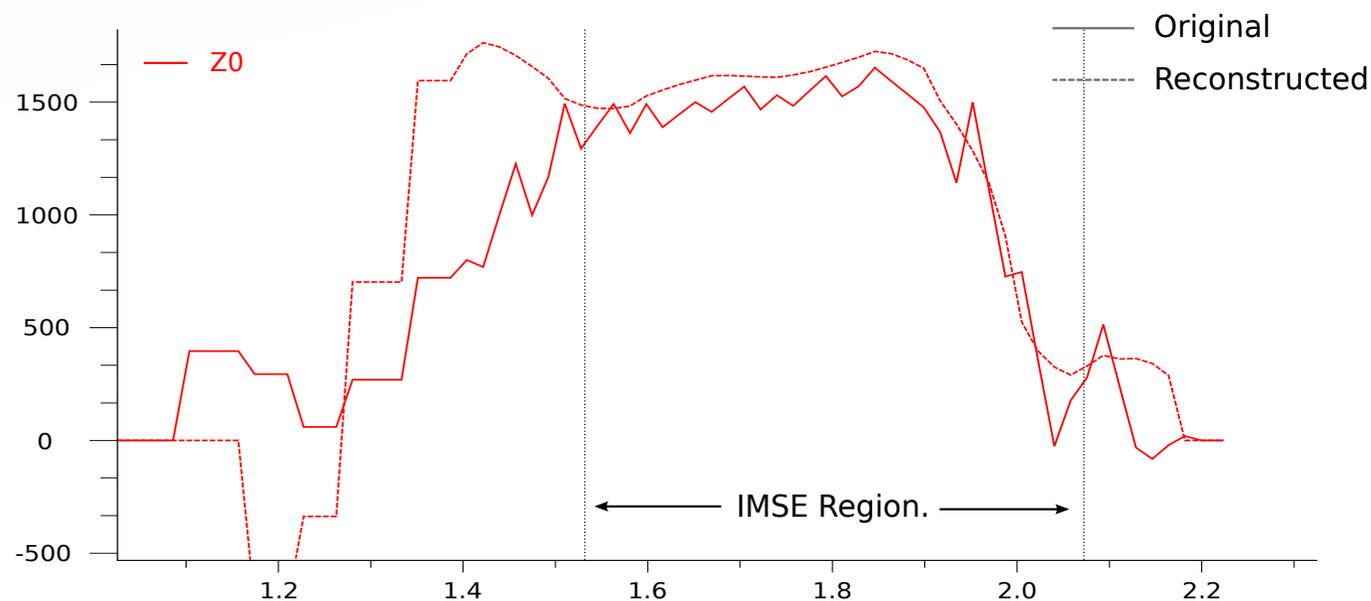
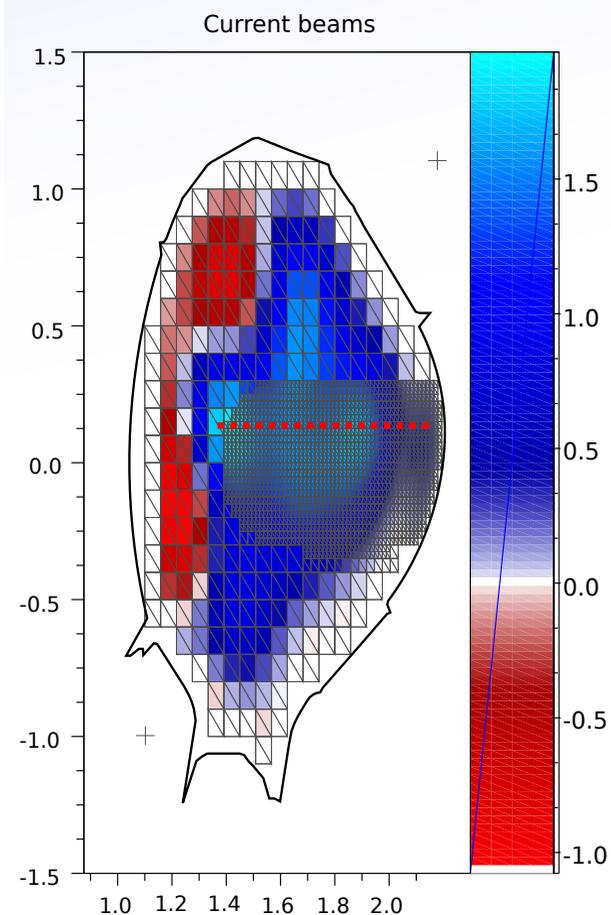


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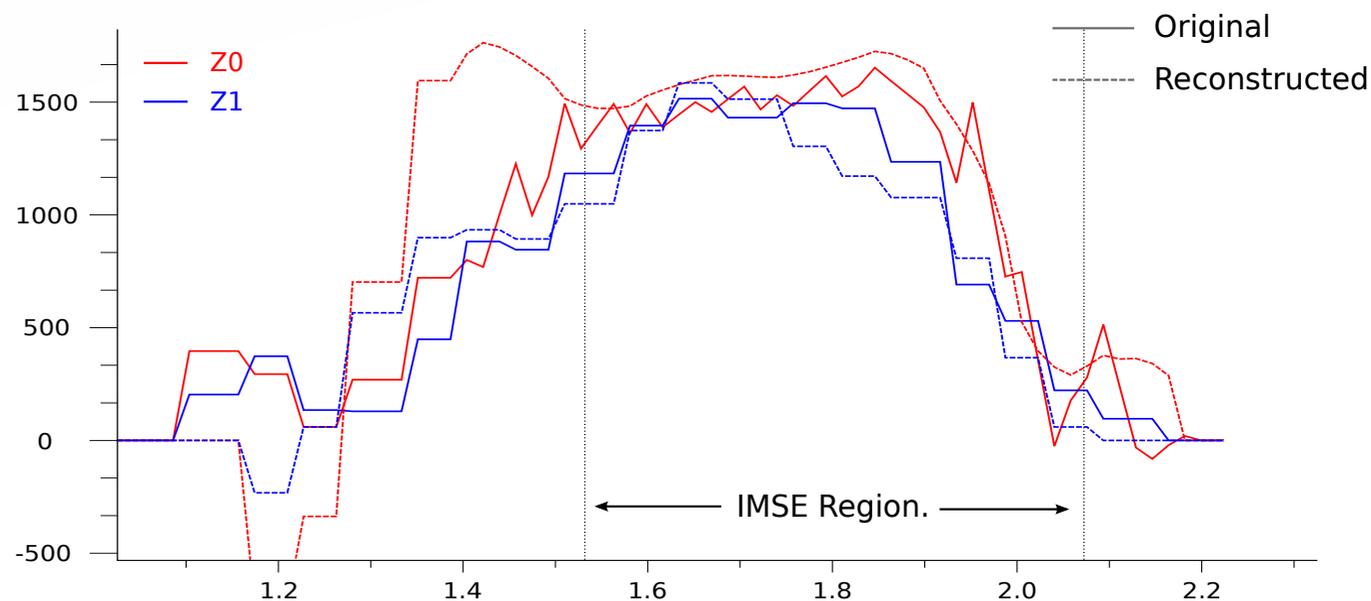
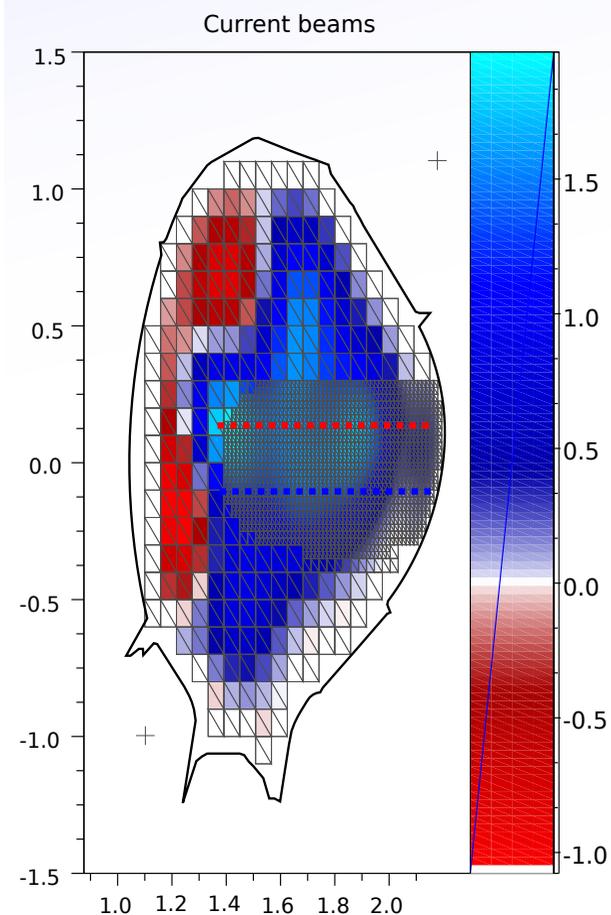


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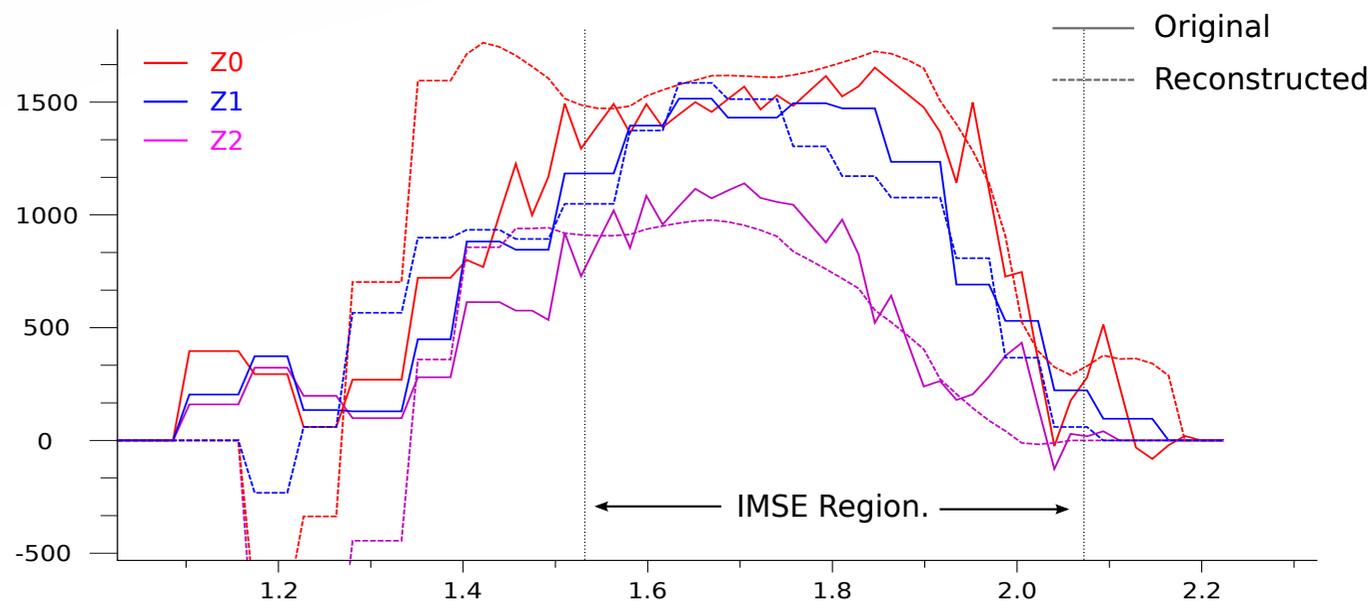
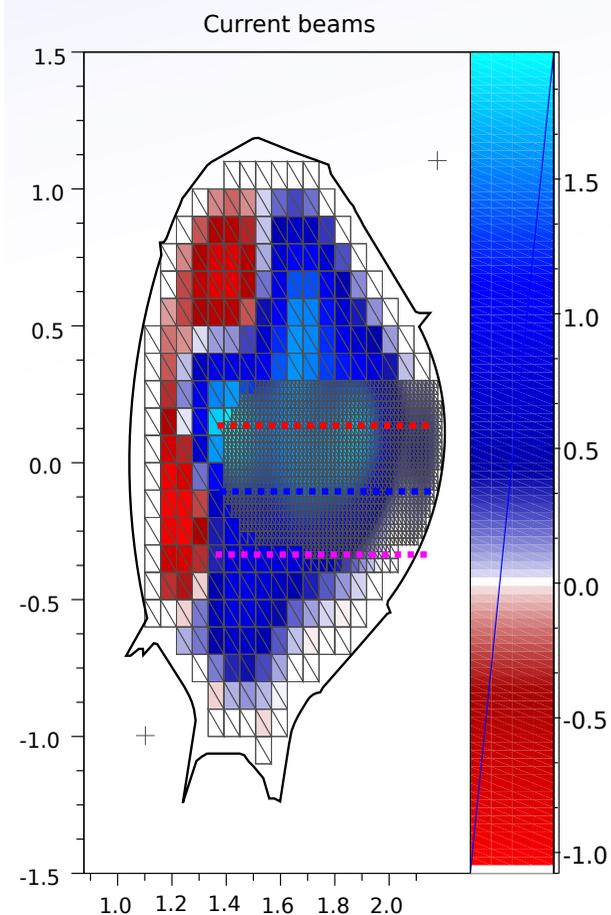


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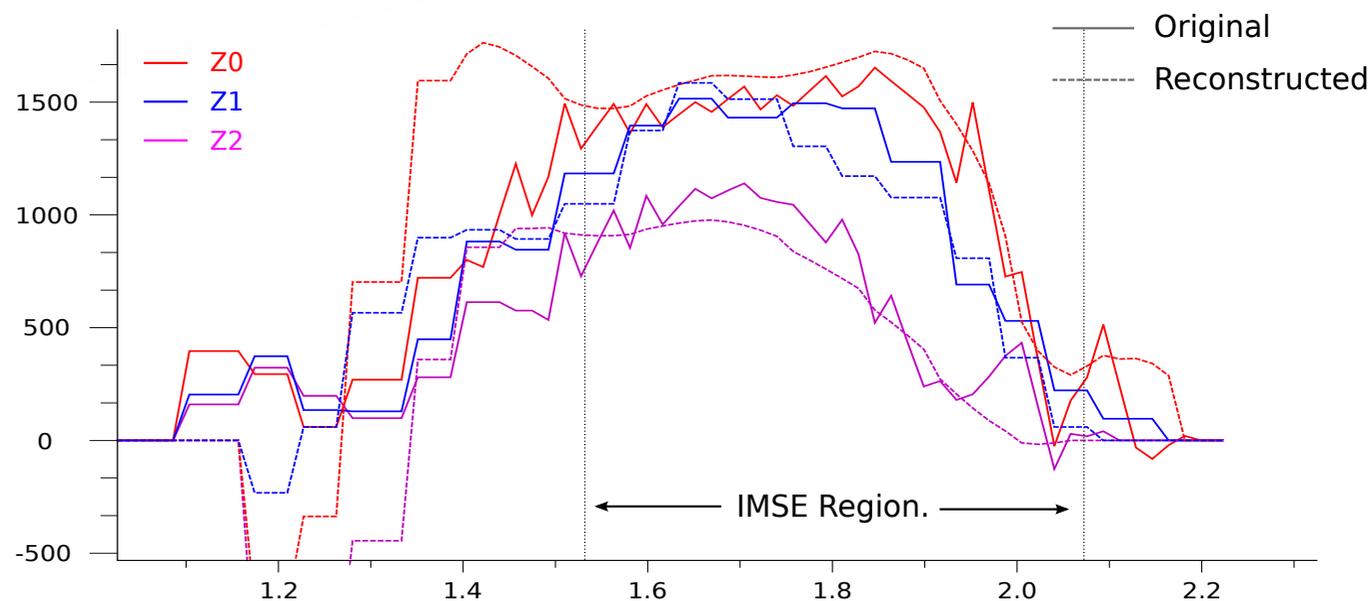
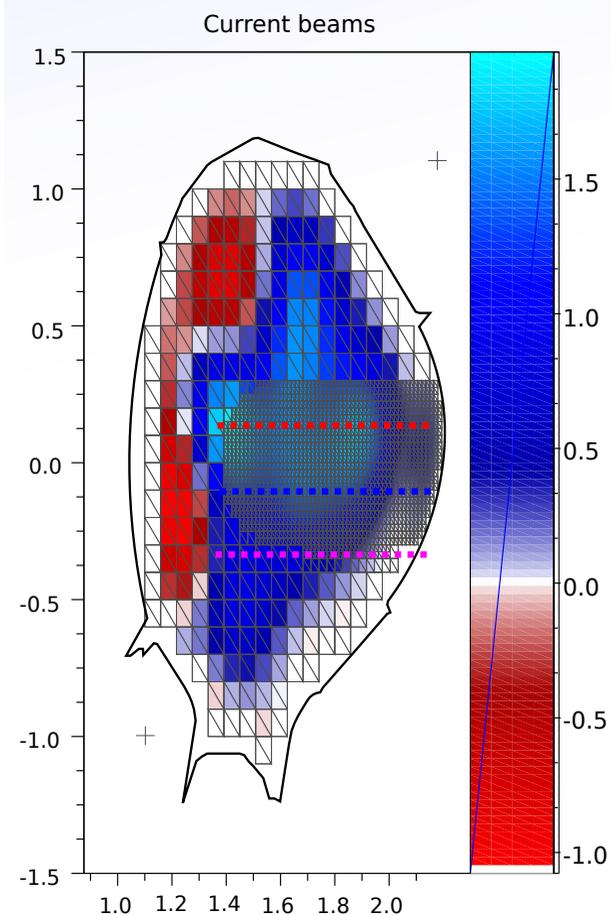


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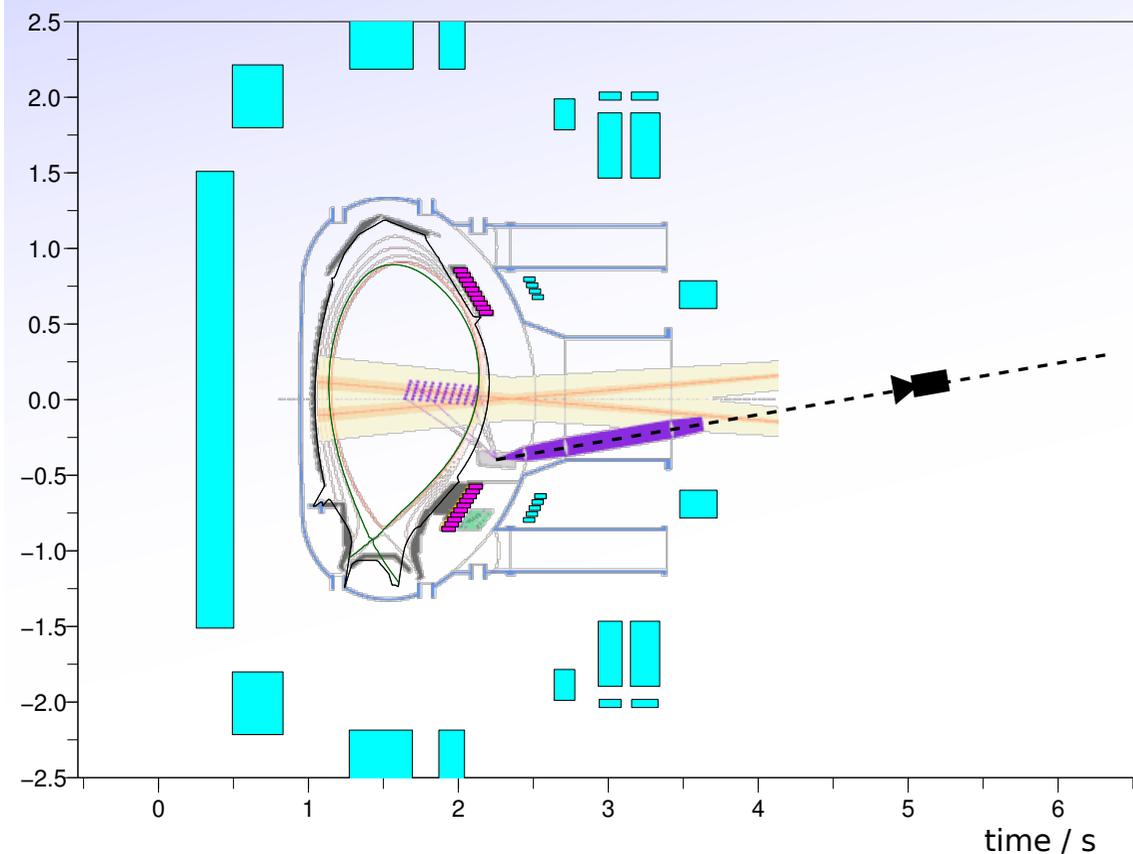
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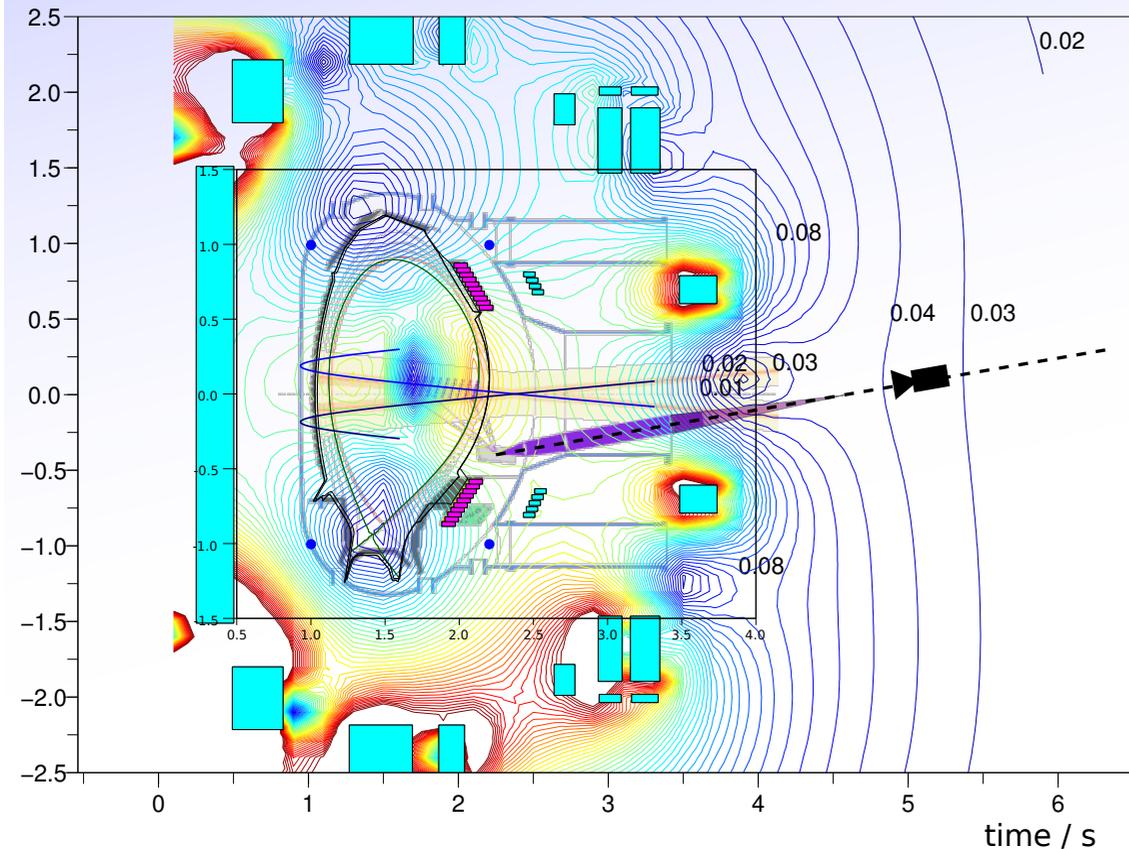
Initial results indicate that it is possible to recover the $j_\phi(R, Z)$ to at least a good resolution for studying the bulk plasma (e.g. testing different equilibrium models etc). This is MUCH better than is currently possible. Resolution is processing limited - Higher resolution may be possible, but computation cost rises with resolution as $\sim n^4$.

Other progress (Hardware)



Ideally, we want to fix the camera and optic plates directly to the viewing optics (no fibre etc).

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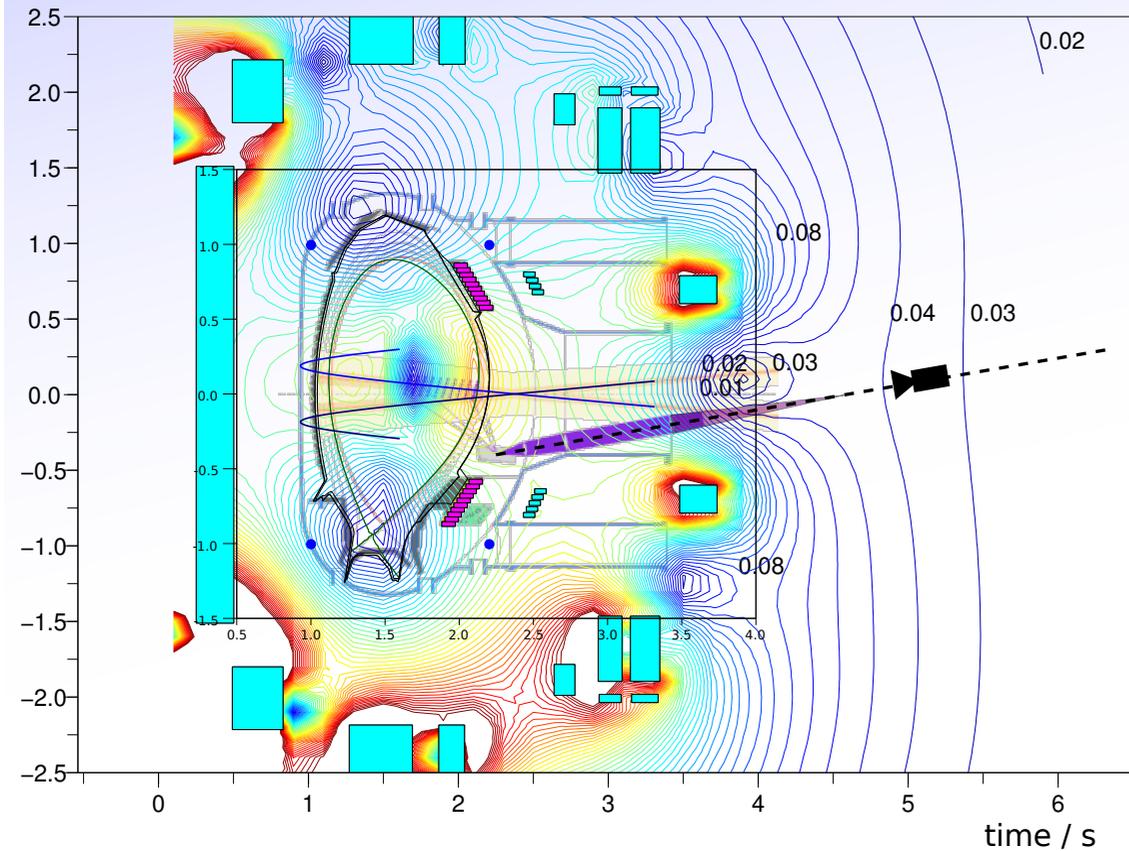


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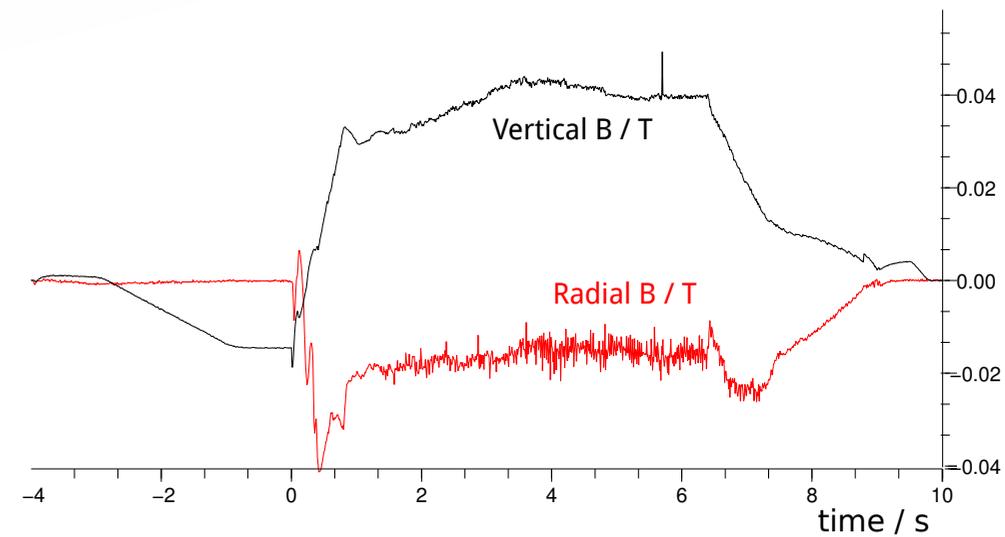
For the highest plasma current ($I_p = 1.2\text{MA}$),
 $|B| < 50\text{mT}$:

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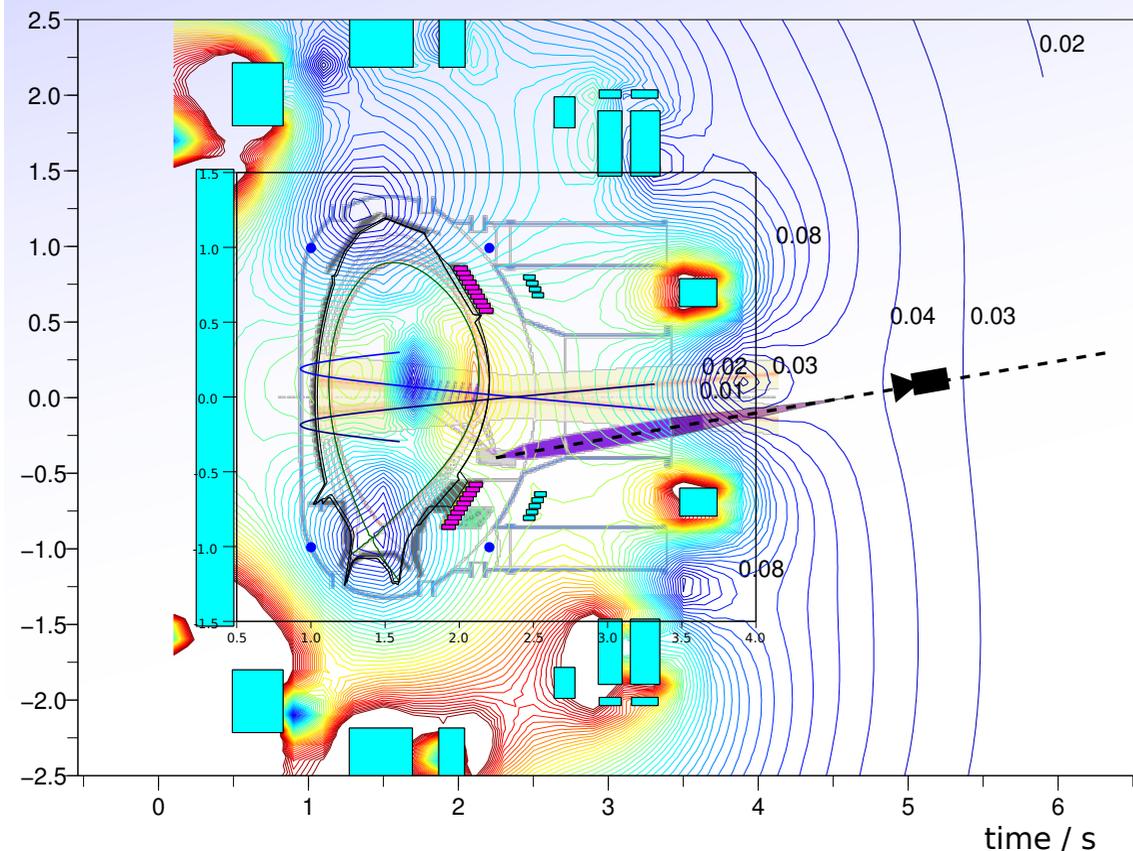


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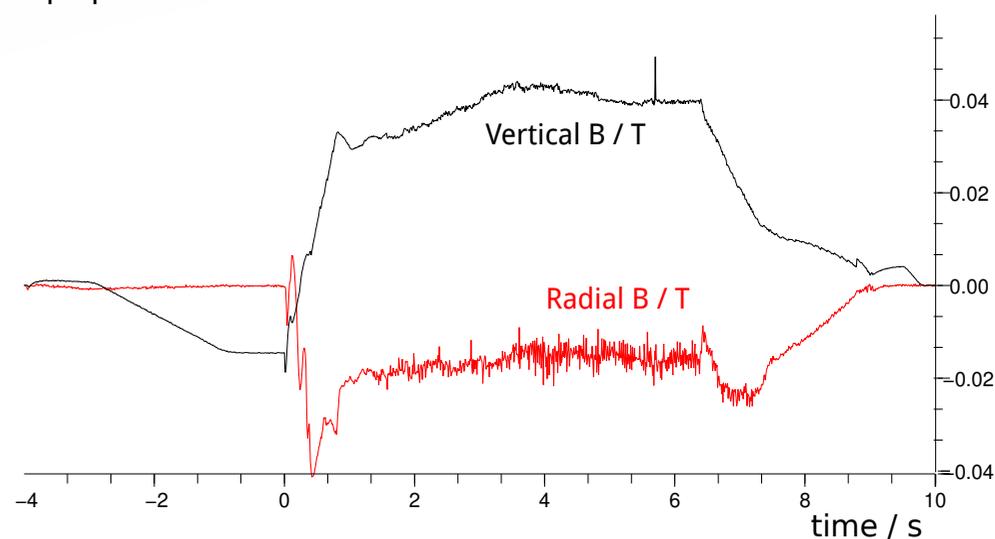


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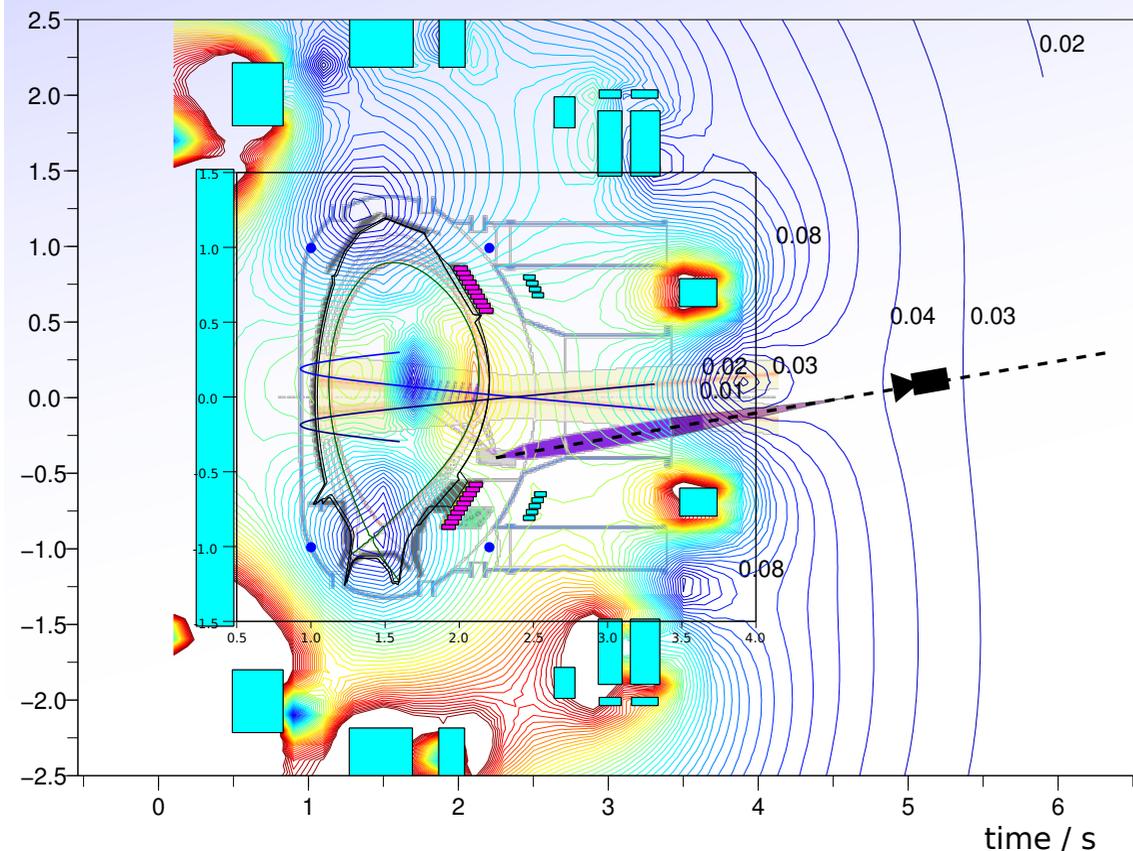
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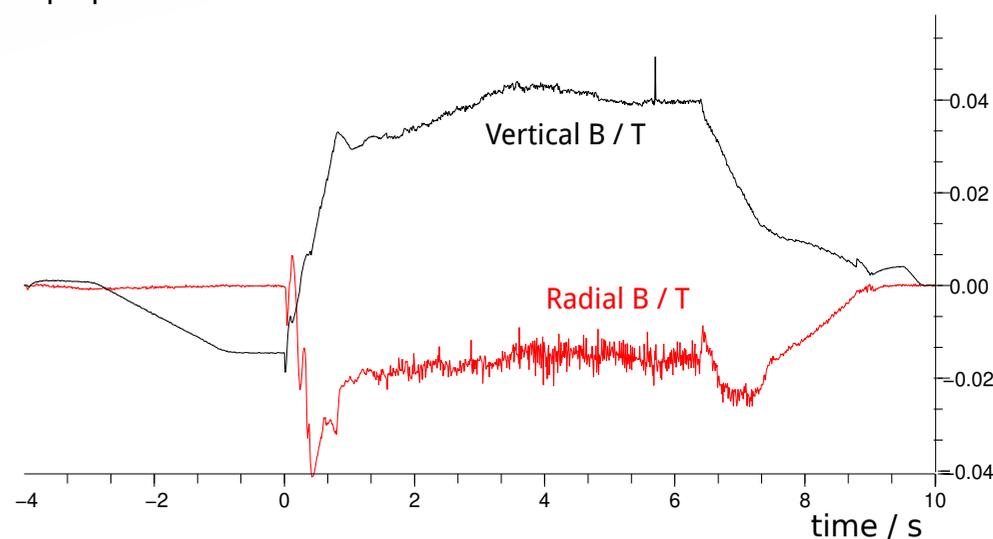
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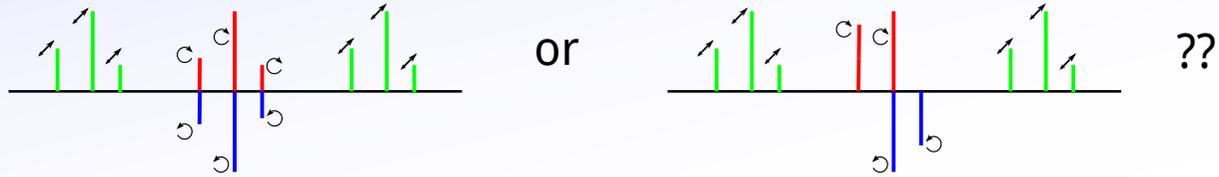
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- Faraday rotation due the field in the Savart plates will not be a problem, but the main delay plate might be. (I'm assuming Lithium Niobate, but I can't find a Verdet constant for it in the Literature. Any suggestions?)

Other progress (Model)

Various other effects have been corrected in the forward model:

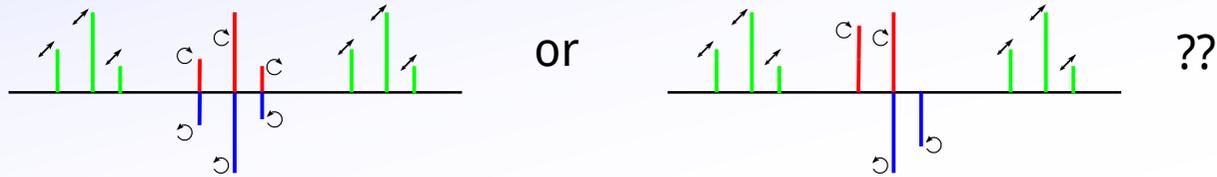
- Detail of Stark splitting and component polarisations. (Thanks to R. Reimer for pointing this out)



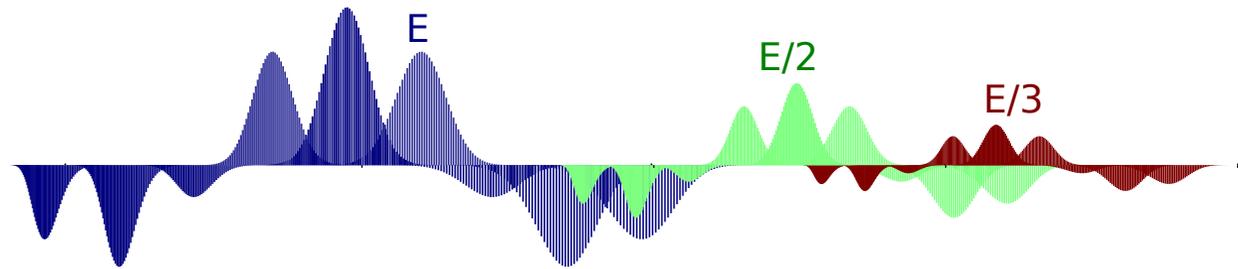
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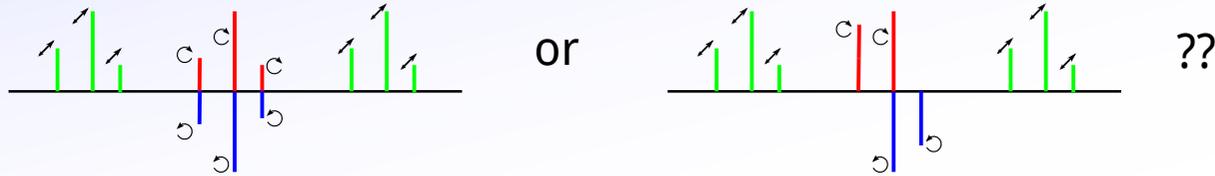
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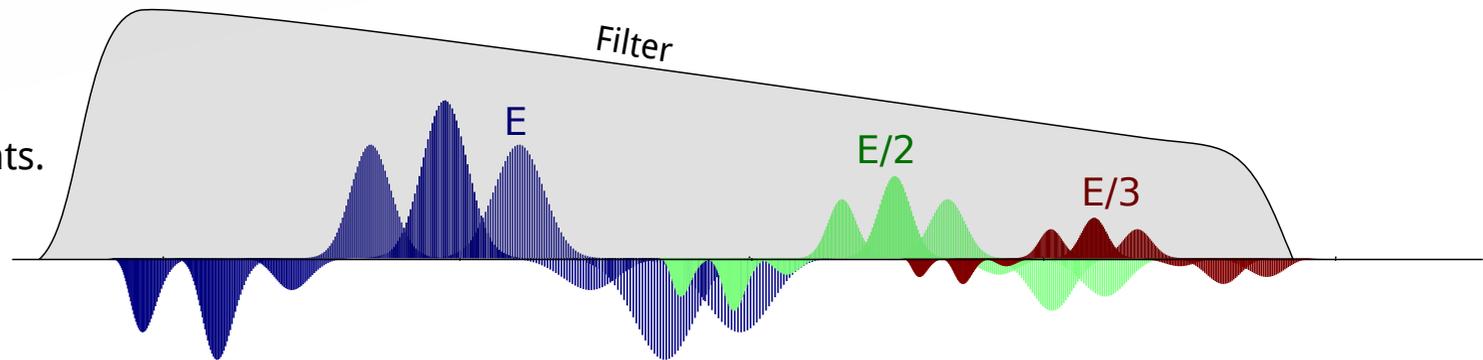
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- Non-uniform filter pass-band.

- Asymmetries in Stark components.

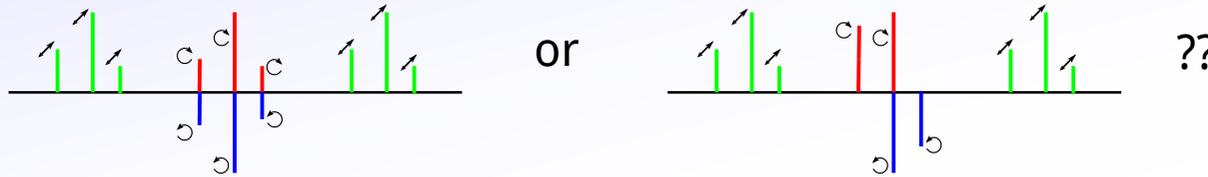


Some of these significantly effect the image phases, but the polarisation angle (from the amplitude) remains unaffected.

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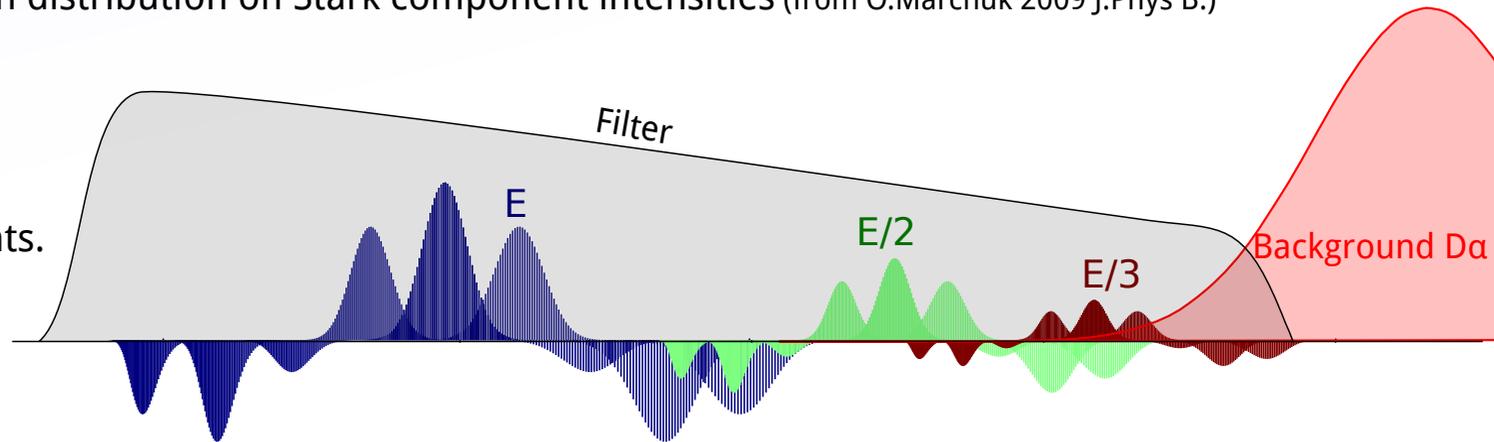
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Things still to add:

- Background D-Alpha and FIDA. (These will only reduce S/N).
- CCD noise (other than photon statistics).
- Viewing optics.